

Non-adiabatic Molecular Dynamics using Hefei-NAMD

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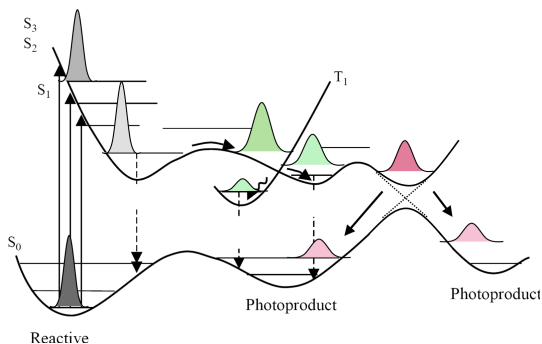
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Beyond Born-Oppenheimer Approximation

$$\Psi(\mathbf{r}, \mathbf{R}, t) = \Omega_j(\mathbf{R}, t)\Phi_j(\mathbf{r}; \mathbf{R}); \quad \hat{\mathcal{H}}_{el}(\mathbf{r}; \mathbf{R})\Phi_j(\mathbf{r}; \mathbf{R}) = E_j(\mathbf{R})\Phi_j(\mathbf{r}; \mathbf{R})$$



Method beyond BO approximation

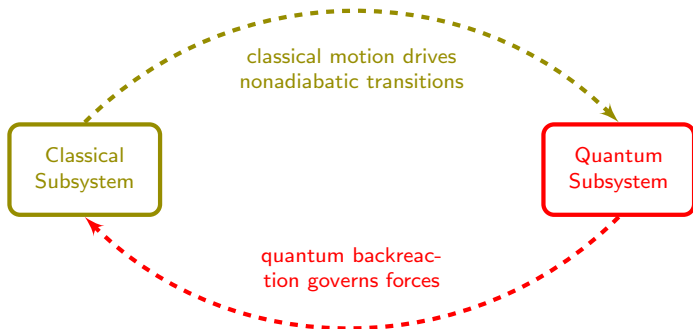
Quantum Dynamics

- MCTDH

Mixed Quantum-Classical

- Trajectory surface hopping (TSH)
- Ehrenfest dynamics
- ...

$$i\hbar \frac{\partial \Phi(\mathbf{r}, \mathbf{R}, t)}{\partial t} = \hat{\mathcal{H}}_{el}(\mathbf{r}, \mathbf{R}) \Phi(\mathbf{r}, \mathbf{R}, t); \quad \Phi(\mathbf{r}, \mathbf{R}, t) = \sum_j C_j(t) \psi_j(\mathbf{r}, \mathbf{R})$$



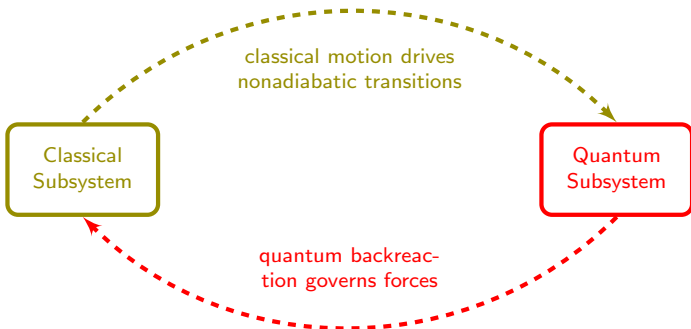
Surface Hopping

$$M_\alpha \ddot{\mathbf{R}}_\alpha = -\nabla_\alpha E_k^{el}(\mathbf{R})$$

Ehrenfest Dynamics

$$M_\alpha \ddot{\mathbf{R}}_\alpha = -\nabla_\alpha \langle \hat{\mathcal{H}}_{el}(\mathbf{r}, \mathbf{R}) \rangle$$

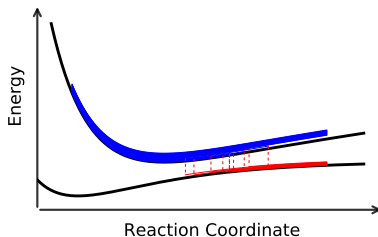
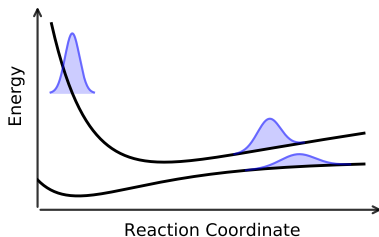
$$i\hbar \frac{\partial \Phi(\mathbf{r}, \mathbf{R}, t)}{\partial t} = \hat{\mathcal{H}}_{el}(\mathbf{r}, \mathbf{R}) \Phi(\mathbf{r}, \mathbf{R}, t); \quad \Phi(\mathbf{r}, \mathbf{R}, t) = \sum_j C_j(t) \psi_j(\mathbf{r}, \mathbf{R})$$



Surface Hopping

$$M_\alpha \ddot{\mathbf{R}}_\alpha = -\nabla_\alpha E_k^{el}(\mathbf{R})$$

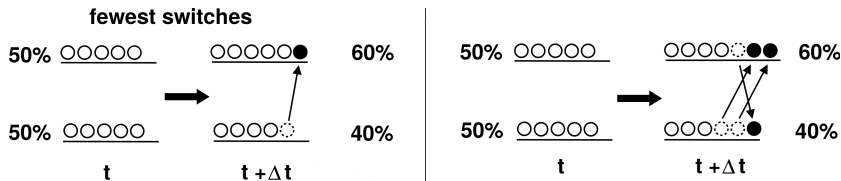
Trajectory Surface Hopping



Basic Idea

- Electronic wavefunction evolves according to time-dependent Schrödinger equation.
 - $\Phi(\mathbf{r}, \mathbf{R}, t) = \sum_j C_j(t) \psi_j(\mathbf{r}, \mathbf{R}) \Rightarrow i\hbar \dot{C}_j(t) = \sum_k C_k(t) [H_{jk} - i\hbar \dot{\mathbf{R}} \cdot \mathbf{d}_{jk}]$
 - $H_{jk} = \langle \psi_j | \hat{\mathcal{H}}_{el} | \psi_k \rangle$
 - Non-adiabatic Couplings (NAC): $\mathbf{d}_{jk} = \langle \psi_j | \nabla_{\mathbf{R}} | \psi_k \rangle$
- An ensemble of **independent** nuclear trajectories is considered.
- Each trajectory propagates on **one single** potential energy surface at any given time.
 - $M_{\alpha} \ddot{\mathbf{R}}_{\alpha} = -\nabla_{\alpha} E_k^{el}(\mathbf{R})$
- Hops of trajectories between electronic states is possible.
 - Tully's Fewest-Switches algorithm et al.

Fewest-Switches Algorithm



Assumptions

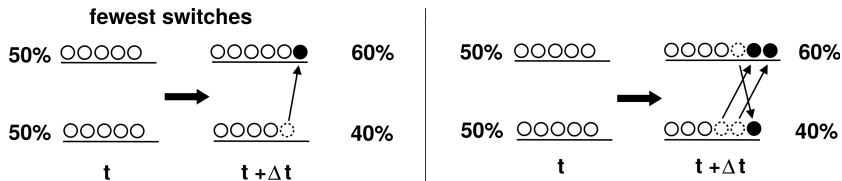
- Ensemble of independent trajectories have same coefficients $C_j(t)$.
- Internal consistency condition $N_j(t) \propto C_j^*(t) C_j(t) = \rho_{jj}(t)$.
- Hops from j to different $k \neq j$ are independent.
- Overall trajectory hops should be minimum.

Fewest-Switches: Hopping Probability

Transition from current state j to state $k \neq j$ is allowed only if population of state j is **decreasing**.

$$P_{jk}(t, \Delta t) = \max\left(-\frac{2 \int_t^{t+\Delta t} dt \left[\hbar^{-1} \text{Im}(\rho_{jk} H_{jk}) - \text{Re}(\rho_{jk} \mathbf{d}_{jk} \cdot \dot{\mathbf{R}}) \right]}{\rho_{jj}}, 0\right)$$

Fewest-Switches Algorithm

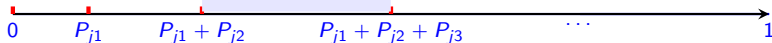


Assumptions

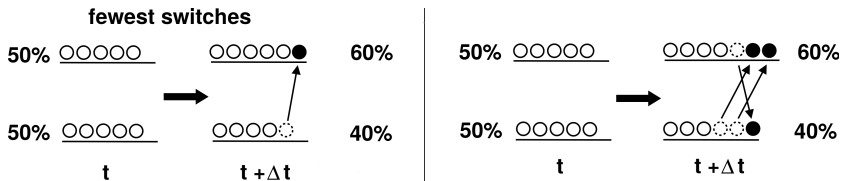
- Ensemble of independent trajectories have same coefficients $C_j(t)$.
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- Hops from j to different $k \neq j$ are independent.
- Overall trajectory hops should be minimum.

Fewest-Switches: Which state to hop

$$\sum_{l=1, l \neq j}^{k-1} P_{jl} < \xi < \sum_{l=1, l \neq j}^k P_{jl} \quad \text{then } j \rightarrow k$$



Fewest-Switches Algorithm



Assumptions

- Ensemble of independent trajectories have same coefficients $C_j(t)$.
- Internal consistency condition $N_j(t) \propto C_j^*(t) C_j(t) = \rho_{jj}(t)$.
- Hops from j to different $k \neq j$ are independent.
- Overall trajectory hops should be minimum.

Fewest-Switches: After trajectory hops

- Energy should be conserved after the hop e.g. by rescaling the nuclear velocity in the direction of non-adiabatic coupling.
 - $E_j + K_j = E_k + K_k \Rightarrow \mathbf{R}_\alpha^k = \mathbf{R}_\alpha^j - M_\alpha^{-1} \lambda \mathbf{d}_{jk}$
- Hops **rejected** when energy conservation not satisfied.

FSSH combined with TDDFT

single particle representation

$$\psi_p(\mathbf{r}, t) = \sum_j c_j(t) \phi_j(\mathbf{r}; \mathbf{R})$$

time-dependent Kohn-Sham equation

$$i\hbar \frac{\partial \psi_p(\mathbf{r}, t)}{\partial t} = \mathcal{H}(\mathbf{r}; \mathbf{R}) \psi_p(\mathbf{r}, t)$$

$$i\hbar \dot{c}_j(t) = \sum_k c_k(t) (\varepsilon_k \delta_{jk} - i\hbar d_{jk})$$

$$d_{jk} = \langle \phi_j | \frac{\partial}{\partial t} | \phi_k \rangle$$

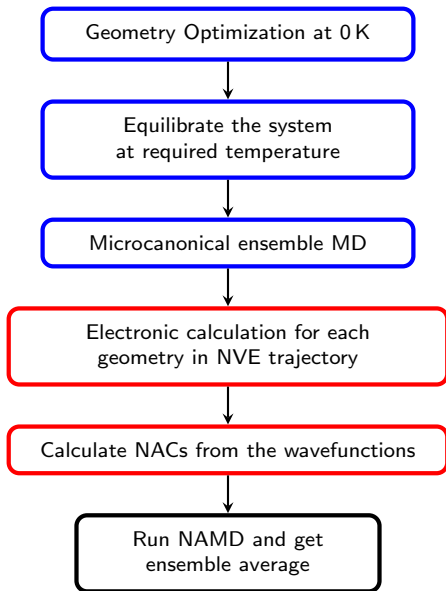
FSSH hopping probability

$$P_{jk}(t, \Delta t) = \max \left[\frac{2 \int_t^{t+\Delta t} dt [\text{Re}(\rho_{jk} d_{jk})]}{\rho_{jj}}, 0 \right]$$

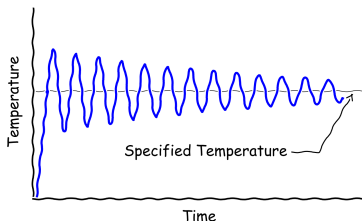
Classical-Path
Approximation

$$G_{jk} = P_{jk} B_{jk}; \quad B_{jk} = \begin{cases} \exp \left\{ -\frac{\varepsilon_j - \varepsilon_k}{k_B T} \right\}, & \text{if } \varepsilon_j > \varepsilon_k \\ 1, & \text{if } \varepsilon_j \leq \varepsilon_k \end{cases}$$

Work Flow



- 1 Geometry optimization: to find an energy minimum configuration.
 - Supercell should be large enough because BZ is only sampled at Γ point.
- 2 Molecular dynamics equilibration run — canonical ensemble,



usually by velocity rescaling method, e.g. set $SMASS = -1$ in VASP.¹

- considered equilibrated when temperature fluctuation within $\pm 10\%$ of the specified temperature.
 - check the configuration before proceeding to next step.
- 3 Molecular dynamics production run — microcanonical ensemble.
 - The equilibrated configuration i.e. CONTCAR of the last step is used as input of this run.
 - The duration of MD is determined by the specific problem.

¹ <https://cms.mpi.univie.ac.at/wiki/index.php/SMASS>

Example INCARs for molecular dynamics using VASP

Example INCAR for equilibration run

```
General
SYSTEM = Your Job Name
PREC   = Med
ISPIN  = 1
ISTART = 0
ICHARG = 1
Ionic Relaxation
ISIF   = 2
EDIFFG = -1E-2
Electronic relaxation
NPAR   = 4
ISMEAR = 0
SIGMA  = 0.1
ALGO   = Fast
NELMIN = 4
NELM   = 120
EDIFF  = 1E-5
Molecular Dynamics
ISYM   = 0      # turn off symmetry
IBRION = 0
NSW    = 500    # No. of ionic steps
POTIM  = 1      # time step 1.0 fs
SMASS  = -1     # velocity rescaling
NBLOCK = 4      # velocity rescaled
          # every NBLOCK step
TEBEG  = 300    # start temperature
TEEND  = 300    # end temperature
Writing Flag
NWRITE = 1      # make OUTCAR small
LWAVE  = .FALSE.
LCHARG = .FALSE.
```

Example INCAR for production run

```
General
SYSTEM = Your Job Name
PREC   = Med
ISPIN  = 1
ISTART = 0
ICHARG = 1
Ionic Relaxation
ISIF   = 2
EDIFFG = -1E-2
Electronic relaxation
NPAR   = 4
ISMEAR = 0
SIGMA  = 0.1
ALGO   = Fast
NELMIN = 4
NELM   = 120
EDIFF  = 1E-6
Molecular Dynamics
ISYM   = 0      # turn off symmetry
IBRION = 0
NSW    = 5000   # NSW*POTIM fs
POTIM  = 1      # time step 1.0 fs
SMASS  = -3     # Microcanonical
NBLOCK = 1      # XDATCAR contains
          # positions of each step
Writing Flag
NWRITE = 1      # make OUTCAR small
LWAVE  = .FALSE. # WAVECAR not needed
LCHARG = .FALSE. # CHG not needed
```

- ④ Extract the atomic positions from OUTCAR or XDATCAR (NBLOCK should be 1). This can be easily done by a python script using ASE², e.g.

```
#!/usr/bin/env python
import os
from ase.io import read, write

# This may take a little while, since OUTCAR may be large for long MD run.
CONFIGS = read('/path/to/OUTCAR', format='vasp-out', index=':')

#####
# or you can read positions from XDATCAR, which is faster since it only contains
# the coordinates. You may have to process the XDATCAR beforehand when VASP
# forget to write some lines. Executing the following line on shell prompt.
#####
# sed -i 's/^\s*$/Direct configuration=/' XDATCAR
#####
# CONFIGS = read('/path/to/XDATCAR', format='vasp-xdatcar', index=':')

NSW      = len(CONFIGS)           # The number of ionic steps
NSCF     = 2000                  # Choose last NSCF steps for SCF calculations
NDIGIT   = len("{:d}".format(NSCF)) #
PREFIX   = './run/'             # run directories
DFORM    = "%/%0%dd" % NDIGIT    # run directories format
for ii in range(NSCF):          # write POSCARs
    p = CONFIGS[ii - NSCF]
    r = (PREFIX + DFORM) % (ii + 1)
    if not os.path.isdir(r): os.makedirs(r)
    write('{:s}/POSCAR'.format(r), p, vasp5=True, direct=True)
```

SCF calculations II

- 2 With each extracted configuration, perform an SCF calculation to obtain the WAVECAR files for subsequent NAMD calculations. Example parallel job script for this run:

```
START=1          # start configuration
END=2000         # end configuration
for i in `seq ${START} ${END}`
do
  (( j = i - 8 ))
  ii='printf "%04d" $i'
  jj='printf "%04d" $j'
  if [[ -d "run/${ii}" ]]; then
    cd run/${ii}
    if [[ -f RUNNING || -f ENDED ]]; then # Job still running or has ended, skip it
      cd ../../
      continue
    fi
    touch RUNNING
    echo "####_RUNNING_in_DIR:_RUN/${ii}"
    sleep 0.4
    [[ -s ../${jj}/CHGCAR ]] && cp ../${jj}/CHGCAR .
    $OPEN_MPI $IB_FLAG -np $NCPUS -machinefile ${HOSTFILE} $VASP_EXEC # VASP RUN
    if grep 'Total_CPU' OUTCAR >& /dev/null; then
      touch ENDED
    else
      rm ENDED 2> /dev/null
    fi
    rm RUNNING CHG vasprun.xml
    cd ../../
  fi
done
```

²<https://wiki.fysik.dtu.dk/ase/>

- ① Prepare input files for NAMD calculations, for example,

inp

```
&NAMDPARA
  BMIN      = 325      ! bottom band index
  BMAX      = 340      ! top band index
  NBANDS    = 388      ! number of bands

  NSW       = 2000
! number of ionic steps
  POTIM     = 1        ! MD time step
  TEMP      = 100      ! temperature

  NSAMPLE   = 100      ! number of samples
  NAMDTIME  = 1000     ! time for NAMD run
  NELM      = 1000     ! electron time step
  NTRAJ     = 5000     ! SH trajectories
  LHOLE     = .FALSE.  ! hole/electron SH

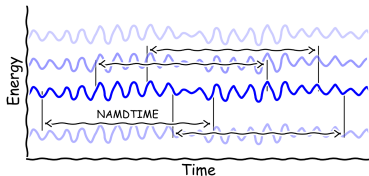
  RUNDIR    = "/path/to/run_dir/"
! LCPEXT   = .TRUE.
/ ! NOTE THE SLASH HERE!
```

INICON

```
87 329
519 330
13 330
533 329
541 329
542 329
558 329
59 329
62 329
65 329
...

```

- Column 1: initial time.
- Column 2: initial band index.
- No. of rows: NSAMPLE in inp.



- ② The program will first calculate non-adiabatic couplings (NACs) and store the NACs in the file COUPCAR. If there is already the COUPCAR, then the program will read COUPCAR and perform NAMD calculations.

The key quantity in the NAMD is the non-adiabatic couplings (NACs) and can be calculated by finite difference method.

$$\begin{aligned} \mathbf{d}_{jk} &= \langle \phi_j | \frac{\partial}{\partial \mathbf{t}} | \phi_k \rangle = \langle \phi_j | \nabla_{\mathbf{R}} | \phi_k \rangle \cdot \dot{\mathbf{R}} \\ &= \frac{\langle \phi_j(\mathbf{t}) | \phi_k(\mathbf{t} + \Delta \mathbf{t}) \rangle - \langle \phi_j(\mathbf{t} + \Delta \mathbf{t}) | \phi_k(\mathbf{t}) \rangle}{2\Delta \mathbf{t}} \end{aligned} \quad (1)$$

- $\phi_j(\mathbf{t})$ and $\phi_j(\mathbf{t} + \Delta \mathbf{t})$ are Kohn-Sham (KS) pseudo-wavefunctions (WFC) from WAVECAR at different time steps. Note that $\phi_j(\mathbf{t})$ are not orthonormalized for PAW PS.
- For Gamma-only VASP, $\phi_j(\mathbf{t})$ are real functions. In addition, there is an arbitrary ± 1 phase factor for each KS-WFC, which may affect the evaluation of Eq. (1).
- The inner product in the Eq. (1) can be obtained simply by summing over the products of plain-wave coefficients³

$$\langle \phi_j | \phi_k \rangle = \sum_{\mathbf{G}} c_j^*(\mathbf{G}) \cdot c_k(\mathbf{G}) \quad (2)$$

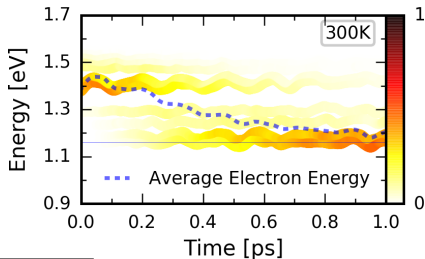
³For Gamma-only WFCs, the plain-wave coefficients other $G = 0$ are multiplied by a factor of $\sqrt{2}$.

Output files & postprocessing

- There are two types of output files: SHPROP.* and PSICT.*, where the suffix corresponds to the initial time in INICON, i.e. column 1.
 - PSICT.*: time-dependent propagation coefficients $c_i(t)$.
 - SHPROP.*: population of the selected states. The final results are obtained by averaging over these files.

```
#####  
#!/usr/bin/env python  
import numpy as np  
from glob import glob  
  
inFiles = glob('/path/to/SHPROP.*')  
outData = np.array([np.loadtxt(inf) for inf in inFiles])  
np.savetxt('shprop_aver.dat', np.average(outData, axis=0))  
#####
```

- The following plot shows the time evolution of the state populations.⁴



⁴script can be found here: https://github.com/QijingZheng/UsefulPythonSnippet/blob/master/poen_fssh.py

Non-adiabatic and Adiabatic Charge Transfer

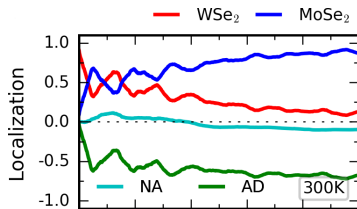
Define **charge localization** in a region V as

$$\begin{aligned}\int_V \rho_{PE}(\mathbf{r}, t) d\mathbf{r} &= \int_V |\psi_{PE}(\mathbf{r}, t)|^2 d\mathbf{r} \\ &= \sum_{ij} c_i^*(t) c_j(t) \int_V \phi_i^*(\mathbf{r}, \mathbf{R}) \phi_j^*(\mathbf{r}, \mathbf{R}) d\mathbf{r}\end{aligned}$$

The contribution to the change of the above term can be separated into two terms

$$\frac{d \int_V \rho_{PE}(\mathbf{r}, t) d\mathbf{r}}{dt} = \sum_{ij} \left\{ \frac{d(c_i^* c_j)}{dt} \int_V \phi_i^* \phi_j d\mathbf{r} + c_i^* c_j \frac{d \int_V \phi_i^* \phi_j d\mathbf{r}}{dt} \right\}$$

Example plot showing the NA/AD contribution to ET.⁵



⁵script can be found here: https://github.com/QijiangZheng/UsefulPythonSnippet/blob/master/spatial_localization_fssh.py

- The compilation of the code is very simple, no special libraries needed.

```
~/tmp/namd_compilation 34 files, 524Kb
zqj@node100 0 17:21:33 Sat Jun 09 $ make clean
rm -f *.mod *.a namd
rm -f prec.o lattice.o wave.o fileio.o couplings.o hamil.o TimeProp.o SurfHop.o main.o namd

~/tmp/namd_compilation 15 files, 100Kb
zqj@node100 0 17:21:36 Sat Jun 09 $ make
gfortran -g -O2 -c prec.f90
gfortran -g -O2 -c lattice.f90
gfortran -g -O2 -c wave.f90
gfortran -g -O2 -c fileio.f90
gfortran -g -O2 -c couplings.f90
gfortran -g -O2 -c hamil.f90
gfortran -g -O2 -c TimeProp.f90
gfortran -g -O2 -c SurfHop.f90
gfortran -g -O2 -c main.f90
gfortran -g -O2 -o namd prec.o lattice.o wave.o fileio.o couplings.o hamil.o TimeProp.o SurfHop.o main.o

~/tmp/namd_compilation 34 files, 524Kb
zqj@node100 0 17:21:39 Sat Jun 09 $ ll namd
-rwxr-xr-x 1 zqj penguin 96682 Jun 9 17:21 namd
```

- Yunhai Li⁶ from Southeast University also write an MPI parallelized version of the code.

⁶liyunhai1016@hotmail.com

Thank you!