

$$= \frac{1}{2\pi a} e^{\frac{i}{2} L} \langle 0 | e^{-i(\Phi_+(x+\epsilon) - \Phi_+(x)) + i(\Phi_-(x+\epsilon) - \Phi_-(x))} | 0 \rangle$$

$$= \langle 0 | e^{i(\Phi_+(x+\epsilon) - \Phi_+(x))} e^{i\Phi_- - i\Phi_+} | 0 \rangle$$

$$\rho \sim \frac{1}{2\pi a} e^{-i\Phi(x+\epsilon) + i\Phi(x)} \sim \frac{1}{2\pi a} e^{-i(\frac{\partial\Phi}{\partial x})\epsilon} = \frac{1}{2\pi a} (1 - i\frac{\partial\Phi}{\partial x}\epsilon + \frac{1}{2}(\frac{\partial\Phi}{\partial x})^2 \epsilon^2 + \dots) = \frac{1}{2\pi a} - \frac{1}{2\pi}(\partial_x \Phi)$$

为什么 $\psi \sim e^{i\Phi}$ 是正确的?

1) 反对易

2) Hilbert space

3) $Z_F = Z_B$ (Bloch, 1933)

$$\star \quad \{\psi(x), \psi^\dagger(y)\} = \delta(x-y)$$

$$\star \quad \text{关联函数一样: } \langle \psi^\dagger(x) \psi(y) \rangle = \frac{1}{2\pi a} \langle e^{-i\phi(x)} e^{i\phi(y)} \rangle$$

⇓

应用

字典

$$\mathcal{L} = \frac{1}{2}(\partial\phi)^2 + A \cos(\beta\phi)$$

$$[\Phi(x), \Phi(y)] = \ln \left(\frac{x-y-ia}{y-x-ia} \right) = \Phi_+(x) + \Phi_-(x)$$

$$\Leftrightarrow \Phi_+[\Phi_+(x), \Phi_-(y)] + [\Phi_-(x), \Phi_+(y)]$$

$$\Leftrightarrow [\Phi_+(x), \Phi_-(y)] = \ln(x-y-ia)$$

1) 关联函数一样

2) 给出 $[\Phi(x), \Phi(y)]$ 表达式

3) 翻译 $\psi \rightarrow \Phi$ 场 明确物理意义

$$\langle \psi^\dagger(x) \psi(y) \rangle = \frac{1}{L} \sum_{q>0} e^{-iq(x-y)} \langle 0 | c_q^\dagger c_q | 0 \rangle$$

$$= \frac{1}{L} \sum_{q>0} e^{-iq(x-y) + aq} \int_{-\infty}^0 e^{Aq} dq = \frac{1}{A}$$

$$= \frac{1}{2\pi} \int_{-\infty}^0 e^{[a-i(x-y)]q} dq = \frac{1}{2\pi} \frac{1}{a-i(x-y)} = \frac{1}{2\pi} \frac{1}{x-y+ia}$$

$$= \frac{1}{2\pi a} \langle e^{-i\phi(x)} e^{i\phi(y)} \rangle = \frac{1}{2\pi a} e^k$$

$$e^k = \frac{ia}{x-y+ia}$$

$$\text{or } k = \ln(ia) - \ln(x-y+ia)$$

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) g = \delta(x-y)$$

$$g \sim \frac{1}{2\pi} \ln r$$

$$\Phi = i \sum_{q>0} \sqrt{\frac{2\pi}{Lq}} (e^{-iqx} b_q^\dagger - e^{iqx} b_q)$$



Remarks:

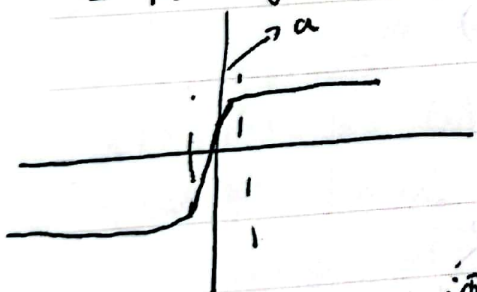
1) $x \sim a+at$
 $p \sim i(a-at)$

2) $\rho \sim \left(\frac{\partial \Phi}{\partial x}\right) \sim \sum_{q>0} \rho_q e^{-iqx} + \rho_q^\dagger e^{iqx}$

3) 为什么 $\frac{1}{L}$

第一步: $[\Phi(x), \Phi(y)] = - \sum_{q>0} \frac{2\pi}{Lq} [e^{-iqx} b_q^\dagger - e^{iqx} b_q, e^{-iqy} b_q^\dagger - e^{iqy} b_q]$
 $= - \frac{2\pi}{L} \sum_{q>0} \left(\frac{1}{q} e^{-iq(x-y)-aq} - \frac{1}{q} e^{iq(x-y)-aq} \right)$ ① $q = \frac{2\pi n}{L}$ ② 求导不为0
 $= - \int_0^{+\infty} dq \frac{e^{Aq} - e^{Bq}}{q} = \int_0^{+\infty} \frac{e^{Aq} - e^{Bq}}{a} dq$

$A = i(x-y) - a$ $B = -i(x-y) - a$
 ① $= - \sum_{n>0} \left(\frac{1}{n} e^{(-i(x-y)-a)\frac{2\pi n}{L}} - \frac{1}{n} e^{(i(x-y)-a)\frac{2\pi n}{L}} \right)$
 $\ln(1-x) = - \sum_{n>0} \frac{1}{n} x^n = \ln(1 - e^{\frac{2\pi}{L}(-i(x-y)-a)}) = \ln(1 - e^{\frac{2\pi}{L}(i(x-y)-a)})$
 $= \ln\left(\frac{2\pi}{L}(i(x-y)+a)\right) - \ln\left(\frac{2\pi}{L}(i(y-x)+a)\right)$
 $= \ln(x-y-ia) - \ln(y-x-ia) = \ln\left(\frac{x-y-ia}{y-x-ia}\right)$



关联函数: $\langle e^{-i\Phi(x)} e^{i\Phi(y)} \rangle$ 合 $\rightarrow \hat{n} = \langle \dots \rangle$
 Normal ordering $: A := A - \langle 0|A|0 \rangle$ $e^{i\Phi} := e^{\hat{\Phi}_+} e^{\hat{\Phi}_-}$

正规化: $\langle 0|e^{-i\Phi(x)+i\Phi(y)} e^{\frac{1}{2}[\Phi(x), \Phi(y)]} |0 \rangle = e^{\frac{1}{2}[\Phi(x), \Phi(y)]} \langle 0|e^{-i\Phi(x)+i\Phi(y)} |0 \rangle$
 $e^A = e^{A_+ + A_-} = e^{A_+} e^{A_-} e^{-\frac{1}{2}[A_+, A_-]}$
 $= e^{\frac{1}{2}[\Phi(x), \Phi(y)]} e^{-\frac{1}{2}[A_+, A_-]} \langle 0|e^{A_+} e^{A_-}|0 \rangle = \frac{i}{2\pi(x-y-ia)}$
 $K = \frac{1}{2} [\Phi(x), \Phi(y)] - \frac{1}{2} [A_+, A_-] = \frac{1}{2} [\Phi_+(x), \Phi_+(y), \Phi_-(x) - \Phi_-(y)] = \frac{1}{2} \ln\left(\frac{x-y-ia}{y-x-ia}\right) + \frac{1}{2} [2\ln(-ia) - \ln(x-y-ia) - \ln(y-x-ia)] = \ln(-ia) - \ln(y-x-ia)$



$$\langle \psi^\dagger(x) \psi(y) \rangle = \frac{i}{2\pi} \frac{1}{x-y+ia} \quad \langle \psi(y) \psi^\dagger(x) \rangle = -\frac{i}{2\pi} \frac{1}{y-x+ia}$$

$$\langle \{ \psi(y), \psi^\dagger(x) \} \rangle = \delta(x-y) = \frac{i}{2\pi} \left(\frac{1}{x-y+ia} + \frac{1}{y-x+ia} \right)$$

$$= \frac{i}{2\pi} \left(\frac{1}{x-y+ia} - \frac{1}{x-y-ia} \right) = \delta(x-y)$$

X X Z model

$$-J \sum_i b_i^x b_{i+1}^x + b_i^y b_{i+1}^y + \Delta \sum_i b_i^z b_{i+1}^z \quad \text{Jordan - Wigner}$$

$$H = -t \sum_i (c_i^\dagger c_{i+1} + h.c.) + \mu c_i^\dagger c_i + U c_i^\dagger c_i c_{i+1}^\dagger c_{i+1}$$

$$H_0 = [-2t \cos(ka) + \mu] C_k^\dagger C_k \quad \mu=0$$

$$(-2t \cos(ka))' = 2ta \sin(ka) \quad ka = \frac{\pi}{2}$$

$$= 2ta = v_F$$

$$H_0 = v_F \psi_R^\dagger (-i \frac{\partial}{\partial x}) \psi_R - v_F \psi_L^\dagger (-i \frac{\partial}{\partial x}) \psi_L = v_F k C_k^\dagger (k - v_F k) d_k^\dagger d_k$$

$$\mathcal{L} = \frac{1}{2} [(\partial_t \phi_R)^2 - v_F^2 (\partial_x \phi_R)^2] + \frac{1}{2} [(\partial_t \phi_L)^2 - v_F^2 (\partial_x \phi_L)^2]$$

$$\sum_i c_i^\dagger c_{i+1} \leftrightarrow \int \psi^\dagger(x) \left(\frac{\partial}{\partial x} \right) \psi \quad c_i \sim \psi(x=ia)$$

$$\sum_i c_i^\dagger c_i = \int \psi^\dagger(x) \psi(x) dx \quad c_i = \Lambda \psi(x=ia)$$

$$= \Lambda^2 \sum_i (\psi^\dagger(x=ia) \psi(x=ia))$$

$$\sum_i -t c_i^\dagger c_{i+1} = -ta \sum_i \psi^\dagger(x=ia) \psi(x=(i+1)a) + h.c.$$

$$= -t \int \psi^\dagger(x=ia) \psi(x=ia+a) dx + h.c.$$

$$\simeq -t \int \psi^\dagger \frac{\partial^2}{\partial x^2} \psi$$

$$\sum_i c_i^\dagger c_i c_{i+1}^\dagger c_{i+1} = (\sqrt{a})^4 \sum_i \psi^\dagger(x=ia) \psi(x=ia) \psi^\dagger(x=ia+a)$$

$$\psi(x=ia+a) = a \int dx \psi^\dagger(x) \psi(x) \psi^\dagger(x+a) \psi(x+a)$$

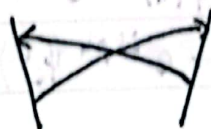
$$\psi = e^{ik_F x} \psi_R + e^{-ik_F x} \psi_L \quad \psi \sim \frac{F_i}{\sqrt{2\pi a}} e^{i\Phi_i}$$

$$= \left(\cancel{p} \leftarrow + \leftarrow \right) \times \left(\cancel{p} \leftarrow + \leftarrow \right)$$

$e^{-i2k_F x} \quad e^{i2k_F x} \quad e^{-i2k_F(x+a)} \quad e^{i2k_F(x+a)}$

forward scattering $\sim (\partial \phi^2) \sim p^2$

backward scattering



Wengu



$$(s \cdot s + s \cdot f + f \cdot s + f \cdot f)$$

$$h.c. + \psi_L^\dagger \psi_R = \frac{F_L F_R}{2\pi a} e^{-i\Phi_L + i\Phi_R} + \frac{F_R F_L}{2\pi a} e^{-i\Phi_R + i\Phi_L}$$

$$= \frac{F_L F_R}{2\pi a} (e^{i\Phi_L - i\Phi_R} - e^{i\Phi_R - i\Phi_L}) = \frac{i F_L F_R}{\pi a} \sin(\Phi_L - \Phi_R)$$

今天目的：如何直接写出来：

$$\left. \begin{matrix} L(\psi) \\ H(\psi) \end{matrix} \right\} \Rightarrow \left. \begin{matrix} L(\Phi) \\ \text{物理意义} \end{matrix} \right\}$$

字典： ψ 相互作用 \Rightarrow Bosonized L

关键：normal ordering

① Jordan-Wigner string

$$e^{i\pi \int_{-\infty}^x \rho(x') dx'} \quad \phi = \pi \int_{-\infty}^x \rho(x') dx' \Leftrightarrow \frac{\partial \phi}{\partial x} = \pi \rho$$

$$\psi^\dagger = \frac{F}{\sqrt{2\pi a}} e^{-i\phi} \quad (F^\dagger = F \quad F^2 = 1) \quad \psi^\dagger \psi = \frac{F^2}{2\pi a} e^{-i\phi(x)} e^{i\phi(x)}$$

$$\rho = \lim_{a \rightarrow 0} \psi^\dagger(x+a) \psi(x) = \frac{F^2}{2\pi a} e^{-i\phi(x+a)} e^{i\phi(x)} = \frac{1}{2\pi a}$$

$$= \frac{1}{2\pi a} e^{-i\phi(x+a) + i\phi(x) + \frac{1}{2} [\phi(x+a), \phi(x)]} = \frac{1}{2\pi a} e^{\frac{1}{2} [\phi(x+a), \phi(x)]} e^{-i(\frac{\partial \phi}{\partial x})}$$

$$= \frac{1}{2\pi a} e^{\frac{1}{2} [\psi(x+a), \psi(x)]} (1 - i a \frac{\partial \phi}{\partial x})$$

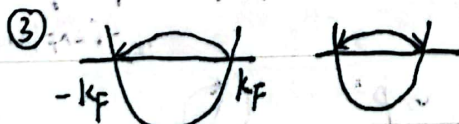
$$= \text{const} - \frac{i\gamma}{2\pi a} e^{\frac{1}{2} [\phi(x-y), \phi(x)]} (\frac{\partial \phi}{\partial x})$$

$$\approx \text{const} + \frac{1}{2\pi} (\frac{\partial \phi}{\partial x})$$

$$\rho = \psi^\dagger(x+y) \psi(x) := A = \langle 0 | A | 0 \rangle = \frac{1}{2\pi} (\frac{\partial \phi}{\partial x})$$

字典：① $\rho \rightarrow \rho := \psi^\dagger \psi \rightarrow (\frac{\partial \phi}{\partial x})$

② Luttinger liquid $(\psi^\dagger \partial_x \psi) \leftrightarrow (\partial_x \phi)^2$
 1次 \rightarrow 2次 (ϕ^2)



背散射

pairing

$$V \cos(2k_F x) \psi^\dagger(x) \psi(x) \quad \psi(x) = e^{ik_F x} \psi_R + e^{-ik_F x} \psi_L$$

$$V \cos(2k_F x) (e^{-ik_F x} \psi_R^\dagger + e^{ik_F x} \psi_L^\dagger) (e^{ik_F x} \psi_R + e^{-ik_F x} \psi_L)$$

