

Review

精确对角化.

基矢 | 基矢而设

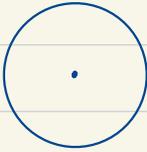
$$|\star \text{其它} \psi = \sum_n C_n \phi_n(x)$$

Lanczos 方法

可以处理大的系统 } 三对角矩阵. $N \sim 10^8$
 $g_0 \rightarrow g_1 \rightarrow g_2 \rightarrow \dots \rightarrow g_N$ } $O(N^2)$ | 不占内存

精确对角化中的对称性.

e.g.



有守恒量守恒, 或其它守恒量.

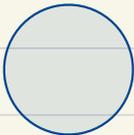
今天的例子: 柱对称的量子点能级, 80, 90年还十分重要.

P.S. 有磁轨道耦合; 和标引. 绝缘体有关.

这里利用 Bessel 函数求解

其它的基矢选择 可能要选择一些特殊函数.

首先我们考虑一个简单的情形



$$H = -\frac{\hbar^2}{2m} \left(\frac{d^2}{dx^2} + \frac{d^2}{dy^2} \right) + U(x, y), \quad U(x, y) \text{ 有旋转对称性}$$

$$\text{需要我们去解 } H\psi(x, y) = E\psi(x, y)$$

我们选择在坐标来处理这个问题.

$$\begin{aligned} H &= -\frac{\hbar^2}{2m} \left(\frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} \right) + U(r) \\ &= -\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} \right) + U(r) \end{aligned}$$

我们使用边界条件 $\psi(R, \theta) = 0$, 则方程可分离变量.

$$\psi(r, \theta) = \varphi_n(r) e^{in\theta} \quad \text{单值性.}$$

$$\Rightarrow H\psi = E\psi$$

$$\Rightarrow \left\{ \begin{aligned} &-\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} - \frac{n^2}{r^2} \right) \varphi_n(r) = E_n \varphi_n(r) \\ &\quad \quad \quad \uparrow \text{一个一维的方程.} \end{aligned} \right.$$

$$\varphi_n(r) = 0. \quad \varphi_n(r) \text{ 有界.}$$

无量纲化. 取 $r = \lambda x$

$$\Rightarrow -\frac{\hbar^2}{2m} \left(\frac{1}{\lambda^2} \frac{\partial^2}{\partial x^2} + \frac{1}{\lambda^2 x} \frac{\partial}{\partial x} - \frac{n^2}{\lambda^2 x^2} \right) \varphi_n = E_n \varphi_n$$

$$\text{取 } \lambda = \frac{\hbar^2}{2m E_n}$$

$$\Rightarrow -\left(\frac{\partial^2}{\partial x^2} + \frac{1}{x} \frac{\partial}{\partial x} - \frac{n^2}{x^2} \right) \varphi_n = \varphi_n$$

$$\Rightarrow x^2 \varphi_n'' + x \varphi_n' + (x^2 - n^2) \varphi_n = 0$$

Bessel function

$$x^2 f'' + x f' + (x^2 - \nu^2) f = 0$$

$f \sim J_\nu(x)$ 和 $Y_\nu(x)$
需类 奇=类

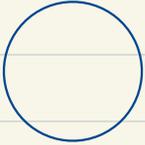
求解出本征值

$$\psi_n(\sqrt{\frac{\hbar}{2mE}} x) = 0 \Leftrightarrow J_n\left(\frac{R}{\sqrt{\frac{\hbar}{2mE}}}\right) = 0$$

at $r=R$

$$\Rightarrow \frac{R}{\sqrt{\frac{\hbar}{2mE}}} = X_{nj} \quad J_n(X_{nj}) = 0$$

★



\Rightarrow 可以求解. 按 n 分类.

★ Bessel function 可以做为基矢

$$\psi_{jn} = \sum_k C_{nk} J_n\left(\frac{\lambda_{nk} r}{R}\right)$$

可以在 $r=R$ 时, 天然满足 $\psi_{jn}(r=R) = 0$. 以及角动量守恒的条件.

Bessel function 正交性.

$$\int_0^R J_n(X_{nj} \frac{r}{R}) J_n(X_{nk} \frac{r}{R}) r dr = \begin{cases} 0, & j \neq k \\ \frac{R^2}{2} J_{n+1}^2(X_{nj}) \end{cases}$$

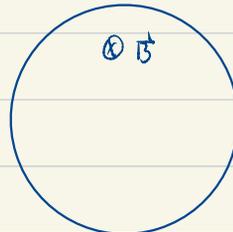
$$H = \begin{pmatrix} M + \frac{1}{2}k_x^2 + \frac{1}{2}k_y^2, & \alpha(k_x - i k_y) \\ \alpha(k_x + i k_y), & -M - \frac{1}{2}(k_x^2 + k_y^2) \end{pmatrix}$$

$$k_x = -i \frac{\partial}{\partial x}$$

$$k_y = -i \frac{\partial}{\partial y}$$

如果加磁场 $\vec{k} \rightarrow \vec{k} - e\vec{A}$

- 复原性
- ① 具有旋转对称性.
 - ② 2×2 矩阵
 - ③ SOC, 磁场



+ \vec{A}

选择 $\vec{A} = \frac{B}{2}(-y, x)$ 的对称规范, 柱坐标 $\begin{cases} y = r \sin \theta \\ x = r \cos \theta \end{cases}$

$$\begin{aligned} \textcircled{1} \quad (\vec{k} - e\vec{A})^2 &= k^2 + e^2 A^2 - e(\vec{k} \cdot \vec{A} + \vec{A} \cdot \vec{k}) & \begin{cases} p_x A_x = [p_x, A_x] + A_x p_x \\ p_y A_y = [p_y, A_y] + A_y p_x \end{cases} \\ &= -\nabla^2 + e^2 A^2 - ie(\vec{\nabla} \cdot \vec{A} + \vec{A} \cdot \vec{\nabla}) \\ &= -\left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2}\right) + \frac{e^2 B^2}{4} - ie \underbrace{\left[-y \frac{\partial}{\partial x} + x \frac{\partial}{\partial y}\right]} \\ &= -\left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2}\right) + \frac{e^2 B^2}{4} + ie B \frac{\partial}{\partial \theta} \end{aligned}$$

$$\textcircled{2} \quad k_x - i b y - e A_x + i e A_y = e^{-i\theta} \left[-i \left(\frac{\partial}{\partial r} - i \frac{1}{r} \frac{\partial}{\partial \theta} \right) + i \frac{B r e}{2} \right]$$

$$\begin{cases} H_{11} = M - \frac{1}{2} \left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} - \frac{B^2 e^2}{4} r^2 - i e B \frac{\partial}{\partial \theta} \right) \\ H_{12} = d e^{-i\theta} \left((-i) \left(\frac{\partial}{\partial r} - i \frac{1}{r} \frac{\partial}{\partial \theta} \right) + i \frac{B e r}{2} \right) \Rightarrow \text{引起角量的混合} \\ H_{21} = H_{12}^* \\ H_{22} = -M + \frac{1}{2} \left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} - \frac{B^2 e^2}{4} r^2 - i e B \frac{\partial}{\partial \theta} \right) \end{cases}$$

$$\Rightarrow \psi(r, \theta) = \begin{pmatrix} \psi_n(r) e^{i n \theta} \\ \psi_{n+1}(r) e^{i(n+1)\theta} \end{pmatrix}$$

