

因为场 $(\phi(x,t))$ 同样是无穷多自由度，考虑将 $S[x]$ 美化
到均论中

$$\langle q_1 q_2 | e^{-iHt} | q'_1 q'_2 \rangle = \int Dq_1 Dq_2 e^{iS[q_1, q_2]}$$

$$\lim_{N \rightarrow \infty} \langle q_1 \dots q_N | e^{-iHt} | q'_1 \dots q'_N \rangle = \int D\phi e^{iS[\phi]}$$

只有变换到动量空间，才有可能继续做下去

2.28

$$\prod_i \int D\phi(x_i) = J \prod_K \int D\phi(K) \quad \phi(K) = \sum_i e^{ikx_i} \phi(x_i)$$

变为 Gauss integral Feynmann diagram
无穷维只有 Gauss 积分有解析闭形势

Path integral

$$\begin{array}{cccc} \textcircled{1} \text{ Boson Field} & \textcircled{2} \text{ Fermion Field} & \textcircled{3} \text{ Gauge Field} & \textcircled{4} \text{ Spin Field} \\ \int D\psi D\psi^+ & & \int DA & \end{array}$$

Boson field: Gauss integral (Ref. A. Zee)

$$\int dx e^{-Ax^2} = \sqrt{\frac{\pi}{A}}$$

$$\int e^{-Ax^2 + Jx} dx = \dots = \sqrt{\frac{\pi}{A}} e^{\frac{J^2}{4A}} \rightarrow \text{左、右 Taylor series}$$

$$\int e^{-Ax^2} x^{2n} dx = \frac{J^{2n}}{(2n)!} \text{ 右边展开 } \sum_{n=0}^{+\infty} \int e^{-Ax^2} \frac{J^{2n}}{(2n)!} x^{2n} dx$$

$$\langle x^{2n} \rangle = \frac{(2n)!}{n! 2^n} \left(\frac{J}{2A}\right)^n$$

$$\text{可用 } 2n \text{ 个虚而配对, 强度为 } \frac{1}{2A} \text{ 倍. } \frac{G_n^2 G_{n-2}^2 \dots G_2^2}{n!} \left(\frac{J}{2A}\right)^n$$



source / perturbation
FUTURE PEOPLE

$$\frac{\int e^{-Ax^2 + \bar{J}x} x^2 dx}{\int e^{-Ax^2} dx} = \dots = \frac{1}{2A}$$

① $\frac{\int dx_1 \dots dx_n e^{-\frac{1}{2}x^T A x} x_i x_j}{\int dx_1 \dots dx_n e^{-\frac{1}{2}x^T A x}} = \langle x_i x_j \rangle$ Gauss integral
 \Rightarrow multivariable

② $\frac{\int d\bar{z} d\bar{z} e^{-\frac{1}{2}\bar{z}^T A \bar{z}} \bar{z} \bar{z}}{\int d\bar{z} d\bar{z} e^{-\frac{1}{2}\bar{z}^T A \bar{z}}} = \langle \bar{z} \bar{z} \rangle$ \Rightarrow Complex

③ $\frac{\int d\bar{z}_1 d\bar{z}_2 \dots d\bar{z}_i d\bar{z}_j e^{-\frac{1}{2}\bar{z}^T A \bar{z}} \bar{z}_i \bar{z}_j}{\int e^{-\frac{1}{2}\bar{z}^T A \bar{z}}} = \langle \bar{z}_i \bar{z}_j \rangle \Rightarrow$

计算: $\int dx_1 \dots dx_n e^{-\frac{1}{2}x^T A x} = I$

变换 $x = Oy$ $O \in O(N)$

$y = O^T x$ Jacobi det = 1

$I = \int \left(\frac{\partial x_i}{\partial y_j} \right) dy_1 dy_2 \dots dy_n e^{-\frac{1}{2}x^T A x}$

设 A 对角化: $x^T A x = y^T O^T A O y \rightarrow \text{diag}: \begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_n \end{pmatrix}$

$= \int \left(\frac{\partial x_i}{\partial y_j} \right) dy_1 \dots dy_n e^{-\frac{1}{2} \sum_i y_i^2 \lambda_i} = \sqrt{\frac{(2\pi)^n}{\prod \lambda_i}}$



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$$\text{计算 } \int d\lambda_1 \dots d\lambda_n e^{-\frac{1}{2}x^T A x + J^T x}$$

$$e^{-\frac{1}{2}x^T A x + J^T x} = e^{-\frac{1}{2}(x^T + C^T) A (x + C) + J^T x + J^T C}$$

$$= e^{-\frac{1}{2}x^T A x - \frac{1}{2}C^T A C - \frac{1}{2}x^T A C - \frac{1}{2}C^T A x + J^T x + J^T C}$$

$$C = A^{-1}J$$

$$= e^{-\frac{1}{2}x^T A x - \frac{1}{2}C^T A C + J^T C} \quad \text{注: 这几项相消}$$

$$\langle x_i, x_j \rangle = \frac{\frac{\partial}{\partial J_i} \frac{\partial}{\partial J_j} \int e^{-\frac{1}{2}x^T A x + J^T x} Dx}{\int e^{-\frac{1}{2}x^T A x} Dx}$$

$$= \left(\frac{\partial}{\partial J_i} \frac{\partial}{\partial J_j} e^{\frac{1}{2}J^T + A^{-1}J} \right) \Big|_{J=0}$$

$$= \left(\frac{\partial}{\partial J_i} \frac{\partial}{\partial J_j} e^w \right) \Big|_{J=0}$$

$$\left(\frac{\partial}{\partial J} e^w = \frac{\partial w}{\partial J} e^w \right) \quad \frac{\partial}{\partial J_i} \frac{\partial}{\partial J_j} e^w = \frac{\partial}{\partial J_i} \left(\frac{\partial w}{\partial J_j} e^w \right)$$

$$= \frac{\partial w}{\partial J_i} \frac{\partial w}{\partial J_j} e^w + \frac{\partial^2 w}{\partial J_i \partial J_j} e^w$$

推广到复数

$$e^{-\frac{1}{2}z^T A z + J^T z + z^T J} \quad z = \begin{pmatrix} z_1 \\ \vdots \\ z_n \end{pmatrix} \quad z^T = (z_1^T \dots z_n^T)$$

$$\Rightarrow \int Dz D\bar{z} e^{-\frac{1}{2}z^T A z + J^T z + \frac{1}{2}A^{-1}J}$$



相互作用与微扰

$$\int dx e^{-\frac{1}{2}Ax^2 - \frac{\lambda}{4!}x^4} = e^{-F}$$

$$= \sum_{n=0}^{\infty} \int dx e^{-\frac{1}{2}Ax^2} \left(-\frac{\lambda}{4!}\right)^n \frac{1}{n!} x^{4n}$$

3.3.

$$\mathcal{L} = \left(\frac{\partial \phi}{\partial T}\right)^2 - V^2 \left(\frac{\partial \phi}{\partial x}\right)^2 - m^2 \phi^2 - \frac{\lambda}{4!} \phi^4$$

☆ φ 像 r 坐标 / 振动大小

$$\star Z = \int D\phi e^{iS} \quad S = \int \mathcal{L} d\vec{x} dt$$

Path integral

一般为复数，偶尔有反常

☆ 求解：动量空间 $\prod_i d\phi(x_i) \rightarrow \int \prod_k d\phi(k)$

$$\text{J 为常数时不重要。因为 } \langle Q \rangle = \frac{\int P(x) Q dx}{\int P(x) dx}$$

唯一可解 N 维积分 \Rightarrow Gauss 分布

恒等式方程

$$\textcircled{1} \quad \frac{\int dx e^{-\frac{A}{2}x^2} x^2}{\int dx e^{-\frac{A}{2}x^2}} = \frac{1}{A} \quad \textcircled{2} \quad \frac{\int d\vec{z} d\bar{\vec{z}} e^{-A\vec{z}\bar{\vec{z}}}}{\int d\vec{z} d\bar{\vec{z}} e^{-A\vec{z}\bar{\vec{z}}}} = \frac{1}{A}$$

发散问题

目的：说明 ∞ 处不在e.g. 点电荷产生电场的能量 $E \propto \int \frac{e^2}{r^4} d^3\vec{r} = 9\pi(\frac{1}{r}) \Big|_0^{+\infty}$ 例 1. 2 dim $\delta(\vec{x})$ potential

ref. Nyeo. Am.J.Phys. 2000.



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