

2021. 9. 20.

目的: $\{-\frac{\hbar^2}{2m}\nabla^2 + V(x)\}\psi(x) = E\psi(x)$

Method: 有限差分方法.

要求 1D, 2D Schrödinger Equation

缺点: 内存随维度增长而增长.

3D: $\underline{N^3}$

难以接受.

步骤: ① 无量纲化待求解的方程.

② 离散化无量纲后的方程.

$$\begin{cases} \frac{\partial \psi(x)}{\partial x} \approx \frac{\psi(x+h) - \psi(x)}{h} + o(h) \\ \frac{\partial^2 \psi(x)}{\partial x^2} \approx \frac{\psi(x+h) - 2\psi(x) + \psi(x-h)}{h^2} + o(h^2) \end{cases}$$

取一个有限的 h , 而不是 $h \rightarrow 0$ (有限差分)

Schrödinger 方程 $(-\frac{\hbar^2}{2m}\frac{\partial^2}{\partial x^2} + V(x))\psi(x) = E\psi(x)$
 \Downarrow

$$-\frac{\hbar^2}{2m}(\psi_{i-1} + \psi_{i+1} - 2\psi_i) + V_i\psi_i = E\psi_i$$

$$H|\psi\rangle = E|\psi\rangle$$

$$H = \begin{pmatrix} \frac{1}{2h^2} & -\frac{1}{2h^2} & & & \\ -\frac{1}{2h^2} & \frac{1}{2h^2} & & & \\ & & \ddots & & \\ & & & \ddots & \\ & & & & \ddots \end{pmatrix} + \begin{bmatrix} V_1 \\ V_2 \\ \vdots \\ V_N \end{bmatrix}$$

特点: (1d)

1) 稀疏矩阵, 三对角矩阵.

2) $\hbar \rightarrow 0$ 时, 动能似乎发散, 如何理解?

作业.

① $V(x) = \frac{1}{2} m \omega^2 x^2$

比较数值结果与精确结果的误差. (与离散长度 h ?

量子数 n ?)

误差 10^{-7} . $h \sim 10^{-3}$ $m \sim 10^{-4}$

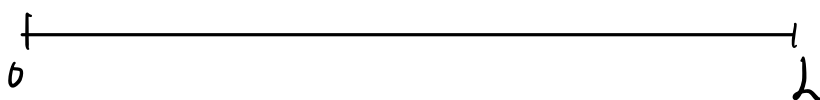
② $V(x) = \frac{1}{2} m \omega^2 x^2 + A \cos(kx + \theta)$

① 数值

② 微扰解 (A 小时)

低能的状态. 精确至二阶

考虑



$$\begin{cases} \psi(0) = \psi(L) = 0 \\ \psi'(0) = \psi'(L) = 0 \\ -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi(x) = E \psi(x) \end{cases}$$

解是 $\psi \sim \sin\left(\frac{2n\pi x}{L}\right)$, $E = \frac{n^2 \pi^2 \hbar^2}{2mL^2}$

离散化

$$H = -\frac{1}{2h^2} \begin{pmatrix} -2 & 1 & & & \\ 1 & -2 & 1 & & \\ & & \ddots & \ddots & \\ & & & \ddots & \ddots \\ & & & & -2 & 1 \\ & & & & 1 & -2 \end{pmatrix}$$

对角化后得到 $\bar{E}(h)$. 而

$$\bar{E} = \lim_{h \rightarrow 0} \bar{E}(h) = \frac{\hbar^2 k^2}{2m}$$

方程形式为

$$-\frac{1}{2h^2} (\psi_{i+1} + \psi_{i-1} - 2\psi_i) = E\psi_i$$

★ 解是什么样子的.

$$\begin{cases} a_{n+1} = x a_n + y a_{n-1} \\ \& a_n = q^n. \end{cases}$$

$$\& \psi_n = e^{iknh}$$

$$\Rightarrow -\frac{1}{2h^2} (e^{ikh} + e^{-ikh} - 2) = E$$

$$\Rightarrow \bar{E}(h) = \frac{1}{h^2} (1 - \cos(kh))$$

$$\begin{aligned} &\approx \frac{1}{h^2} (1 - 1 + \frac{1}{2} k^2 h^2 + o(h^2)) \\ &\stackrel{\lim_{h \rightarrow 0}}{=} \frac{1}{2} k^2 \end{aligned}$$

边界条件.

$$\psi_n = A e^{iknh} + B e^{-iknh}$$

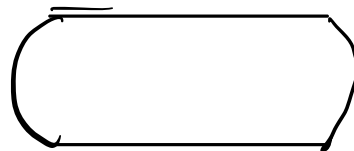
$$\begin{cases} -2\psi_1 + \psi_2 = E\psi_1 \\ \psi_{N-1} - 2\psi_N = E\psi_N \end{cases}$$

\Rightarrow 代入求解出 k

$$\Rightarrow E_n = \frac{\hbar^2 k^2}{2m^2}$$

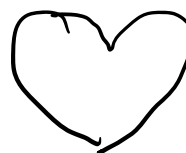
作业. 2D 下推了

1)



操场

2)



心型.

由心势能为零, 波函数在边界上

$$V(x) \begin{cases} 0, & \text{区域内} \\ \infty, & \text{区域外} \end{cases}$$

\wedge
1000 ~ 10000, 足够

二维下. 的偏微分方程.

$$-\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) \psi(x,y) + V(x,y) \psi(x,y) = E \psi(x,y)$$

\Rightarrow 离散化

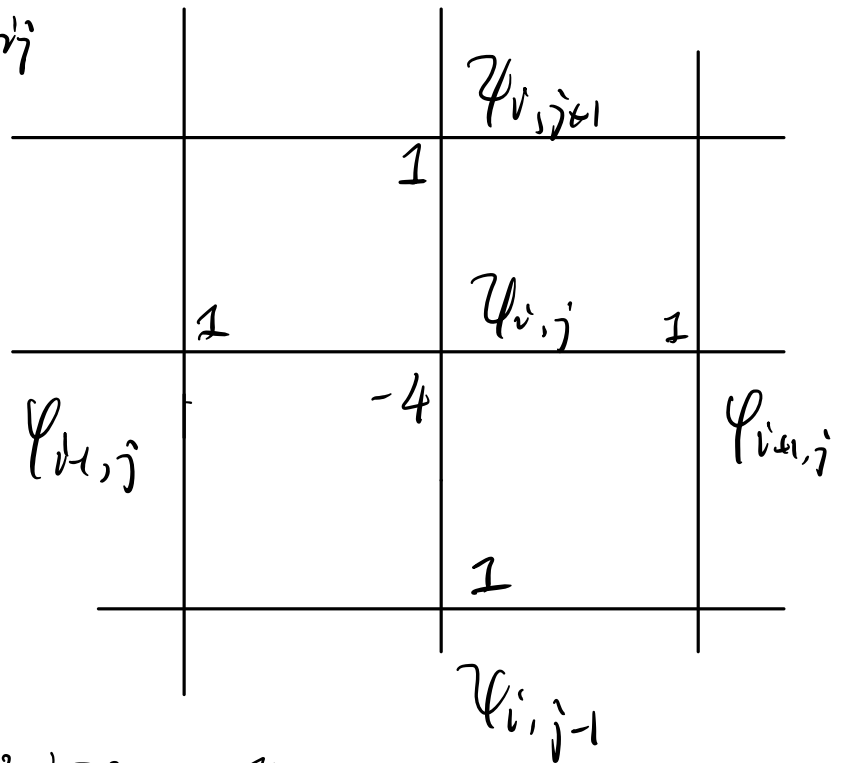
$$-\frac{\hbar^2}{2m} (\psi_{i+1,j} + \psi_{i-1,j} - 2\psi_{i,j} + \psi_{i,j+1} + \psi_{i,j-1} - 2\psi_{i,j})$$

$$+ V_{ij} \psi_{ij} = E \psi_{ij}$$

或者

$$- \frac{1}{\epsilon h^2} (\psi_{i+1,j} + \psi_{i-1,j} + \psi_{i,j+1} + \psi_{i,j-1} - 4 \psi_{i,j})$$

$$+ V_{ij} \psi_{ij} = E \psi_{ij}$$



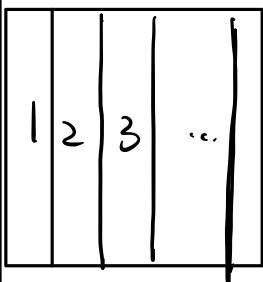
特点: 稀疏矩阵, 非对称阵.

② 矩阵维度 正比于 离散点 N^2 , ($N = \frac{X}{h}$)

eg. X 方向 100 个点, 矩阵维度

$$N = 10000$$

$\psi_{i,j} \rightarrow \text{Reshape}$

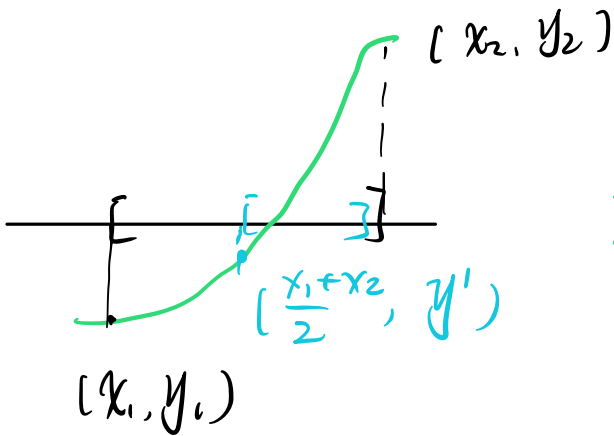
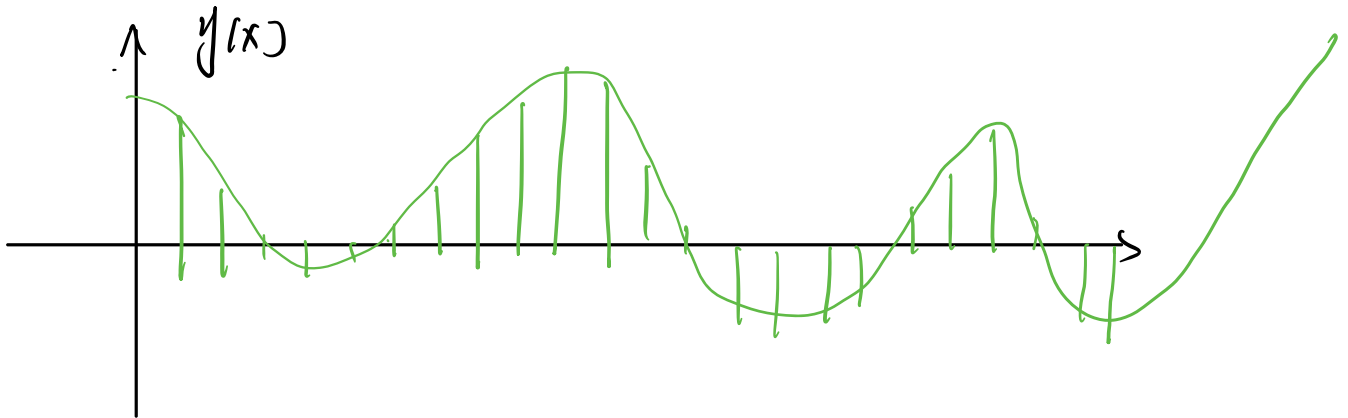
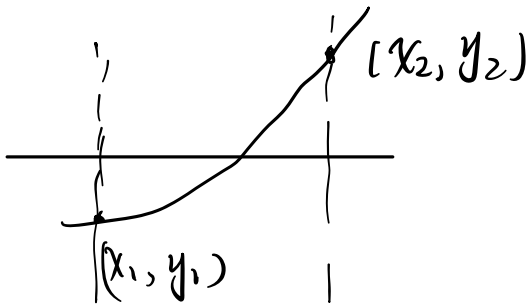


一个技巧

求根 (二分法)

$$x_1 x_2 < 0$$

系统至少存在两个解。

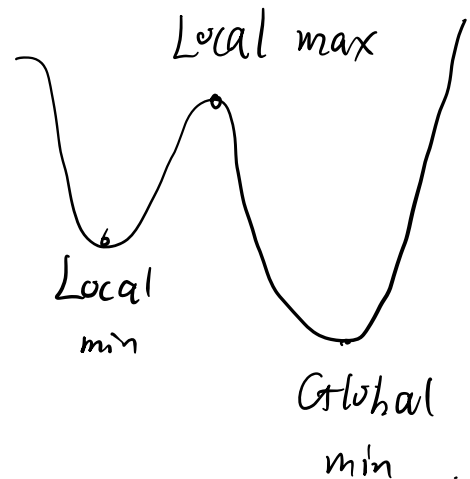


二分法。

不停寻找。还会讲牛顿法。

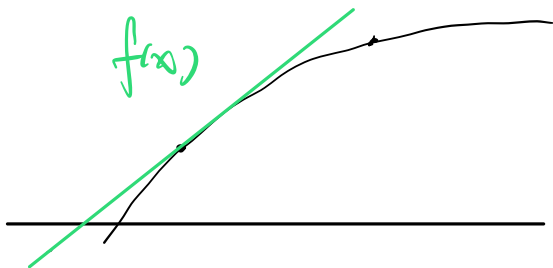
思想 | 求根
| 极值

$$\begin{cases} F(x) \text{ 极小值} \\ \frac{\partial F(x)}{\partial x} = 0 \end{cases}$$



牛顿法

$$f(x) = f(x_0) + f'(x_0)(x - x_0)$$



$$\begin{aligned} f(x) &= 0 \\ \Downarrow \\ f(x_0) + f'(x_0)(x - x_0) &= 0 \\ \Downarrow \\ x &= x_0 - \frac{f(x_0)}{f'(x_0)} \end{aligned}$$

$$\lim_{n \rightarrow \infty} x_n = x^* \Leftrightarrow x^* = x^* - \rho \frac{f(x^*)}{f'(x^*)}$$

$$\Leftrightarrow f(x^*) = 0$$

推广高维

$$f(\vec{x}) = f(\vec{x}_0) + \frac{\partial f}{\partial x^i} (x^i - x_0^i) = 0$$

$$\vec{x} = \vec{x}_0 + \left(\begin{array}{c} \\ \\ \end{array} \right)$$

求解 $\begin{cases} y' = f(x, y) \end{cases}$

$$\left. \begin{array}{l} \\ \end{array} \right\} \frac{\partial^2 u(\vec{x}, t)}{\partial t^2} = a^2 \nabla^2 u(\vec{x}, t)$$

龙格-库塔方法

mma 中 $\begin{cases} DSolve \\ NDSolve \end{cases}$

matlab ode

一个程序

ode45

$$y' = f(x, y)$$

$$y(x_{n+1}) = y(x_n) + f(x_n, y_n)h$$

$$(x_1, y_1)$$

↓

$$(x_2, y_2)$$

↓

⋮

$$x_{n+1} = x_n + h$$

自洽, 隐式的方法

$$y_{n+1} = y_n + \frac{h}{2} f(x_n, y_n) +$$

$$\frac{h}{2} f(x_{n+1}, y_n + hf(x_n, y_n))$$