

(5.19) E-W model } 动量空间的 Brown motion

$$\frac{\partial h}{\partial t} = \nu \nabla^2 h + \eta$$

$$\frac{\partial h(k,t)}{\partial t} = -\nu k^2 h(k,t) + \eta(k,t)$$

$$\langle \eta(k,t) \eta(k',t') \rangle = 2D(k) \delta(t-t') \delta(k+k')$$

(与 pairing 相关)

$$\star \frac{\partial f}{\partial t} = -\mu f + \eta \Rightarrow f(t) = f(0) e^{-\mu t} + \int_0^t e^{-\mu(t-t')} \eta(t') dt'$$

$$\sigma^2 = \langle (h(x,t) - \bar{h})^2 \rangle \propto t^\beta \quad \star \text{Kardar Chap. 9.}$$

1990s - 21世纪初 表面科学、多重分形结构

KPZ 方程

1) KPZ 原始论文 1986

2) Medina, Hwa, Kardar, Zhang, 1989, PRA

3) Parisi, 吴咏时 "不同固定规范下的微扰论" 1980 中国科学

$$\dot{\phi} = \partial^2 \phi - m^2 \phi + g \phi^3 + \eta$$

1979-1983 Parisi \Rightarrow spin glass

4) Toner, Yuhai, Tu PRE 1998

(USTC 少年班, 1983-1987)

$$\partial_t \vec{v} + \lambda_1 (\vec{v} \cdot \nabla) \vec{v} + \lambda_2 (\nabla \cdot \vec{v}) \vec{v} + \lambda_3 \nabla \vec{v}^2 = \alpha v - \beta |\vec{v}|^2 \vec{v} - \nabla p + \dots + \vec{\eta}$$

2020 Onsager prize

C.N. Yang

Jason Ho (冷原子)

Yuhai Tu 生物物理

$$\text{矢量 } \vec{v} = \begin{pmatrix} v_1 \\ v_2 \\ v_3 \end{pmatrix}$$

非线性

关键

$$Z = \int D\phi e^{-S}$$

$$\langle \phi(x) \phi(x') \rangle = \frac{\int D\phi \phi \phi e^{-S}}{\int D\phi e^{-S}}$$

↑ 平均: 场在 S 下平均, 本质上给一个分布

随机数

1) 构造 S

2) 不用 S, 直接求平均

$$\langle \eta \eta \rangle = \dots$$

KPZ eq 1986

首先概念:
$$\frac{\partial h}{\partial t} = \nu \nabla^2 h + \frac{\lambda}{2} (\nabla h)^2 + \eta$$

$$\langle \eta(\vec{x}, t) \eta(\vec{x}', t') \rangle = 2D \delta(\vec{x} - \vec{x}') \delta(t - t')$$



扫描全能王 创建

性质 1: 伽利略不变性

$$\begin{cases} t' = t \\ \vec{x}' = \vec{x} + \lambda \vec{\epsilon} t \\ h' = h + \vec{\epsilon} \cdot \vec{x} \end{cases}$$

性质 2: Burgers eq at $\eta = 0$

可解

对于 $\frac{\partial h}{\partial t} = \nu \partial^2 h + \frac{\lambda}{2} (\partial h)^2$

可定义 $W = e^{\frac{\lambda}{2\nu} h}$

$$\frac{\partial W}{\partial t} = \frac{\lambda}{2\nu} \frac{\partial h}{\partial t} W$$

$$\frac{\partial W}{\partial x} = \frac{\lambda}{2\nu} \frac{\partial h}{\partial x} W$$

$$\frac{\partial^2 W}{\partial x^2} = \left(\frac{\lambda}{2\nu}\right)^2 \left(\frac{\partial h}{\partial x}\right)^2 W + \frac{\lambda}{2\nu} \frac{\partial^2 h}{\partial x^2} W$$

$$\frac{\partial W}{\partial t} = \text{const.} \cdot \frac{\partial^2 W}{\partial x^2}$$

反代换 $h = \frac{2\nu}{\lambda} \ln W$

$$V(\vec{x}) = \eta(\vec{x}, t)$$

同理 $\frac{\partial h}{\partial t} = \nu \partial^2 h + \frac{\lambda}{2} (\partial h)^2 + V(x) \Rightarrow \frac{\partial W}{\partial t} = \nu \partial^2 W + \boxed{V(x)} W$

随机势中扩散问题

性质 3: 定义 $\vec{v} = -\nabla h$

$$\frac{\partial \vec{v}}{\partial t} + \lambda \vec{v} \cdot \nabla \vec{v} = \nu \nabla^2 \vec{v} - \nabla \eta$$

Navier - Stokes eq.

Remark: $\nabla \eta$ 是随机数

Colored noise
Correlated noise

$$\begin{cases} \langle \eta(\vec{x}, t) \nabla \eta(\vec{y}, t) \rangle \sim \delta(x-y) \delta(t-t') (\nabla)_i (\nabla)_j \\ \langle \eta(\vec{k}, \omega) \eta(\vec{k}', \omega') \rangle \sim f(k, k') \delta(k+k') \delta(\omega+\omega') \end{cases}$$

求解

$$\frac{\partial h}{\partial t} = \nu \partial^2 h + \frac{\lambda}{2} (\partial h)^2 + \eta$$

无解析解, 只能用 perturbation, $\lambda = 0$

平均 $\langle \eta \eta \rangle$ 等效 $\int D\phi \phi \phi e^{-S} / \int D\phi e^{-S}$

背景

钱伟长 \Rightarrow 边界层理论

多重尺度分析 multiple scaling method



$\ddot{y} + y = \varepsilon (y - \frac{1}{3} y^3)$ Rayleigh problem

$$y = y_0 + \varepsilon y_1 + \varepsilon^2 y_2 + \dots$$

$$\ddot{y}_0 + y_0 + \varepsilon (\ddot{y}_1 + y_1) + \varepsilon^2 (\ddot{y}_2 + y_2) = \varepsilon [y_0 + \varepsilon y_1 + \varepsilon^2 y_2 - \frac{1}{3} (y_0 + \varepsilon y_1 + \varepsilon^2 y_2)^3]$$

$$\varepsilon^0: \ddot{y}_0 + y_0 = 0 \Rightarrow y_0 = A \cos(t + \theta)$$

$$\varepsilon^1: \ddot{y}_1 + y_1 = y_0 - \frac{1}{3} y_0^3 \Rightarrow \ddot{y}_1 + y_1 = A \sin(t + \theta) + \dots$$

$$\varepsilon^2: \ddot{y}_2 + y_2 = y_1 - y_0^2 y_1$$

如果对 y_1 做 Fourier transformation
每一项抵消

$$\ddot{y}_1 + y_1 = A \cos[(1+\eta)t + \theta]$$

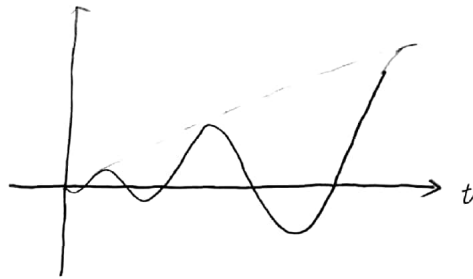
$$y_1 = K \cos[(1+\eta)t + \theta]$$

$$y_1 = \frac{A}{(1-(1+\eta)^2)} \cos[(1+\eta)t + \theta]$$

$$\sim \cos(t + \theta) + t \sin(t + \theta)$$

可能与奇异微扰理论相关

Singular perturbation theory



(简要理解
受迫振动形式)

$$\frac{\partial h}{\partial t} = \nu \nabla^2 h + \frac{\hbar}{2} (\nabla h)^2 + \eta$$

$$h(x, t) = \sum_{\vec{k}} e^{i\vec{k}x} h(\vec{k}, t)$$

$$\eta(x, t) = \sum_{\vec{k}} e^{i\vec{k}x} \eta(\vec{k}, t) \quad \text{讨论 } \eta(\vec{k}, t) \text{ 性质}$$

$$\langle \phi_{\vec{k}} \phi_{-\vec{k}} \rangle \neq 0 \quad \text{or} \quad \langle \phi_{\vec{k}} \phi_{\vec{q}} \rangle = \frac{\delta_{\vec{k}+\vec{q}}}{k^2 + m^2}$$

$$\sum_{\vec{k}} \left(\frac{\partial}{\partial t} h(\vec{k}, t) \right) e^{i\vec{k}x} = - \sum_{\vec{k}} \nu k^2 e^{i\vec{k}x} h(\vec{k}, t) + \eta(\vec{k}, t) e^{i\vec{k}x} + \sum_{\vec{k}} \sum_{\vec{q}} \frac{\hbar}{2} (iq)(i(k-q)) e^{i\vec{k}x} h(\vec{q}, t) h(\vec{k}-\vec{q}, t)$$

取 $e^{i\vec{k}x}$:

$$\frac{\partial h(\vec{k}, t)}{\partial t} = -\nu k^2 h(\vec{k}, t) + \eta(\vec{k}, t) - \sum_{\vec{q}} \frac{\hbar}{2} q(k-q) h(\vec{q}, t) h(\vec{k}-\vec{q}, t)$$



$$\langle \eta(k, t) \eta(k', t') \rangle$$

$$\left(\begin{aligned} \eta(x, t) &= \sum_k e^{ikx} \eta(k, t) \\ \eta(k, t) &= \int e^{-ikx} \eta(x, t) dx \end{aligned} \right.$$

$$= \int e^{-ikx - ik'y} \frac{\langle \eta(x, t) \eta(y, t') \rangle}{2D \delta(x-y) \delta(t-t')} dx dy$$

$$= \int e^{-i(k+k')x} 2D dx \delta(t-t')$$

$$= (2\pi)^D 2D \delta(k+k') \delta(t-t') \quad \leftarrow \text{明确物理意义}$$

解: $G_0 = e^{-\nu k^2 t} \Theta(t)$ step function

$$h(k, t) = \frac{G_0(k, t) h(k, 0) + \int_0^t G_0(k, t-\tau) \eta(k, \tau) d\tau}{-\frac{\lambda}{2} \sum_q q(k-q) \int_0^t G_0(k, t-\tau) h(q, \tau) h(k-q, \tau) d\tau}$$

$\lambda=0$, E-W mode 解

迭代方程 if $\lambda=0$.

Parisi - Wu model.

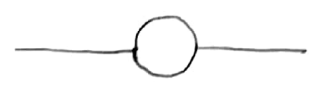
$$\mathcal{F} \left(\begin{aligned} \phi &= \partial^2 \phi - m^2 \phi + g \phi^2 + \eta & \langle \eta(x, t) \eta(x', t') \rangle &= 2\delta(x-x') \delta(t-t') \\ \phi(k, t) &= -k^2 \phi(k, t) - m^2 \phi(k, t) + g \sum_q \overset{? \text{ const}}{\phi(q, t) \phi(k-q, t)} + \eta(k, t) \end{aligned} \right.$$

① $g=0$

$$\phi(x, t) = \int G(x-y, t-\tau) \eta(y, \tau) dy d\tau \Rightarrow \langle \phi(x, t) \phi(x', t') \rangle_{\text{平均}} \left\langle \begin{matrix} S \\ \text{disorder} \end{matrix} \right.$$

$$\langle \phi(x, t) \phi(x', t') \rangle = \int G(x-y, t-\tau) G(x'-y', t'-\tau')$$

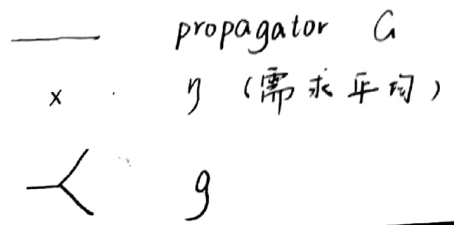
$$\langle \eta(y, \tau) \eta(y', \tau') \rangle dy dy' d\tau d\tau' = 2 \int G(x-y, t-\tau)^2 dy d\tau$$



② $g \neq 0$

解用 propagator 表示

$$\phi(x, t) = \int G(x-y, t-\tau) [\eta(y, \tau) + g \phi^2(y, \tau)] dy d\tau$$



$$\phi = \frac{x}{G\eta} + \text{diagram} + \text{diagram}$$

$$\int G(x-y, t-\tau) \eta(y, \tau) d\tau dy$$

$$= G\eta$$

$$\langle \phi^2 \rangle = \text{diagram} \propto g^0 D \quad D = \text{平均}$$

$$+ \text{diagram} \propto g^1 D^2$$

$$+ \text{diagram} \propto g^1 D^2$$

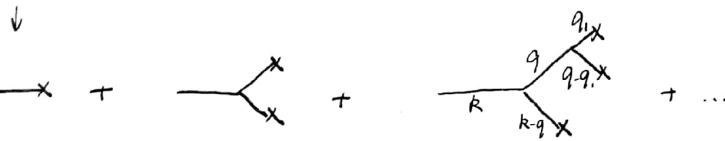
$$+ \text{diagram} \propto g^1 D^2$$

+ ...

KPZ eq $h = G_0 \eta - \frac{\lambda}{2} q(k-q) G_0 h h$

迭代

$$= G_0 \eta - \frac{\lambda}{2} q(k-q) [G_0 \eta - \frac{\lambda}{2} q(k-q) G_0 h h] [G_0 \eta - \frac{\lambda}{2} q(k-q) G_0 h h]$$



如何平均 1) $\int \langle h^2(x, t) \rangle dx = \sigma^2 \propto t^2$

$$\int_0^{\Lambda} \Lambda e^{-dl} dq + \int_{\Lambda}^{\Lambda} \Lambda e^{-dl} dq$$

Shang-keng Ma

马上度

