

无序 \rightarrow 平均
 \uparrow
 $\langle \dots \rangle \rightarrow$ 平均

相似点 $\langle x \rangle = \int D\phi x e^{-S} / \int D\phi e^{-S} Z$
 $= \int D\phi x P(\phi)$, $P(\phi) = \frac{e^{-S}}{Z}$.

热力学平均 \Leftrightarrow 平衡态方式
 平衡 $\Leftrightarrow D = k_B T \mu$.

无序 $\langle \eta \eta \rangle \sim \delta(x-x') \delta(t-t')$

Parisi 工作: (不用固定规范的做法微扰)

$\mathcal{L} = \frac{1}{2}(\dot{\phi})^2 - \frac{g}{2!}\phi^3 - \frac{m}{2}\phi^2$, $\partial \mathcal{L} / \partial(\partial \mu \phi) = \partial \mathcal{L} / \partial \phi$

$\ddot{\phi} = \partial^2 \phi - m^2 \phi - \frac{g}{2} \phi^2 - \frac{1}{\mu} \dot{\phi} + \xi$. overdamped regime $\dot{\phi} \sim 0$.

$\dot{\phi} = \partial^2 \phi - m^2 \phi - g \phi^2 + \eta$

$g=0$, $\frac{\partial \phi(k,t)}{\partial t} = -\mu_k \phi(k,t) + \eta(k,t)$, $\mu_k = k^2 + m^2$

方差, $t \rightarrow \infty$
 $\bar{\phi} = 0$.

$\phi(k,t) = e^{-\mu_k t} \phi(k,0) + \int_0^t e^{-\mu_k(t-t')} \eta(k,t') dt'$

确定平均 $\phi(k,t)$ 涨落随机
 对方差无影响

t 足够大时, $\phi(k,t) = \int_0^t e^{-\mu_k(t-t')} \eta(k,t') dt'$

$\langle \phi(-k,t) \phi(k,t) \rangle = \int_0^t e^{-\mu_k(t-t_1)} e^{-\mu_k(t-t_2)} \langle \eta(k,t_1) \eta(-k,t_2) \rangle$

$= 2D \int_0^t e^{-2\mu_k(t-t')} dt'$

$t \rightarrow \infty$ 时 $\frac{\partial D}{\partial \mu_k} = \frac{D}{k^2 + m^2}$ 取 $D=1 \Leftrightarrow \langle \phi(-k,t) \phi(k,t) \rangle = \frac{1}{k^2 + m^2}$

非平衡系统 (dynamical)

平均下来 propagator 一样. 都是 $G_0 = \frac{1}{k^2 + m^2}$.

经典场论中 $\int D\phi e^{-\beta H} \sim \int D\phi e^{-\beta(k^2 + m^2)} \phi_k^* \phi_k$

$\langle \phi_k^* \phi_k \rangle = \frac{1}{\beta(k^2 + m^2)} = \frac{k_B T}{k^2 + m^2}$

$\dot{\phi} = \mu(\partial^2 \phi - m^2) + \eta \Leftrightarrow \mu_k = \mu(k^2 + m^2)$





Einstein relation

$$\langle \phi(-k, t) \phi(k, t) \rangle = \frac{D}{\mu(k^2 + m^2)} = \frac{k_B T \mu}{\mu(k^2 + m^2)} \quad (D = k_B T \mu)$$

$$\langle \chi \rangle = \langle \phi_k^* \phi_k \rangle = \frac{k_B T}{k^2 + m^2} \rightarrow \text{相变 (critical phys)}$$

$$\langle \phi(k, t) \phi(k, t) \rangle = \frac{k_B T}{k^2 + m^2} \quad \text{相变 (dynamical critical phys)}$$

KP2: 1986 $g\phi^2 \rightarrow g(\nabla\phi)^2$

$$\mu=1, D=1 \quad \left| \quad \dot{\phi} = \partial^2 \phi - m^2 \phi - g\phi^2 + \eta, \quad \dot{\phi}(k, t) = -(k^2 + m^2)\phi(k, t) + g \int \phi(q, t) \phi(k-q, t) \dagger \eta(k, t) \right.$$

$t \rightarrow \infty$, 对应是 frequency 区间什么物理 ($\omega \rightarrow 0$)

$$f(t) = \sum_{\omega} f_{\omega} e^{i\omega t}, \quad \frac{1}{t} \int_0^t f(t) dt = f_0 \quad (\text{低频有效}) \quad \text{低能物理}$$

接下来重点: 1) 认识 Feynman 图, 了解规则.

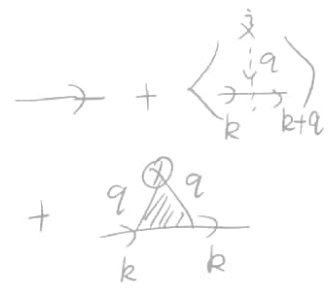
2) 如何做平均 \rightarrow *

$$\begin{cases} H = H_0 + V(\vec{x}) \\ V(\vec{x}) = \sum_i v_i \delta(\vec{x} - \vec{x}_i) \end{cases} \quad \begin{cases} \langle v_i \rangle = 0 \\ \langle v_i v_j \rangle = 2D \delta_{ij} \end{cases}$$

phys. transport.

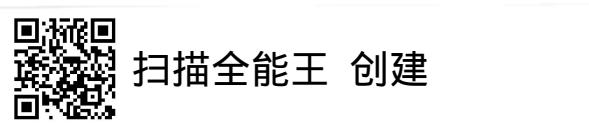
$$S = S_{\psi} + S_{\phi} + S_g + S_{\lambda} + S_I$$

\uparrow 2D Dirac \uparrow scalar field \uparrow $\frac{1}{4} \phi^4$
 \uparrow $\frac{1}{2}(\partial\phi)^2 - \frac{m}{2}\phi^2$ \uparrow $\int v_i(x) \psi^\dagger \psi dx$



Replica trick \Leftrightarrow spin Glass

- 无序 \Rightarrow 平均
- 1) Replica trick, $n \rightarrow 0$
 - 2) Brown motion, $\langle \eta \eta \rangle = 2D \delta(x-x') \delta(t-t')$
 - 3) Transport $\langle G \rangle = \frac{1}{N} \sum_i G(z_i)$
- * $\langle \eta \eta \eta \eta \rangle \Rightarrow$ All pairings.



$$\int_{(x,t)}^{G_0} \int_{(x',t')} \eta \Leftrightarrow G\eta \Rightarrow G * \eta \text{ 卷积}$$

$$\dot{\phi} = \partial\phi - m^2\phi^2 + g\phi^3 + \eta$$

$$\phi(x,t) = \int G(x-x', t-t') [\eta(x',t') + g\phi^2(x',t')] dx' dt'$$

$$= \int G(x-x', t-t') \eta(x',t') dx' dt' + g \int G(x-x', t-t') dx' dt' \int G(x'-x_2, t'-t_2) [\eta + g\phi^2]$$

$$G(x'-x_2, t'-t_2) [\eta + g\phi^2] dx_2 dt_2 dx_1 dt_1 + \dots$$

$$\phi(x,t) = \text{---}x + \frac{1}{(x,t)} \int \eta(x',t') \text{---}x$$

$$\phi = \text{---}x + \text{---}x + \text{---}x + \dots \Leftrightarrow \begin{cases} \langle \phi(x,t) \rangle \\ \langle \phi(x,t_x) \phi(y,t_y) \rangle \end{cases}$$

$$\langle \phi^2(x,t) \rangle = \int \frac{1}{k} \mu_k = \frac{1}{k} \frac{1}{k^2 m^2}$$

$$\langle \phi(x,t_x) \phi(y,t_y) \rangle = \frac{(x',t')}{(x,t_x) (y,t_y)} = \frac{1}{k^2 m^2} e^{ik(x-y)} \text{ 相差一个傅利叶变换}$$

$$\langle \phi(x) \phi(y) \rangle \sim \frac{1}{k} \frac{1}{k^2 m^2} e^{ik(x-y)}$$

	$\text{---}x$	$\text{---}x$	$\text{---}x$	$\text{---}x$
$\text{---}x$	$\text{---}x$	0	0	$\text{---}x$ (等边)
$\text{---}x$	0	2 $\text{---}x$	0	
		0		

$$G(x-x', t-t') = \frac{1}{i\omega - (k^2 + m^2)} e^{i(kx - \omega t)}$$

$$\frac{D(k)}{i(k^2 + m^2)} \propto \frac{k}{k^2 + m^2} \sim \frac{1}{k}$$

$D(k)$ 重要

$\partial^4 \phi$ 为什么不讨论? 这项对低能有效模型不起作用 \therefore 其 $\propto k^4 \phi$

Navier - Stoke eq.

