

## Casimir effect.

$$1D: \quad E = \hbar c k \quad , \quad k_n = \frac{n\pi}{a}$$

$$E = \frac{\hbar c}{2} \sum_n k_n = \frac{\hbar c}{2} \sum_{n=1}^{\infty} \frac{n\pi}{a}$$

ref. Zeta function. / 正则化.

$$\zeta(s) = \sum_{n=1}^{\infty} \frac{1}{n^s}$$

$$\zeta(-1) = -\frac{1}{12}$$

$$\Rightarrow E = \frac{\hbar c}{2} \cdot \left(-\frac{\pi}{12a}\right) = -\frac{\hbar c \pi}{24a}$$


驻波.

$$\begin{aligned} E^f &= \frac{1}{2} \hbar c \sum_k |k| = \frac{1}{2} \hbar c \frac{L}{2\pi} \int_0^{\infty} dk |k| \\ &= \frac{\hbar c}{2} \frac{L}{\pi} \int_0^{+\infty} dk \cdot k \end{aligned}$$

$$\begin{aligned} \tilde{E} &= \frac{E^f}{L} \cdot a = \frac{\hbar c a}{2\pi \epsilon^2} \int_0^{+\infty} \epsilon k \, d\epsilon \, e^{-\epsilon k} \quad \text{|||||} \\ &= \frac{\hbar c a}{2\pi \epsilon^2} \quad \leftarrow \text{没有板时真正的能量.} \end{aligned}$$

$$\text{又} \quad E = -\frac{\hbar c}{2} \sum_{n=1}^{+\infty} \frac{n\pi}{a}$$

$$= -\frac{\hbar c}{2} \sum_{n=1}^{+\infty} k_n \left| \begin{array}{l} e^{-k_n \epsilon} \\ f(k_n \epsilon) \\ e^{-k_n / \lambda} \end{array} \right. \quad (\epsilon = 1/\lambda)$$

$$= \frac{\hbar c a}{2\pi \epsilon^2} - \frac{\hbar c \pi}{24a}$$

有板存在时增加的能量.

$$\Rightarrow \text{Casimir Force} = -\frac{\partial E}{\partial a} = -\frac{\hbar c \pi}{24a^2}$$

耦合常数与选择的能标有关系.

$$E = -\frac{\hbar^2 \lambda^2}{2m} e^{-\frac{\hbar \lambda}{\lambda}}$$

可以观测到的能量越高, 可以相互作用的激发态越多.

$\phi^4$  理论

声子场.

超子 / 超流.

Spin / 磁性.

$$\mathcal{L} = \left(\frac{\partial\phi}{\partial t}\right)^2 - \left(\frac{\partial\phi}{\partial x}\right)^2 - \frac{\lambda}{4!}\phi^4$$

$$Z = \int D\phi e^{-iS[\phi]}$$

$$S = \int \mathcal{L} dx, \quad x = (\vec{x}, t).$$

$$\mathcal{L} = \frac{1}{2} \left(\frac{\partial\phi}{\partial t}\right)^2 - \frac{m^2}{2} \left(\frac{\partial\phi}{\partial x}\right)^2 - \underbrace{\frac{\lambda}{4!}\phi^4}_{-V}$$

$$= T - V.$$

Ref. Peskin. Chap 9.

Shankar. Chap 4.

Two views of renormalization

$m$ : mass.

$\lambda$ : interaction.

$$Z = \int D\phi e^{-iS[\phi]} \quad \text{如何计算?}$$

考虑 1 维的情形:

$$S = \int dx \mathcal{L}_0,$$

$$\mathcal{L}_0 = \frac{1}{2} \left(\frac{\partial\phi}{\partial t}\right)^2 - \frac{m}{2} \left(\frac{\partial\phi}{\partial x}\right)^2 - \frac{\mu}{2} \phi^2,$$

$\phi$ : real field.

$$\phi(x) = \frac{1}{V} \sum_{\mathbf{k}} \phi_{\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{x}} \Leftrightarrow \phi(x) = \phi^*(x). \quad \phi(\mathbf{k}) = \phi^*(\mathbf{k}).$$

$$\Rightarrow S = \int ( ) \phi_{\mathbf{k}} \phi_{\mathbf{k}'} d\mathbf{k} \leftarrow \text{计算如下:}$$

$$\phi(x) = \frac{1}{V} \sum_{\mathbf{k}} \phi(\mathbf{k}) e^{i\mathbf{k}\cdot\mathbf{x}}, \quad \mathbf{k}\cdot\mathbf{x} = \vec{k}\cdot\vec{x} - \omega t.$$

Ref. Parseval 定理

$$\int \frac{1}{2} \left(\frac{\partial\phi}{\partial t}\right)^2 dx = \frac{1}{2V^2} \sum_{\mathbf{k}, \mathbf{k}'} \phi_{\mathbf{k}} \phi_{\mathbf{k}'} (-i\omega_{\mathbf{k}}) (-i\omega_{\mathbf{k}'}) \int e^{i(\mathbf{k}+\mathbf{k}')\cdot\mathbf{x}} \cdot x dx$$

$$= \frac{1}{2V^2} (2\pi)^d \sum_{\mathbf{k}, \mathbf{k}'} \phi_{\mathbf{k}} \phi_{\mathbf{k}'} (-i\omega_{\mathbf{k}}) (-i\omega_{\mathbf{k}'}) \delta(\mathbf{k}+\mathbf{k}')$$

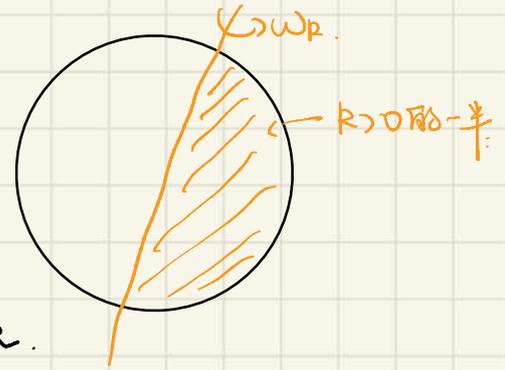
( $d$  为积分维度.)

和差化积分.  
 $\sum_R \omega = \frac{1}{2\pi} \int dk$

$$= \frac{1}{2V^2} (2\pi)^d \frac{V}{(2\pi)^d} \sum_R \int \phi_R \phi_{R'} (-i\omega_R) (-i\omega_{R'}) \delta(k+k') dR'$$

$$= \sim$$

原则上:  $k \rightarrow -k$   
 $\phi(k) \rightarrow \phi^*(k)$   
 $\phi_{-k} \rightarrow \phi_R^*$



$k \rightarrow -k$  没变化  
 所以(x2) =  $\frac{1}{V} \sum_{R \geq 0} \phi_R \phi_{-R} (-i\omega_R) (-i\omega_{-R}) dk$ .

$\omega_k = V|k|$   
 $\omega_k = \omega_{-k}$

$(\frac{\partial}{\partial t})$  贡献的仅仅是  $-i\omega_R$

同理  $\Rightarrow (\frac{\partial}{\partial x})$  贡献  $k$ .

$\Rightarrow S = \frac{1}{V} \sum_{R \geq 0} (-\omega_R^2 + m^2 \vec{R}^2 - \mu) \phi_R^* \phi_R$   
 (系数 1/2 消失了) 复数场 (且, 为对角)

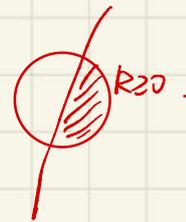
总结起来. 在  $\vec{R}$  空间中.

$$Z = \int \prod \phi_R e^{iS[\phi_R]}$$

$$S_0 = \sum_R \left( \frac{Q_R}{V} \right) \phi_R^* \phi_R \left. \begin{array}{l} = 0 \\ = \infty \end{array} \right\}$$

$$\prod_R D\phi_R = \prod_{R \geq 0} D\phi_R^* \phi_R$$

$$\int dx_1 \dots dx_n e^{-\sum_i b_i x_i^2}$$



在计算关联时会有用:

$$\langle \phi(x) \phi(y) \rangle = \frac{\int D\phi e^{iS_0} \phi(x) \phi(y)}{\int D\phi e^{iS_0}}$$

$$= \sum_{q, q'} e^{iq \cdot x} e^{iq' \cdot y} \frac{\int e^{iS_0} \phi_q \phi_{q'} D\phi_q D\phi_{q'}}{\int e^{iS_0} D\phi_q D\phi_{q'}}$$

$$= \sum_{q, q'} e^{iq \cdot x + iq' \cdot y} \frac{1}{a} \delta(q, q')$$

$$\left( a = \frac{Q_R}{V} \right)$$

$$= \frac{1}{V} \sum_q e^{iq(x-y)} \frac{1}{Q_R}$$

2021.3.10 第三周第二节课.

Parseval formula.

$$f(x) = \frac{a_0}{\sqrt{2}} + \sum_{n=1}^{\infty} a_n \cos(nx) + \sum_{n=1}^{\infty} b_n \sin(nx)$$

$$\langle f, f \rangle = \int f(x)^2 dx$$

or  $\int f(x)^2 dx = \int f(k)^2 dk$ .

今天的内容: 1. Green's function.

Propagator.  $D(x-y) = \frac{1}{V} \sum_{\vec{k}} ( )$   
 图表示.  $\underbrace{\hspace{2cm}}_{\text{成对出现}}$

2. Interaction.

Cumulant Expansion (Moment Expansion).

物理: linked cluster expansion.

意义:

$$\phi(x) = \frac{1}{V} \sum_{\vec{k}} e^{i\vec{k} \cdot \vec{x}} \phi(\vec{k}), \quad \vec{k} = (k_0, \vec{k}). \text{ 对 } \vec{k} \text{ 之前的 } \omega_{\vec{k}}. \quad x = (t, -\vec{x}).$$

$$S_0 = \int \mathcal{L}_0 dx$$

$$\int \frac{1}{2} \left( \frac{\partial \phi}{\partial t} \right)^2 dx = \frac{1}{2V^2} \sum_{\vec{k}, \vec{k}'} \phi(\vec{k}) \phi(\vec{k}') (ik_0)(ik_0') \int e^{i(\vec{k} + \vec{k}') \cdot \vec{x}} dx$$

$$= \frac{1}{2V^2} \sum_{\vec{k}, \vec{k}'} \phi(\vec{k}) \phi(\vec{k}') (ik_0)(ik_0') \delta(\vec{k} + \vec{k}') \cdot (2\pi)^d$$

$$= \frac{1}{2V^2} \sum_{\vec{k}} \phi(\vec{k}) \frac{V}{(2\pi)^d} \int \phi(\vec{k}') (ik_0)(ik_0') \delta(\vec{k} + \vec{k}') d\vec{k}'$$

要求  $k' = -k$ .  
 $\rightarrow k_0' = -k_0$

$$= + \frac{1}{2V} \frac{1}{(2\pi)^d} \sum_{\vec{k}} \phi(\vec{k}) \phi(-\vec{k}) k_0^2$$