

应用 Liquid Crystal, XY model, FIBA, QWZ model

图像：几何

Topo 不变量

$\pi_d(X)$
典型空间
简单 {
 S^1
Lie group
 RP^d, CP^d

自发对称破缺

Kibble-Zurek 机制
quench

π_1 kink

Fiber Bundle 纤维丛

$\psi(\vec{x})$

H_a^2

底空间
 $x, k \cdot$

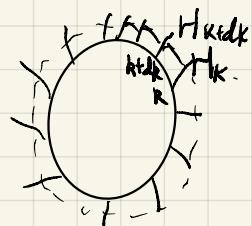
$\psi(\vec{x}) \in C^n$
 $H_a^2 \in A_{n \times n}$
 \vec{x}
Fiber

$\pi_d(M) = \pi_d(U/H)$

H : little group +

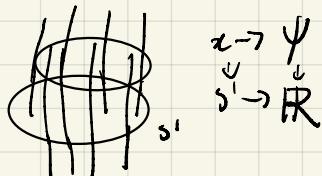
isotropic group

指向弱手

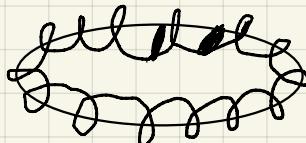
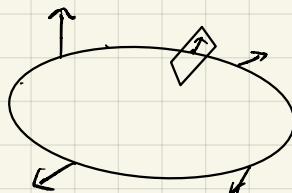


$\psi_k \in R$

$s' \rightarrow R$



$s' \rightarrow s'$



图像：几何

$$S^1 \rightarrow S^1 \Rightarrow \pi_1(S^1) = \mathbb{Z}$$

$$\mathbb{Z}^d = \mathbb{Z}^n$$



$$\vec{x} \in S^d \quad \psi_{\vec{x}} \in X$$

$\psi_{\vec{x}}$ 随空间 (\vec{x}, \vec{r}, \dots) 变化而变化 \Rightarrow 几何

空间： $\vec{x} \rightarrow$ 空间， M 是什么？

① Liquid crystal $\vec{n}(\vec{x}) \quad \vec{n}: \vec{x} \rightarrow S^d, RP^d$

② QM $\psi_{\vec{x}} \in \mathbb{C}, \mathbb{C}^n, BEC$, Ground state

③ Topo $H_K \sim U(N+M)/U(N) \times U(M)$

④ SSB $M \cong G/H$

SSB

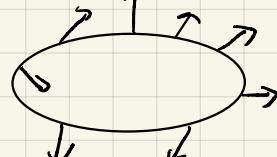
$\vee \xrightarrow{\text{quench}} \cup$

$M \cong G/H$

$\uparrow \downarrow$

墨西哥帽

$s(x): \vec{x} \rightarrow \mathbb{Z}_2$



$\overline{T\uparrow\uparrow J\downarrow\downarrow}$

局部基态

整体不是

几个结论

正合:

$$\pi_d(\mathcal{U}/\mathcal{H}) \sim \pi_d(\mathcal{U}), \pi_d(\mathcal{H})$$

$A \rightarrow B \rightarrow C$ 是正合的

正合序列 exact sequence

短正合 short

$$0 \rightarrow A \rightarrow B \rightarrow 0$$

长正合 long

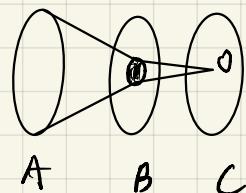
$$0 \rightarrow A \rightarrow B \rightarrow C \rightarrow 0$$

由许多短正合序列表构成

$$0 \rightarrow A_1 \rightarrow B_1 \rightarrow 0 \rightarrow A_2 \rightarrow B_2 \rightarrow C_2 \rightarrow 0 \rightarrow \dots$$

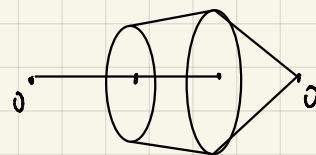
前的像 = 后的核

$$\text{Im}(A \rightarrow B) = \ker(B \rightarrow C)$$



$$0 \rightarrow A \rightarrow B \rightarrow 0 \Rightarrow A \cong B$$

同构



$$0 \rightarrow A \rightarrow B \rightarrow C \rightarrow 0 \Rightarrow B \cong A \otimes C$$

$$B/A \cong C$$

同态基本定理: $f: A \rightarrow B$

$$\text{Im } f = A / \ker f$$

找 $\pi_d(\mathcal{U}/\mathcal{H})$

查看		
H	\mathcal{U}	$M \cong (\mathcal{U}/H)$
π_3	$\pi_3(H) \rightarrow \pi_3(\mathcal{U}) \rightarrow \pi_3(M)$	
π_2	$\pi_2(H) \rightarrow \pi_2(\mathcal{U}) \rightarrow \pi_2(M)$	
π_1	$\pi_1(H) \rightarrow \pi_1(\mathcal{U}) \rightarrow \pi_1(M)$	正合序列
π_0		

$$\rightarrow \pi_d(H) \rightarrow \pi_d(\mathcal{U}) \rightarrow \pi_d(M) \rightarrow \pi_{d-1}(H) \rightarrow \pi_{d-1}(\mathcal{U}) \rightarrow \dots$$

几何相 (绝热过程) \Rightarrow Berry phase.

绝热过程: (热力学 $dU = dQ + dW$)

$dQ = 0$ adiabatic process

经典力学 \Rightarrow 绝热不变量 $I = \frac{1}{2\pi} \oint p dq$ 意义?

原子物理

QM

QM 中 $I = \frac{1}{2\pi} \oint p dq$

\downarrow

$$\frac{1}{2\pi} \oint \langle \psi | p | \psi \rangle dq$$

\downarrow

$$\frac{1}{2\pi} \oint \langle \psi | \frac{\partial}{\partial \vec{R}} | \psi \rangle d\vec{R}$$

期待
① 不变量

② 相移
③ 几何属性

$$\frac{1}{2\pi} \oint p dq = n \hbar \text{ 量子化}$$

Recall Bohr model

几何相
Berry phase

$$\frac{1}{2\pi} \oint p dq \stackrel{\text{stokes}}{=} \frac{1}{2\pi} \oint dP_i dq = \frac{\pi \hbar^2}{2\pi}$$

推导几何相: $H(\vec{x}, \vec{R}(t))$ 外部参数

$$\begin{cases} H(\vec{x}, \vec{R}(t)) \\ R(t) = R(t_0) \end{cases}$$

$$i \frac{\partial}{\partial t} \Psi(\vec{x}, t) = H \Psi(\vec{x}, t).$$

$$\text{e.g. } H = \frac{p^2}{2m} + \frac{1}{2} m (\omega + \alpha R)^2 x^2$$

$$H = \frac{p^2}{2m(\omega + \alpha R)} + \frac{1}{2} m \omega^2 x^2$$

$$H = (\vec{p} + \alpha \vec{R}) \cdot \vec{p}$$

Adiabatic process

$$H(\vec{x}, \vec{R}(t)) \Rightarrow \text{本征态}$$

\vec{x}

$$|\psi\rangle \rightarrow |\psi\rangle e^{i\theta}$$

1. 满足 Schrödinger eq.

不能跟 x 有关.

2. Adiabatic condition

$$|\Psi_n(\vec{x}, t)\rangle = |\Psi_n(\vec{x}, \vec{R}(t))\rangle e^{i\theta(t)}$$

$$\Rightarrow i \frac{\partial}{\partial t} |\Psi_n(\vec{x}, \vec{R}(t))\rangle e^{i\theta(t)} - \dot{\theta}(t) |\Psi_n(\vec{x}, \vec{R}(t))\rangle = H |\Psi_n\rangle e^{i\theta}$$

$$\Rightarrow \langle \Psi_n | i \frac{\partial}{\partial t} |\Psi_n\rangle - \dot{\theta} = \langle \Psi_n | H | \Psi_n \rangle = \lambda_n$$

$$\Rightarrow \dot{\theta} = \langle \varphi_n | i \frac{\partial}{\partial \theta} | \varphi_n \rangle - \lambda_n.$$

$$\Theta(T) - \Theta(0) = \int_0^T \dot{\theta} dt = - \int_0^T \lambda_n dt + i \int_0^T \langle \varphi_n | i \frac{\partial}{\partial \theta} | \varphi_n \rangle dt$$

$$\frac{\partial}{\partial \theta} = \frac{\partial}{\partial R} \varphi_n \cdot \frac{\partial R}{\partial \theta}$$

$$= - \int_0^T \lambda_n dt - T i \oint \langle \varphi_n | i \frac{\partial}{\partial R} | \varphi_n \rangle \cdot d\vec{R}$$

$$= - \int_0^T \lambda_n dt + \oint \langle \varphi_n | i \nabla_{\vec{R}} | \varphi_n \rangle \cdot d\vec{R}$$

動力学相位

几何相位

Berry phase

γ

圖像 1983

$$\text{作業: } H = B_x \sigma_3 + B_y \cos \theta \sigma_x + B_z \sin \theta \sigma_y$$

$\gamma = \text{solid angle}$

$$\left. \begin{array}{l} \textcircled{1} |\psi(\theta)\rangle \\ \textcircled{2} \frac{\partial}{\partial \theta} |\psi(\theta)\rangle \\ \textcircled{3} \int_0^{2\pi} \langle \psi(\theta) | i \frac{\partial}{\partial \theta} | \psi(\theta) \rangle d\theta \end{array} \right\} \text{立体角}$$