

(5.26)

KPZ 方程重整化

已讨论: ① Brown motion

② Linear E-W model - k-space Brown motion

③ Parisi - Wu $\Leftrightarrow \phi^3$ theory 图表示

$$h(k, t) = -[r + vk^2] h(k, t) + \eta + g \int \frac{q}{q} h(q, t) h(k-q, t)$$



④ $\langle h^2(x, t) \rangle \rightarrow t^\sigma$

今天讲 | Medina - Hwa - Kardar - Zhang 1989
 | 1986年 KPZ eq.

New: Correlated disorder

一般形式: $\langle \eta(x, t) \eta(x', t') \rangle = \int \int \delta(x-x') \delta(t-t')$
 \downarrow 可以推广为非白噪声
 $\int \int f_1(x-y) f_2(t-t')$
 or
 $[\int \delta(x-x')] \delta(t-t')$

k-space.

$$\langle \eta(k, \omega) \eta(k', \omega') \rangle = \int \int \delta(k+k') \delta(\omega+\omega') \frac{|k|^\sigma}{|k|}$$

可以改变临界值

eg: Schrödinger eq. dispersion relation $\frac{1}{vk^2+r} \sim \frac{1}{k^2}$

Dirac eq. $\frac{1}{Nvk^2+r} \sim \frac{1}{k}$

主要方程: $h(x, t) = v \nabla^2 h + \eta + \frac{g}{2} (\nabla h)^2$ 如果是 $\frac{g}{2} h^2 \Leftrightarrow \phi^4$ theory (Parisi - Wu).

动量空间:

$$h(x, t) = \sum_{k, \omega} e^{i(\vec{k} \cdot \vec{x} - \omega t)} h(k, \omega)$$



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关心 $e^{i(k \cdot \vec{x} - \omega t)}$ 系数.

$$-i\omega h(k, \omega) = -v^2 k^2 h(k, \omega) + \eta(k, \omega) + \frac{g}{2} \sum_{q, \Omega} (iq) [i(k-q)] h(q, \Omega) h(k-q, \omega-\Omega)$$

$$(vk^2 - i\omega) h(k, \omega) = \eta(k, \omega) - \frac{g}{2} \sum_{q, \Omega} q(k-q) h(q, \Omega) h(k-q, \omega-\Omega)$$

定义 $G_0(k, \omega) = \frac{1}{vk^2 - i\omega}$

$$h(k, \omega) = G_0(k, \omega) \eta(k, \omega) - \frac{g}{2} G_0(k, \omega) \sum_{q, \Omega} q(k-q) h(q, \Omega) h(k-q, \omega-\Omega)$$

若多次迭代, 表达式复杂

key point: 读懂图

1) IF $g=0$.

$$h(k, \omega) = G_0(k, \omega) \eta(k, \omega)$$

2) IF $g \neq 0$.

$$h(k, \omega) = G(k, \omega) \eta(k, \omega) \quad (\text{类似 Dyson eq.})$$

$$h(k, \omega) = \frac{k, \omega}{G} \otimes \eta(k, \omega) = \frac{k, \omega}{G_0} \times + \frac{k, \omega}{G_0} \begin{array}{l} \diagup \text{---} h(q, \Omega) \\ \diagdown \text{---} h(k-q, \omega-\Omega) \end{array}$$

$$\begin{array}{l} \diagup \\ \diagdown \end{array} = -\frac{g}{2} q(k-q)$$

进一步迭代:

$$h(k, \omega) = \frac{k, \omega}{G_0} \times + \frac{k, \omega}{G_0} \begin{array}{l} \diagup \times \\ \diagdown \times \end{array} + \frac{k, \omega}{G_0} \begin{array}{l} \diagup \times \\ \diagdown \times \\ \diagdown \times \end{array} + \frac{k, \omega}{G_0} \begin{array}{l} \diagup \times \\ \diagdown \times \\ \diagdown \times \\ \diagdown \times \end{array} + \dots$$

怎么做 RG?





不同的平均

原来 $\langle \phi \phi \rangle$

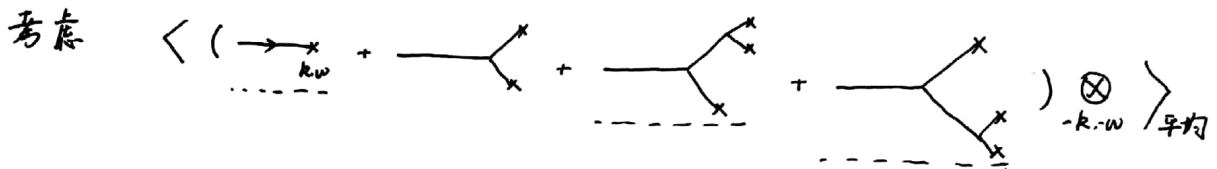
现在 $\langle \eta \eta \rangle$

$$G_0(k, \omega) = \frac{1}{v k^2} \quad (\text{真实值})$$

$$G(k, \omega) = \frac{1}{D k^2} \quad (\text{观测值})$$

$$h(k, \omega) = \text{---} \times = G(k, \omega) \eta(k, \omega)$$

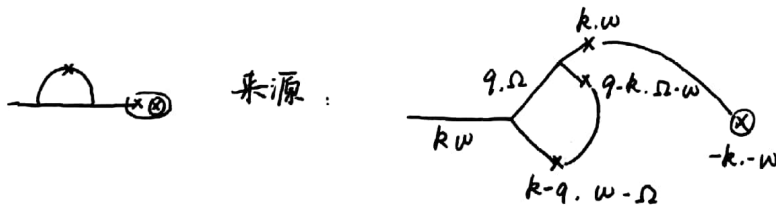
$$\begin{aligned} \langle h(k, \omega) \eta(-k, -\omega) \rangle &= G(k, \omega) \langle \eta(k, \omega) \eta(-k, -\omega) \rangle = G(k, \omega) D(k, \omega) \\ &= \text{---} \times \otimes \end{aligned}$$



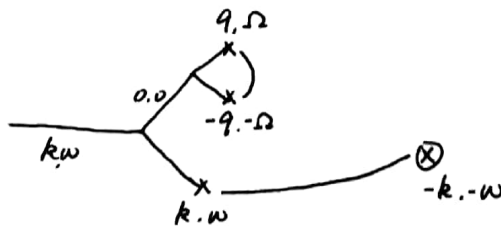
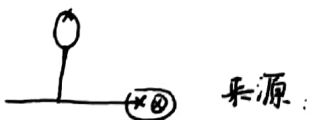
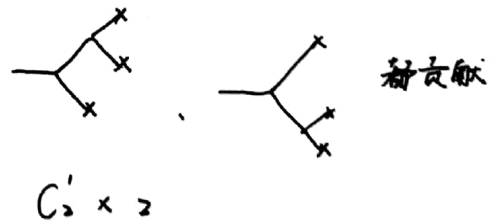
只有一三两项能配对有贡献

$$\langle h(k, \omega) \eta(-k, -\omega) \rangle = \text{---} \times \otimes + \text{---} \times \otimes + \text{---} \times \otimes$$

(可忽略在动量不为0)



拉直并考虑简并度



在动量为0费曼图不被允许时可不考虑费曼图

拉直并考虑简并度

2.



两边约去 $D(k, \omega)$

$$\frac{1}{G} = \frac{k, \omega}{G_0} + 4 \frac{\text{Diagram}}{G_0 \bar{\Sigma} G_0} \quad \bar{\Sigma} \leftarrow \text{自能}$$

规定动量方向

$\langle \eta(q, \Omega) \eta(-q, -\Omega) \rangle = 2D(q, \Omega)$

$+ \text{const.} \quad \text{Diagram} \quad + \dots$

$G_0 \bar{\Sigma} G_0 \bar{\Sigma} G_0$

$$G(k, \omega) = G_0(k, \omega) + 4 \left(-\frac{g}{2}\right)^2 \underline{G_0^2(k, \omega)} \bar{\Sigma} \quad \text{5个传播子}$$

$$\bar{\Sigma} = \sum_{q, \Omega} q(k-q) k q G_0(q, \Omega) G_0(k-q, \omega-\Omega) G_0(-q, -\Omega) 2D(q, \Omega)$$

讨论 g 很小

$$G = \frac{1}{\tilde{\nu} k^2}, \quad \text{set } \omega=0 \quad (t \rightarrow \infty)$$

$$G_0 = \frac{1}{\nu k^2} \quad \tilde{\nu} = \nu + \delta \nu$$

$$G = \frac{1}{(\nu + \delta \nu) k^2}$$

$$\frac{1}{(\nu + \delta \nu) k^2} = \left(\frac{1}{\nu k^2} \right) + 4 \left(-\frac{g}{2}\right) \left(\frac{1}{\nu k^2} \right) \bar{\Sigma}$$

$$\frac{1}{A+x} = \left(\frac{1}{A} \right) - \frac{x}{A^2} \quad A = \nu k^2 \quad x = -g^2 \bar{\Sigma}$$

$$\boxed{\tilde{\nu} k^2 = \nu k^2 - g^2 \bar{\Sigma}}$$

$$\tilde{\nu} k^2 = \nu k^2 - g^2 \sum_{q, \Omega} q(k-q) k q G_0(q, \Omega) G_0(-q, -\Omega) G_0(k-q, \omega-\Omega) \langle \eta \eta \rangle$$

$$= \nu k^2 - g^2 \left(\frac{1}{2\pi}\right)^d \int_{(-\infty, +\infty)} dq d\Omega \frac{(kq \cos \theta - q^2) k q \cos \theta 2D(q, \Omega)}{(\nu q^2 - i\Omega)(\nu q^2 + i\Omega)(\nu(k-q)^2 + i(\omega-\Omega))}$$

一般化: $\tilde{\nu} k^2 = \nu k^2 - A k^2 \int_0^\Lambda f(q) dq \quad q^{d-1} \quad f(q) \sim q^{-2p-3}$

回顾 ϕ^4 theory & $2D \delta$ potential.

$$\tilde{\nu}_R = \nu_\Lambda - A_\Lambda \int_0^\Lambda f_\Lambda(q) q^{d-1} dq$$



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$$\begin{aligned}
 v_R &= v_\Lambda - A_\Lambda \int_0^\Lambda f_\Lambda(q) q^{d-1} dq \\
 &= v_{\Lambda-d\Lambda} - A_{\Lambda-d\Lambda} \int_0^{\Lambda-d\Lambda} f_{\Lambda-d\Lambda}(q) q^{d'} dq \\
 &= v - \frac{\partial v}{\partial \Lambda} d\Lambda
 \end{aligned}$$

\tilde{v}_R : 重整化
 观测值
 与 Λ 无关

直接给出 RG-flow 结果: $\frac{dv}{d\ell} = v [z - z - K d' \frac{\lambda^2 D}{v^3} \frac{d-2+f(\Lambda)}{4d'}]$, $d' = d-1$

$= \alpha v$

$v \sim e^{\alpha \ell}$

讨论 $\langle h(k, \omega) h(-k, -\omega) \rangle = G(k, \omega) G(-k, -\omega) \langle \eta(k, \omega) \eta(-k, -\omega) \rangle$

$= 2D(k, \omega) G(k, \omega) G(-k, -\omega)$

$h =$

		0	
	0		0
		0	$\propto g^4$ (不考虑)

= + +

顶点:

= + + + ...



$$\tilde{g} = g_\Lambda + g_\Lambda^3 \int^\wedge dq dq' d\Omega d\Omega' (\dots)$$

不变

Summary: KPZ eq

$$\left| \begin{array}{l} \frac{dV}{dt} \\ \frac{dD}{dt} \dots \langle \eta \eta \rangle \\ \frac{dg}{dt} \end{array} \right.$$

⇒ 随机微分方程

场论

E-W model (Kardar 书) ⇒ k-space Brown motion

图表示 Parisi 工作 RSB

spin glass 无序 + 多体 $Z = \int D_j e^{-S}$

Replica trick 长程相互作用 排斥

Fermion: SC Fermi Liquid.

Gross - Neveu model.

GN - Yukawa model

Kondo effect.

SYK model

$T_{ijkl} \chi_i \chi_j \chi_k \chi_l$

