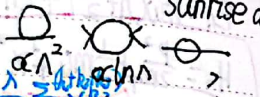
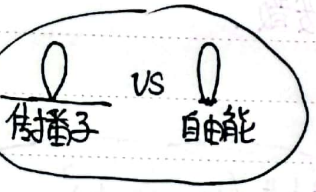


Dyson eq: $g = g_0 + g_0 \Sigma g \Rightarrow$ 右乘 $g^{-1} \Rightarrow 1 - g_0 g^{-1} = g_0 \Sigma \Rightarrow$
 $\begin{cases} g = \frac{1}{p^2 - m^2} \\ g_0 = \frac{1}{p^2 - m_0^2} \end{cases}$
 左乘 $g_0^{-1} \Rightarrow g_0^{-1} - g^{-1} = \Sigma$
 即 $p^2 - m_0^2 - (p^2 - m^2) = \Sigma$
 $m^2 = m_0^2 + \Sigma$
 观测值 \wedge 自由能

$\langle 0 \rangle = \sum_{n=0}^{\infty} \frac{z^n}{n!} \langle OS_1^n \rangle^c$ $S_1 = \frac{-\lambda}{4!} \int \phi^4 dx$
 $\langle \phi^*(p) \phi(p) \rangle$: ① $V=1 \Rightarrow \frac{1}{V}$ 成对出现
 ② 舍弃许多图 (哪些图)

③ 理解 Feynman diag / 基本图  sunrise diag
 $S_1 = \frac{-\lambda}{4!} \int \phi^4 dx = \frac{-\lambda}{4!} \int \frac{1}{k^2} \frac{1}{k^2} \frac{1}{k^2} \frac{1}{k^2} dx$
 $0 = \phi^*(p) \phi(p)$, $\langle 0 \rangle = \langle 0_0 \rangle + \frac{z}{1!} \langle \frac{-\lambda}{4!} \int \phi^4 dx \rangle^c + \dots$
 $= \frac{z}{p^2 - m^2 - i\epsilon} + \frac{z^2}{1!} \langle \frac{-\lambda}{4!} \int \phi^4 dx \rangle^c + \dots$
 $= g_0(p) - \frac{i\lambda}{2} g_0(p) g_0(p) + \frac{(-i\lambda)^2}{(2!)^2} \dots$



再考虑3阶 $\frac{z^3}{3!} \langle \frac{-\lambda}{4!} \int \phi^4 dx \rangle^3$
 $= \frac{z^3}{3!} \langle \frac{-\lambda}{4!} \int \phi^4 dx \rangle^3$
 $= \frac{(-i\lambda)^3}{(2!)^3} \dots$

$\Omega = \frac{-i\lambda}{2} \frac{z}{k^2 - m^2}$

$\Omega \Omega \Omega = g_0 \Sigma g_0 \Sigma g_0$

相互作用: $\frac{-\lambda}{4!}$, $0 = \langle \phi(p_1) \phi(p_2) \phi(p_3) \phi(p_4) \rangle$, p_i 任意
 唯一要求 $\Sigma p_i = 0$
 指任意允许的散射过程 $\langle 0 \rangle_0 = 0$

$\langle 0 \rangle = \langle 0 \rangle_0 + \frac{z}{1!} \langle \frac{-\lambda}{4!} \int \phi^4 dx \rangle^c$
 $= \dots + \frac{z^2}{2!} \langle \frac{-\lambda}{4!} \int \phi^4 dx \rangle^2$
 $= -i\lambda X$

= 阶: $\frac{z^2}{2!} \langle \frac{-\lambda}{4!} \int \phi^4 dx \rangle^2$
 $= \frac{z^2}{2!} \langle \frac{-\lambda}{4!} \int \phi^4 dx \rangle^2$
 $= -\frac{(\Sigma)^2}{2}$
 或 $\langle \phi \phi \phi \phi \rangle \rightarrow 0$

或 $\langle \phi^*(p) \phi(p) \rangle$ 或 $\langle \phi^*(p) \phi(p) \rangle$
 $\frac{z^2}{2!} \langle \frac{-\lambda}{4!} \int \phi^4 dx \rangle^2$
 $\frac{\delta \ln \Omega}{\alpha \Lambda^2 \ln \Lambda}$
 $\alpha = 4$
 $g_0(p) \Sigma_2 g_0(p)$
 $\frac{g_0(p) \Sigma_2 g_0(p)}{g_0(p) \Sigma_2 g_0(p)} = \text{阶数贡献}$

$\langle 0 \rangle = -\text{阶} + \text{阶}$
 $= -i\lambda X - \frac{3\lambda^2}{2} X X \approx -i\lambda_{\text{eff}} X$

其中 $-i\lambda_{\text{eff}} = -i\lambda - \frac{3\lambda^2}{2} X$ 解决方法
 观测到的散射长度

$m^2 = m_0^2 + \frac{g}{\alpha \Lambda^2}$ (平重发散)

2d $S(x)$ potential $\frac{g}{V} \frac{1}{k^2 - \frac{E - \frac{\hbar^2 k^2}{2m}}{m_0}}$, $E = \frac{\hbar^2 \lambda^2}{2m}$
 $\Omega \propto \frac{1}{k^2 - m^2}$



重整化: ① 系统化方法.

$\mu \ll \Lambda$ ② 谁抵消谁

③ 避免复杂的非线性相互作用. (方程)

$\mathcal{L} = \frac{1}{2}(\partial_\mu \phi)^2 - \frac{m^2}{2}\phi^2 - \frac{\lambda}{4!}\phi^4$, $m_0, \lambda_0 \Rightarrow \Lambda$ 在能标 Λ 下定义的

令 $\phi = \sqrt{Z}\phi_R$, $\langle \phi^*(p)\phi(p) \rangle = \frac{Z}{p^2 - m^2} + \text{finite}$

$\mathcal{L} = \underbrace{\left[\frac{1}{2}(\partial_\mu \phi_R)^2 - \frac{m_R^2}{2}\phi_R^2 \right]}_{\mathcal{L}_0 \text{ 观测值 } (\mu, \Lambda) \text{ 实验可见}} + \underbrace{\left[\frac{\lambda_R}{4!}\phi_R^4 + \frac{1}{2}(\delta_Z(\partial_\mu \phi_R)^2 - \delta_m \phi_R^2) - \frac{\delta_\lambda}{4!}\phi_R^4 \right]}_{\text{抵消项 } \delta_Z, \delta_m, \delta_\lambda \text{ 不可测}}$

图 $\langle \phi_R^*(p)\phi_R(p) \rangle = \frac{1}{p^2 - m_R^2}$

$\times = -i\lambda_R$

$\otimes = -i\delta_\lambda$

$\otimes = (\delta_Z p^2 - \delta_m)$

$S_0 = \int \mathcal{L}_0 dx$, $S_I = -\frac{\lambda_R}{4!}\phi_R^4 - \frac{\delta_\lambda}{4!}\phi_R^4 + \frac{\delta_Z}{2}(\partial_\mu \phi)^2 - \frac{\delta_m}{2}\phi^2$

$\langle 0 \rangle = \langle \phi_R^*(p)\phi_R(p) \rangle_0$

$\frac{1}{g(p)} = \frac{1}{g_0(p)} + \frac{1}{\lambda_R} + \frac{\otimes}{\delta_\lambda} + \frac{\otimes}{(\delta_Z p^2 - \delta_m)}$

$\Sigma = \frac{1}{\lambda_R \Lambda^2} + \frac{\otimes}{\delta_\lambda \Lambda^2} + \frac{\otimes}{(\delta_Z p^2 - \delta_m)}$

$m_R^2 = m^2 + \frac{1}{\lambda_R} + \frac{\otimes}{\delta_\lambda} + \frac{\otimes}{(\delta_Z p^2 - \delta_m)}$

有限值 = 0

*: $H = \left[\frac{p^2}{2m} + \frac{1}{2}m\omega_k^2 x^2 + \lambda x^4 + \frac{1}{2}m\delta\omega^2 x^2 \right]_{V_I}$
 $= \hbar\omega_k(n + \frac{1}{2}) + \langle n | V_I | n \rangle$

$\times = \frac{1}{\lambda_R} + \left[\frac{\otimes}{\delta_\lambda} + \frac{\otimes}{\lambda_R} + \frac{\otimes}{\lambda_R} \right] = 0$

则 $\delta_\lambda = -\lambda \alpha$
 $\propto \ln \Lambda$

