

$$t \rightarrow \infty = 2D \left(\frac{1}{2\pi}\right)^d \int_0^\infty \frac{dk k^{d-1}}{v k^2} (1 - e^{-2vk^2 t})$$

$$\downarrow \text{令 } vk^2 t = y^2 \Rightarrow y = \sqrt{vt} k$$

$$= 2D \left(\frac{1}{2\pi}\right)^d \frac{1}{v} (v\sqrt{t})^{d-2} \int_0^{+\infty} \frac{y^{d-1}}{y^2} (1 - e^{-y^2}) dy$$

$$\propto t^{\frac{d-2}{2}} \int_0^{+\infty} \frac{y^{d-1}}{y^2} (1 - e^{-y^2}) dy$$

(5.19) E-W model } 动量空间 Brown motion.

$$\frac{\partial h}{\partial t} = v \nabla^2 h + \eta$$

$$\frac{\partial h(\mathbf{k}, t)}{\partial t} = -v k^2 h(\mathbf{k}, t) + \eta(\mathbf{k}, t)$$

$$\& \frac{\partial f}{\partial t} = -\mu f + \eta \Rightarrow f(t) = f(0) e^{-\mu t} + \int_0^t e^{-\mu(t-t')} \eta(t') dt' \quad (\text{勘误})$$

$$\sigma^2 = \langle (h(\vec{x}, t) - \bar{h}(\vec{x}, t))^2 \rangle = \langle \eta(\mathbf{k}, t) \eta(-\mathbf{k}, t) \rangle = 2D(\mathbf{k}) \delta(t-t')$$

* Kardar's book CH 9. $\propto t^{\frac{d-2}{2}}$

1990s - 21世纪初: 表面科学, 多重分形结构.

↑ 不同 mode 扩散系数不同

KPZ 方程 (Kardar 书元)

Ref: ① KPZ 1986 原文

② Medina, Hwa, Kardar, Zhang, 1989, PRA

③ Parisi, Y.S. Wu: 不用固定规范的微扰论' 1980 << 中国科学 >>

$$\left. \begin{array}{l} \text{Ref: ① KPZ 1986 原文} \\ \text{② Medina, Hwa, Kardar, Zhang, 1989, PRA} \\ \text{③ Parisi, Y.S. Wu: 不用固定规范的微扰论' 1980 << 中国科学 >>} \end{array} \right\} \frac{\partial \phi}{\partial t} = v \partial^2 \phi + \frac{\lambda}{2} (\partial \phi)^2 + \eta$$

$$\dot{\phi} = v \partial^2 \phi - m^2 \phi + g \phi^2 + \eta$$

1979-1983 Parisi Spin Glass

④ Toner, Yuhai Tu. PRE, 1998 (USTC 少: 83-87)

$$\partial_t \vec{v} + \lambda_1 (\vec{v} \cdot \nabla) \vec{v} + \lambda_2 (\nabla \cdot \vec{v}) \vec{v} + \lambda_3 \nabla \vec{v}^2 = v \vec{v} - \beta |\vec{v}|^2 \vec{v} - \nabla p + \dots + \vec{\eta}$$

2020 Onsager Prize.

C.N. Yang
Jason Ho

Yuhai Tu New 生物物理

矢量, 非线性 eq.

$$\text{key point: } z = \int D\phi e^{-S}$$

$$\langle \phi(x) \phi(0) \rangle = \frac{\int D\phi \phi \phi e^{-S}}{\int D\phi e^{-S}} \quad \text{平均的意义: 场在 } S \text{ 下的平均}$$

随机数: 1) 构造 S, 仿照上面 $\xi(x) \phi(x)$

2) 不同 S 求平均, 随机数的平均. $\langle \eta \eta \rangle = ?$

KPZ eq 1986

首先: 概念.

$$\text{Model: } \frac{\partial h}{\partial t} = v \nabla^2 h + \frac{\lambda}{2} (\nabla h)^2 + \eta$$

$$\langle \eta(\vec{x}, t) \eta(\vec{x}', t') \rangle = 2D \delta(\vec{x} - \vec{x}') \delta(t - t')$$

性质: ① 伽利略不变性

$$\begin{cases} t' = t \\ \vec{x}' = \vec{x} + \lambda \vec{e} t \\ \Gamma = \Gamma + \vec{e} \cdot \vec{x} \end{cases}$$

$$\frac{\partial W}{\partial t} = \frac{\partial}{\partial t} \frac{\partial h}{\partial t} W$$

$$\frac{\partial W}{\partial \vec{x}} = \frac{\partial}{\partial \vec{x}} \frac{\partial h}{\partial \vec{x}} W$$

② Burgers eq. at $\eta = 0$.

↳ solvable. 随机场中扩散

$$\frac{\partial W}{\partial \vec{x}} = \left[\left(\frac{\partial}{\partial \vec{x}} \right) \left(\frac{\partial h}{\partial \vec{x}} \right) + \frac{\partial}{\partial t} \left(\frac{\partial h}{\partial \vec{x}} \right) \right] W$$

$$\frac{\partial h}{\partial t} = v \partial^2 h + \frac{\lambda}{2} (\partial h)^2$$

$$\text{define } W = e^{\frac{\lambda}{2v} h}$$

$$\frac{\partial W}{\partial t} = v \partial^2 W, \quad h = \frac{2v}{\lambda} \ln W$$

$$\frac{\partial W}{\partial t} = \nu \partial^2 W + V(\vec{x}) W$$

\uparrow
 $V(\vec{x}) = \eta(\vec{x}, t)$

Remarks: $\nabla \eta$ 是随机数.

$$\langle \nabla \eta(\vec{x}, t) \nabla \eta(\vec{y}, t) \rangle \sim \delta(\vec{x} - \vec{y}) \delta(t - t') \nabla \nabla$$

$$\langle \eta(\vec{R}, \omega) \eta(\vec{R}', \omega) \rangle \sim f(\vec{R}, \vec{R}') \delta(\vec{R} + \vec{R}') f(\omega, \omega')$$

③ Define $\vec{v} = -\nabla h$.

$$\frac{\partial \vec{v}}{\partial t} + \lambda \vec{v} \cdot \nabla \vec{v} = \nu \nabla^2 \vec{v} - \nabla \eta$$

(Navier-Stokes eq.)

求解: $\frac{\partial h}{\partial t} = \nu \partial^2 h + \frac{\lambda}{2} (\partial h)^2 + \eta$

无解析解, 只能微扰做. $\lambda \rightarrow 0$ ✓

平均 $\langle \eta \eta \rangle \rightarrow \int (\partial \phi) \phi \phi e^{-S} / \int \partial \phi e^{-S}$

$\epsilon^0: \dot{y}_0 + y_0 = 0$

eg: $y_0 = A \cos(\omega t + \theta)$

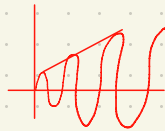
$\epsilon^1: \dot{y}_1 + y_1 = \dot{y}_0 - \frac{\lambda}{2} y_0^2$

$\Rightarrow \dot{y}_1 + y_1 = A \cos(\omega t + \theta) + \dots$

$\epsilon^2: \dot{y}_2 + y_2 = \dot{y}_1 - \lambda y_0^2 y_1$

Fourier trans \times
 $\dot{y}_1 + y_1 = A \cos(\omega + i\eta) t + \theta$
 $y_1 = k \cos(\omega + i\eta) t + \theta$

$\Rightarrow [1 - k(\omega + i\eta)] = A \Rightarrow y_1 = \frac{A}{(1 - i\eta)^2} \cos((\omega + i\eta)t + \theta)$
 $\sim \cos(\omega t + \theta) + t \sin(\omega t)$



背景, 钱伟长: 边界层理论 ✓
 多重尺度分析 ✓
 multiple scaling method

作业: $\ddot{y} + y = \epsilon (\dot{y} - \frac{1}{2} y^3)$ Rayleigh problem
 $y = y_0 + \epsilon y_1 + \epsilon^2 y_2 + \dots$

$\Rightarrow \dot{y}_0 + y_0 + \epsilon (\dot{y}_1 + y_1) + \epsilon^2 (\dot{y}_2 + y_2)$

$= \epsilon [y_0 + \epsilon y_1 + \epsilon^2 y_2 - \frac{1}{2} (y_0 + \epsilon y_1 + \epsilon^2 y_2)^3]$

《奇异微扰理论》✓
 \uparrow
 Singular

$$\frac{\partial h}{\partial t} = \nu \partial^2 h + \frac{\lambda}{2} (\partial h)^2 + \eta$$

$$\frac{\partial}{\partial t} (e^{ikx} h(k, t)) e^{ikx} = -\frac{\nu}{k} \nu k^2 e^{ikx} h(k, t) + \eta(k, t) e^{ikx} + \frac{\nu}{k} \frac{\partial}{\partial t} (\frac{1}{2} (i q) (i(k-q)) e^{ikx} h(q, t) h(k-q, t))$$

$$h(x, t) = \frac{1}{k} e^{ikx} h(k, t)$$

$$\eta(x, t) = \frac{1}{k} e^{ikx} \eta(k, t)$$

比较 e^{ikx} 系数:

讨论 $\eta(k, t)$ 性质

$\Rightarrow \frac{\partial}{\partial t} h(k, t) = -\nu k^2 h(k, t) + \eta(k, t) - \frac{\nu}{2} \frac{\partial}{\partial t} \eta(k-q) h(q, t) h(k-q, t)$

$\langle \phi_k \phi_{-k} \rangle \neq 0$ or $\langle \phi_k \phi_q \rangle = \left(\frac{\delta_{k+q}}{k^2 + q^2} \right)$

Solution: $G_0 = e^{-\nu k^2 t} \Theta(t)$

$\langle \eta(k, t) \eta(k', t) \rangle$

$\eta(x, t) = \frac{1}{k} e^{ikx} \eta(k, t)$

$\eta(k, t) = \int e^{-ikx} \eta(x, t) dx$

$= \int e^{-ikx - iky} \langle \eta(x, t) \eta(y, t) \rangle dx dy$

$= \int e^{-i(k+k')x} \int \delta(x-t) \delta(y-t) dx dy$

$= (2\pi)^d \int \delta(k+k') \delta(t-t') \leftarrow$ 有明确物理意义. 迭代方程: IF $\lambda = 0$.

E-W model 的解 ($\lambda=0$)

$h(k, t) = G_0(k, t) h(k, 0) + \int_0^t G_0(k, t-\tau) \eta(k, \tau) d\tau$

$- \frac{\nu}{2} \frac{\partial}{\partial t} \eta(k-q) \int_0^t G_0(k, t-\tau) h(q, \tau) d\tau$

$h(k-q, \tau) d\tau$

Parisii - Wu model:

$$\dot{\phi} = \partial^2 \phi - m^2 \phi + g \phi^2 + \eta$$

$\langle \eta(x, t) \eta(x', t') \rangle = 2\delta(x-x') \delta(t-t')$

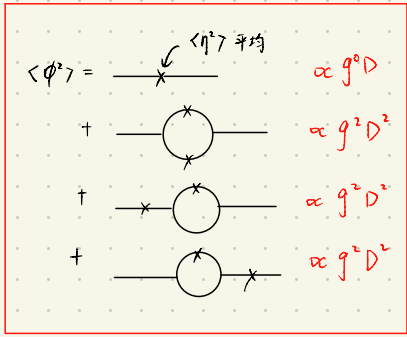
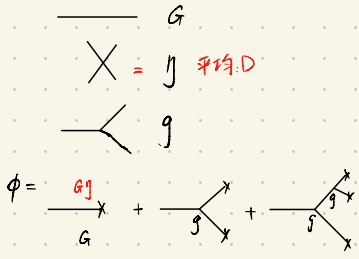
Fourier trans: $\dot{\phi}(k,t) = -k^2 \phi(k,t) - m^2 \phi(k,t) + g \frac{\delta}{\delta \phi} \phi(q,t) \phi(k-q,t) + \eta(k,t)$

讨论: $\phi(x,t) = \int G(x-y, t-\tau) [\eta(y,\tau) + g \phi^2(y,\tau)] dy d\tau$

$g=0$ \downarrow $= \int G(x-y, t-\tau) \eta(y,\tau) dy d\tau \Rightarrow \langle \phi(x,t) \phi(x',t') \rangle_{g=0} \begin{cases} S \\ \text{Disorder} \end{cases}$

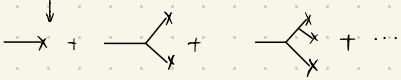
$\langle \phi(x,t) \phi(x',t') \rangle = \int G(x-y, t-\tau) G(x-y', t-\tau') \langle \eta(y,\tau) \eta(y',\tau') \rangle dy dy' d\tau d\tau'$
 $= 2 \int G(x-y, t-\tau)^2 dy d\tau$ ~~—————~~

$g \neq 0$, 上面的解用 propagator 表示.



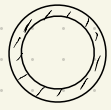
回到 KPZ: $h = \text{horizontal line with cross} + \text{vertex with two outgoing lines} + \dots$

$h = G\eta - \frac{1}{2} g(k-q) G h h$
 $= G\eta - \frac{1}{2} g(k-q) G_0 [G_0 \eta - \frac{1}{2} g(k-q) G_0 h h] [G_0 \eta - \frac{1}{2} g(k-q) G_0 h h]$



如何平均: 1) $\int \langle h^2(x,t) \rangle dx = \sigma^2 \propto t^z$

$\int_0^\wedge e^{-dt} dq + \int_\wedge^\wedge e^{-dt} dq$



Shang-Keng Ma 马上海