

$$T^2 \psi = T \tilde{\psi} = \tilde{\tilde{\psi}} = \lambda \psi, \quad \lambda = e^{i\theta}$$

$$T^2 = e^{i\theta} \Rightarrow U K U K = e^{i\theta} \Rightarrow U U^\dagger = e^{i\theta}$$

Review

分类 - Bott 周期律.

推导: 技巧

Stiefel manifold

短正合序列

$$A_1 \xrightarrow{f_1} A_2 \xrightarrow{f_2} A_3 \xrightarrow{f_3} A_4$$

主题: 为什么对称性很重要

$$M = G/H$$

单粒子

1) H 的分类

2) Topo 相变 \Rightarrow 能带

key: H 属于 () Lie 群 / Lie 代数。

Lie 群: A, KZ, 42
C, D, DI

History: Wigner 1951, 1958

Dyson 1962,

多体物理. \Rightarrow H \Rightarrow 随机矩阵.

矩阵元独立同分布

$$\begin{cases} \text{① } H = H^\dagger \\ \text{② 对称性不变.} \end{cases}$$

下面: Symmetry, 时间反演, 粒子数守恒, 宇称守恒
BdG equation, Bogoliubov - de - Gennes

$H_{BdG} \in \text{Lie Algebra } \mathfrak{g}$

ref. Atland, Zirnbauer PRB, 55, 1142 (1997)

$\langle \psi | \psi \rangle = \langle \hat{A} \psi | \hat{A} \psi \rangle \Rightarrow A \text{ is unitary or Anti-unitary}$

$A = uK$

Time-reversal $T = uK$ 且 $T^2 = \pm 1$

K , 复共轭算符.

$T^2 = 1$, eg $T = K, \sigma_x K$

$T^2 = -1$, eg $T = i\sigma_y K$

注意 $TH = HT$

OR $\begin{cases} T^2 = 1, \text{有什么后果} \\ T^2 = -1, \end{cases}$

实表示 $\Rightarrow H = u \lambda u^\dagger$
Kramers degeneracy \uparrow

U(N) Group

$T^2 = 1$



O(N) Group

$T^2 = 1, TH = HT$

$H|\phi\rangle = \lambda|\phi\rangle, \lambda \in \mathbb{R}$

$T H |\phi\rangle = \lambda T |\phi\rangle = H(T |\phi\rangle) = \lambda(T |\phi\rangle)$

H 本征值为 λ 时, 本征态是 $|\phi\rangle$ 和 $T|\phi\rangle$
怎么判定 $|\phi\rangle$ 和 $T|\phi\rangle$ 关系。

$\langle \phi | T \phi \rangle = \begin{cases} 1, & \text{本征值为单重.} \\ 0, & \text{本征值为两重.} \end{cases}$

$$|\psi\rangle = \begin{pmatrix} a_1 \\ \vdots \\ a_n \end{pmatrix}, |\phi\rangle = \begin{pmatrix} b_1 \\ \vdots \\ b_n \end{pmatrix}$$

区分.
 $\langle A^\dagger \phi | \psi \rangle \quad \langle \phi | A \psi \rangle$

$$\langle \phi | A \psi \rangle = (b_1^*, b_2^*, \dots, b_n^*) U \begin{pmatrix} a_1 \\ a_2 \\ \vdots \\ a_n \end{pmatrix} = \sum b_i^* U_{ij} a_j$$

$$\begin{aligned} |A^\dagger \phi\rangle &= (UK)^\dagger |\phi\rangle = K^\dagger U^\dagger |\phi\rangle = K(U^\dagger |\phi\rangle) \\ &= K(U_{ij}^* b_j) \\ &= K(b_j U_{ij}^*) \\ &= b_j^* U_{ij} \end{aligned}$$

$$\Rightarrow \langle A^\dagger \phi | \psi \rangle = b_j U_{ij} a_i = (\langle \phi | A \psi \rangle)^*$$

证明: $T^2 = -1 \Rightarrow |\phi\rangle \perp T|\phi\rangle$

$$\langle \phi | T\phi \rangle = 0.$$

$$\begin{aligned} \text{证 } \langle \phi | T\phi \rangle &= (\langle T^\dagger \phi | \phi \rangle)^* \\ &= -(\langle T\phi | \phi \rangle)^* \\ &= -\langle \phi | T\phi \rangle \end{aligned}$$

$$T^2 = -1 \Rightarrow \underline{\underline{T^\dagger = -T}}$$

$$\Rightarrow \langle \phi | T\phi \rangle = 0$$

$|\phi\rangle$ 和 $T|\phi\rangle$ 正交。

$$\left\{ \begin{array}{l} T^2 = -1 \Rightarrow \langle \phi | T\phi \rangle = 0, \text{ degenerate} \end{array} \right.$$

Summary

$A = \text{unitary}$

$$\langle \phi | A \psi \rangle = \langle A^\dagger \phi | \psi \rangle$$

$A = \text{Anti unitary}$

$$\langle \phi | A \psi \rangle = (\langle A^\dagger \phi | \psi \rangle)^*$$

$$\left. \right\} \tau^2 = 1 \Rightarrow T|\phi\rangle = |\phi\rangle$$

e.g. $T=K \Rightarrow K|\phi\rangle = |\phi\rangle \Rightarrow$ wavefunction 可取实数.

$T = \sigma_x K$

$$\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} x^* \\ y^* \end{pmatrix} = \begin{pmatrix} y^* \\ x^* \end{pmatrix} = \begin{pmatrix} x \\ y \end{pmatrix} \Leftrightarrow \begin{cases} x = x^* \\ y = y^* \end{cases} \Rightarrow \phi = \begin{pmatrix} x \\ x^* \end{pmatrix}$$

$T = \sigma_z K$

$$\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} x^* \\ y^* \end{pmatrix} = \begin{pmatrix} x^* \\ -y^* \end{pmatrix} = \begin{pmatrix} x^* \\ -y^* \end{pmatrix} = \begin{pmatrix} x \\ y \end{pmatrix} \Leftrightarrow \phi = \begin{pmatrix} x \\ -iy \end{pmatrix}$$

particle-hole

$$C = U_p K$$

$$\begin{cases} TH = HT \\ CH = -HC \end{cases}$$

$$HC|\phi\rangle = -E C|\phi\rangle$$

$$\begin{cases} C^2 = 1 & \text{e.g. } C = \tau_x K \\ C^2 = -1 & C = \tau_y K \end{cases}$$

$\sigma = \text{Spin } \uparrow \downarrow$
 $\tau = \text{particle-hole 空间}$

10 重简并: 为什么是 10

$$3 \times 3 \times 2 = 18$$

chiral symmetry S 手征 \Leftrightarrow 左手, 右手 $\underline{\gamma_5}$

$$\text{e.g. } \begin{cases} TH = HT \\ CH = -HC \end{cases}$$

$$\begin{cases} \{\gamma_5, H_b\} = 0 \\ \gamma_5^2 = 1 \end{cases}$$

$$S = CT = U_c K U_c^\dagger K$$

$$= U C U_T^*$$

为厄米算符

$$\begin{cases} C T H = C H T = -H C T \\ \delta H = -H S \end{cases}$$

$$U Y S U^* = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\begin{cases} \delta^2 H = H S^2 \\ \delta^+ S = 1 \end{cases}$$

$$\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} U H_0 U^* = -h \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$h = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

$$\downarrow \\ h = \begin{pmatrix} 0 & g \\ g & 0 \end{pmatrix}$$

$$H S |\phi\rangle = -E S |\phi\rangle$$

结论: ① 存在 T 对称, C 对称, 一定存在 S 对称性。

② $\begin{cases} \text{存在 T 和 S} \Rightarrow \text{一定存在 C} \\ \text{存在 C 和 S} \Rightarrow \text{一定存在 T} \end{cases}$

③ 例外 $\{S, H\} = 0, T^2 = 0, C^2 = 0$

④ $\delta^2 = 0, T^2 = 0, C^2 = 0$

10 重 T, C

$\begin{cases} 1) \text{ 两个都有} & 4 \\ 2) \text{ 只有一种} & 4 \\ 3) \text{ 只有 S} & 1 \\ 4) \text{ 都没有} & 1 \end{cases}$

BdG Equation 以及 H 的空间

$$H = \sum_{\alpha\beta} h_{\alpha\beta} C_{\alpha}^{\dagger} C_{\beta} + \frac{1}{2} \Delta_{\alpha\beta} C_{\alpha}^{\dagger} C_{\beta}^{\dagger} + \frac{1}{2} \Delta_{\beta\alpha}^* C_{\beta} C_{\alpha}$$

α, β index $\begin{cases} \text{Real space} \\ 1 \quad - \quad - \end{cases}$

1 R space

e.g.

$$H = \epsilon_1 c_1^\dagger c_1 + \epsilon_2 c_2^\dagger c_2 + \Delta c_1^\dagger c_2^\dagger + \Delta^* c_2 c_1 \\ + t c_1^\dagger c_2 + t c_2^\dagger c_1$$

$$= (c_1^\dagger \ c_2^\dagger \ c_1 \ c_2) \left(\begin{array}{cc|cc} \epsilon_1/2 & t/2 & 0 & \Delta/2 \\ t/2 & \epsilon_2/2 & -\Delta/2 & 0 \\ \hline 0 & -\Delta^*/2 & -\epsilon_1/2 & -t/2 \\ \Delta^*/2 & 0 & -t/2 & -\epsilon_2/2 \end{array} \right)$$

$$\begin{pmatrix} c_1 \\ c_2 \\ c_1^\dagger \\ c_2^\dagger \end{pmatrix}$$

$$+ \epsilon_1/2 + \epsilon_2/2$$

$$H_{\text{BdG}} = \frac{1}{2} (c^\dagger, c) \begin{pmatrix} h & \Delta \\ \Delta^\dagger & -h^\dagger \end{pmatrix} \begin{pmatrix} c \\ c^\dagger \end{pmatrix}$$

特点: 厄米性唯一.

2> 保证对角的后粒子层性不变.

$$H_{\text{BdG}} = \frac{1}{2} \begin{pmatrix} h & \Delta \\ \Delta^\dagger & -h^\dagger \end{pmatrix}$$

具有 C 对称性.

$$C = \bar{\sigma}_x K = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} K$$

$$CH = -HC \Rightarrow \bar{\sigma}_x H^* = -H \bar{\sigma}_x$$

$$\begin{aligned} & \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} h^T & \Delta^* \\ \Delta^T & -h \end{pmatrix} \\ &= - \begin{pmatrix} h & \Delta \\ \Delta^\dagger & h^T \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \end{aligned}$$

$\Rightarrow \underline{\Delta^T = -\Delta}$, 对于 BdG Equation 有 particle-hole 对称性.

D 类. ① Why we call it D class

② $H \rightarrow$ Lie Algebraic

$$\left. \begin{aligned} X = iH & \quad X^T = -X = -\bar{\sigma}_x X^T \bar{\sigma}_x \\ Y^T = -Y & = -\bar{\sigma}_x Y^T \bar{\sigma}_x \end{aligned} \right\} [X, Y] = Z$$