

Polarization

$$\nabla \cdot \vec{D} = -\rho(\pi) \quad \nabla \cdot (\partial_t \vec{P} - \vec{j}) = 0 \quad (\vec{E} = 0)$$

$$\Rightarrow \Delta P_\alpha = \int_0^T dt j_\alpha = e \int_0^T \sum_n dt \int_{Bt} \frac{d\vec{q}}{(2\pi)^d} \rho_{q,t}^1$$

$$= eva$$

(证明方法见沈佩清附录)

提示  $\vec{j}$  SSH  $Q = 0, \frac{1}{2}$

$$\frac{\partial P}{\partial t} + \nabla \cdot \vec{j} = 0$$

$$Q = \int P dx$$

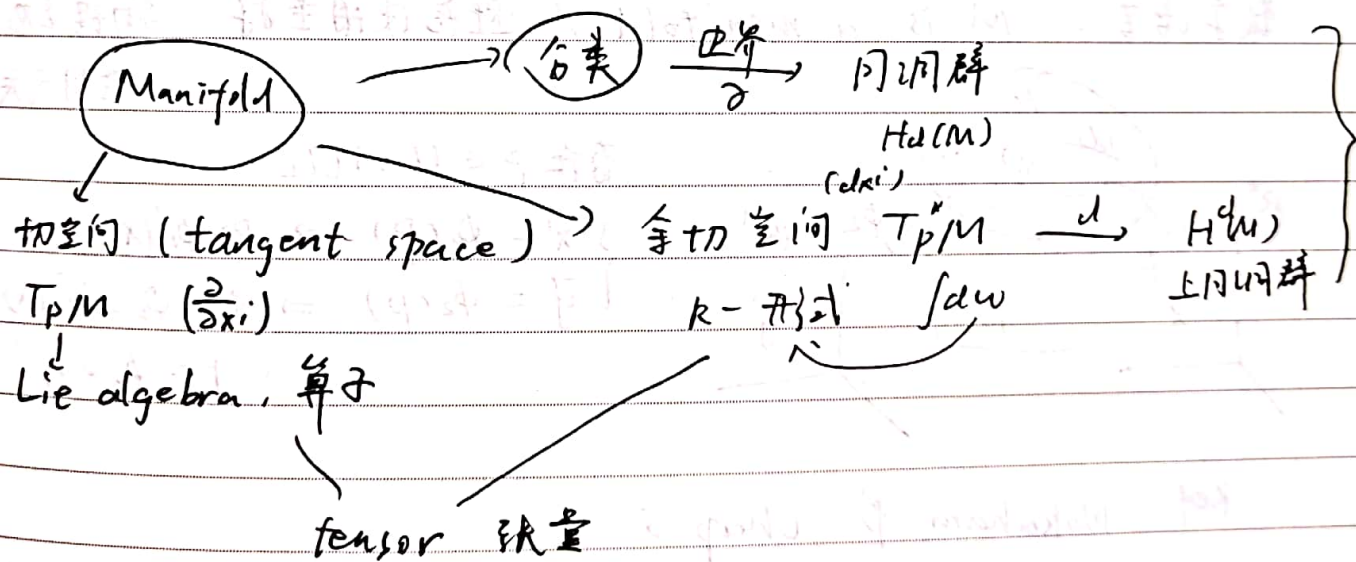
$$\dot{Q} = \int \dot{P} dx = \int \nabla \cdot \vec{j} dx = 0$$

$$Q = \frac{ne^1}{h} \quad n \in \mathbb{Z} \quad \int \vec{B} \cdot d\vec{s}' \in \mathbb{Z}$$

几何相

Zak phase  $\gamma_n = \frac{2\pi}{a} X_{nn}$  polarization

$\int d\omega$  不能整体使用 Stokes 定理



A. Zee

A tensor is sth. transform like a tensor.

Multilinear

for  $\{k\}$   $\langle v, d\omega \rangle = \langle \partial v, \omega \rangle$

$A \sim A + \nabla \phi$   $\omega \sim \omega + d\eta$   $(d\omega = d\omega + d^2\eta = d\omega)$



群  $\left\{ \begin{array}{l} \text{群表示论} \\ \text{群同态同构定理} \end{array} \right. \rightarrow \text{同伦群} \rightarrow \text{Lie群 / 对称性} \rightarrow \text{拓扑学}$   
 Noether 群元 = mapping Topo 分类 无天

问题：分类是什么

$\hookrightarrow$  同伦群 or 其它

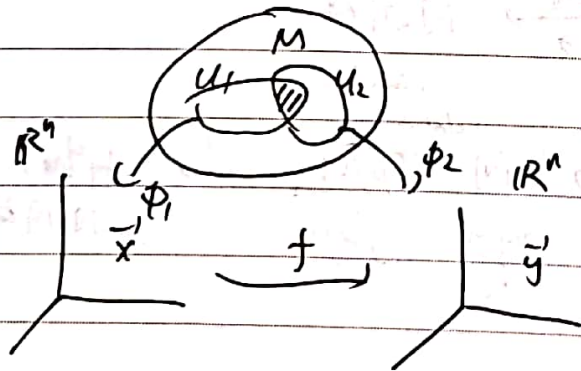
- 理解：
- 1) Topo defect
  - 2) Topo band
  - 3) BEC

### 1) Manifold

局部同胚于  $\mathbb{R}^n$  空间的空间

Homeomorphism

数学语言：M is a manifold / 避免使用坐标 物理规律的坐标无关



存在  $p \in U_1 \cap U_2$

$$\begin{cases} \bar{x} = \phi_1(p) \Rightarrow p = \phi_1^{-1}(\bar{x}) \\ \bar{y} = \phi_2(p) \Rightarrow y = \phi_2 \circ \phi_1^{-1}(\bar{x}) = f(\bar{x}) \end{cases} \Rightarrow f = \phi_2 \circ \phi_1^{-1}$$

Ref. Nakahara  $\S$  Chap 5

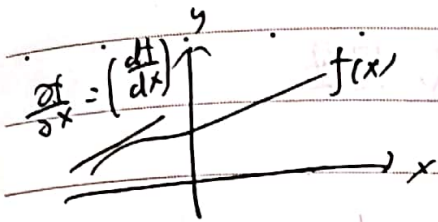
Gauge potential, knots, gravity Chap 3

切空间/切向量

$$f: M \rightarrow \mathbb{R} \quad f(x) \in \mathbb{R} \quad f \in C^\infty(M)$$

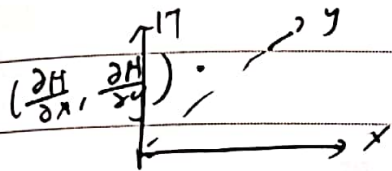
$\mathbb{R}^2$  平面





$\mathbb{R}^2$  平面

$H(x, y)$



定义  $\hat{V} = \frac{\partial}{\partial x}$   $\hat{V} f = \frac{\partial f}{\partial x}$

$\hat{V} : C^\infty(M) \rightarrow \mathbb{R}$

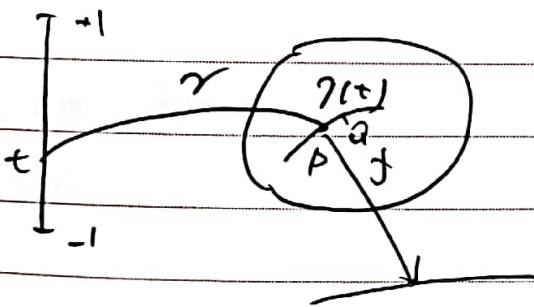
例:  $L_x = y \frac{\partial}{\partial x} - x \frac{\partial}{\partial y}$

$e^{a \frac{\partial}{\partial x}} f(x) \rightarrow f(x+a)$

物理意义

$\hat{V} = (\frac{\partial}{\partial x}, \frac{\partial}{\partial y})$   $\hat{V}(f) = (\frac{\partial}{\partial x} f, \frac{\partial}{\partial y} f)$

考虑  $M$  a manifold  $f : M \rightarrow \mathbb{R}$



$f(p) = f(\gamma(t)) = f \circ \gamma(t)$

$p = \gamma(t)$

$\frac{df}{dt} = \frac{d(f \circ \gamma)}{dt}$  全微分

$= f' \frac{d\gamma}{dt}$   $p = (x_1, \dots, x_n)$

$= \frac{\partial f}{\partial x^i} \frac{dx^i}{dt}$

$= (v^i \frac{\partial}{\partial x^i}) f = \hat{V}(f)$

$\hat{V} = \sum v^i \frac{\partial}{\partial x^i}$   $T_p M$  切空间

1)  $\hat{V}$  和全微分无关  $\hat{V} = \sum v^i \frac{\partial}{\partial x^i} = \sum v^i(y) \frac{\partial}{\partial y^i}$

但  $v^i(x) \rightarrow v^i(y)$  之间存在一个变换

2)  $v^i(p)$   $v^i : M \rightarrow \mathbb{R}$



3)  $\hat{V}$  是线性空间中的元素,  $V^i$  是系数,  $\frac{\partial}{\partial x^i}$  是基

条件:  $e_i$  线性独立

4)  $\frac{\partial}{\partial x^i}$  是线性独立的

线性  $x^i = A^{ij} y^j$        $\frac{\partial}{\partial x^i} = \sum A^{i0} y^j \frac{\partial}{\partial y^j}$

5) 加法性

6)  $\hat{V}, \hat{W} \in T_p M$

$\hat{V} + \hat{W} \in T_p M$        $[\hat{V}, \hat{W}] \in T_p M$

$\hat{V} \hat{M} \in T_p M$

### 余切空间 Cotangent space

$df = f_0 - f_p$  线性无关

$= \sum \left( \frac{\partial f}{\partial x^i} dx^i \right) \Leftrightarrow 1\text{-form} \Rightarrow$  写成线性空间形式

$= \sum \frac{\partial f}{\partial y^\alpha} \left( \frac{\partial x^i}{\partial y^\alpha} \right) \left( \frac{\partial x^i}{\partial y^\beta} \right) dy^\beta$

Basis  $dx^i$  系数为  $\left( \frac{\partial f}{\partial x^i} \right)$

$= \sum \frac{\partial f}{\partial y^\alpha} \left( \frac{\partial x^i}{\partial y^\alpha} \right) dy^\beta$

$= \sum \frac{\partial f}{\partial y^\alpha} dy^\alpha$

$dx^i$  线性无关

### 对偶关系

$\hat{V} = \sum V^i \frac{\partial}{\partial x^i} = (V^1, \dots, V^n) \in T_p(M)$

$\omega = \sum \frac{\partial f}{\partial x^i} dx^i = \left( \frac{\partial f}{\partial x^1}, \dots, \frac{\partial f}{\partial x^n} \right) \in T_p^*(M)$

$\omega(\hat{V})$  内积 inner product

$\langle a, b \rangle = a \cdot b = a^i e_i \cdot b^j e_j$        $e^i \cdot e_j = \delta^i_j$

另一组基  $\omega = df = \frac{\partial f}{\partial x^i} dx^i$        $V = V^i \frac{\partial}{\partial x^i}$

取  $e^i = dx^i$        $e_i = \frac{\partial}{\partial x^i}$

$\langle \omega, V \rangle = V^i \frac{\partial f}{\partial x^j} \langle dx^i, \frac{\partial}{\partial x^j} \rangle = V^i \frac{\partial f}{\partial x^i} \langle e^i, e_j \rangle = V^i \frac{\partial f}{\partial x^i}$

$\langle \omega, \hat{V} \rangle = V^i \frac{\partial f}{\partial x^i} = V(f) = \langle df, V \rangle$



应用: Schutz 数学物理中的几何方法 (Chap 5)  
微分形式的应用 / 物理

热力学:  $du = -p dv + T ds$   
 $d(du) = -dp \wedge dv + dT \wedge ds$

$p = p(V, S) \quad T = T(V, S)$   
 $-\frac{\partial p}{\partial S} = \frac{\partial T}{\partial V}$

数学:  $df = P dx + Q dy$   
 $\frac{\partial P}{\partial y} = \frac{\partial Q}{\partial x}$

电磁力学:  $w = \frac{1}{2} dp^i \wedge dq^i$   
 $dw = 0$

可以用  $T_p M$  做什么?  $\Rightarrow$  tensor  $\Rightarrow$  得证弦方程和生括无关

1) 积分

$\Rightarrow$  line integral, Surface integral, Volume integral

引入张量 (具体): 形式

0-form: $C^\infty(M)$ 函数	$\Omega^0(M)$	$C^n$	
1-form: $\sum f_i dx^i$	$\Omega^1(M)$	$C^n$	$T_p^*(M)$
2-form: $\sum_{i,j} f_{ij} dx^i \wedge dx^j$	$\Omega^2(M)$	$C^n$	$T_p^*(M) \otimes T_p^*(M)$
3-form: $\sum_{i,j,k} f_{ijk} dx^i \wedge dx^j \wedge dx^k$		$C^n$	

$(n-1)$ -form

$n$ -form  $f_{1 \dots n} dx^1 \wedge \dots \wedge dx^n \quad C^n$

$\Omega = \bigoplus_{k=0}^n \Omega^k(M) \quad (0, k)$  形式

线性  $A \otimes B \quad u \in A \quad v \in B \quad uv \in A \otimes B$

张量线性

$W^i V^j \rightarrow W^i \otimes V^j = (W^i \otimes V^j) \tilde{u}^a \tilde{v}^b$



$$dx^i \wedge dx^j \rightarrow w^i w^j \wedge dy^\alpha \wedge dy^\beta \quad \text{乘法: } \wedge$$

张量 (1,2) 维 tensor  $T_p M \otimes T_p M \otimes T_p^0 M$

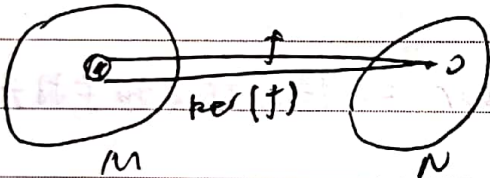
$$\left( v^\sigma \frac{\partial}{\partial x^\sigma} \right) \otimes \left( \frac{\partial f}{\partial x^i} dx^i \right) \otimes \left( \frac{\partial f}{\partial x^\alpha} dx^\alpha \right) = v^\sigma \frac{\partial f}{\partial x^i} \frac{\partial f}{\partial x^\alpha} \left( \frac{\partial}{\partial x^\sigma} \otimes dx^i \otimes dx^\alpha \right)$$

(0, k) 形式

$$\Omega^0(M) \xrightarrow{d} \Omega^1(M) \xrightarrow{d} \Omega^2(M) \xrightarrow{d} \Omega^3(M) \rightarrow \dots \rightarrow \Omega^n(M) \xrightarrow{d} 0$$

$$a = \frac{10^2}{h} \quad \oint \langle \psi | \nabla_R | \psi \rangle dR = \oint \langle \psi | d | \psi \rangle$$

线性代数



$$f([x]) = f(x) = y \quad [x] = \{x + x' \mid x' \in \ker f\}$$

$$\dim M = \dim(\text{im } f) + \dim(\ker f)$$

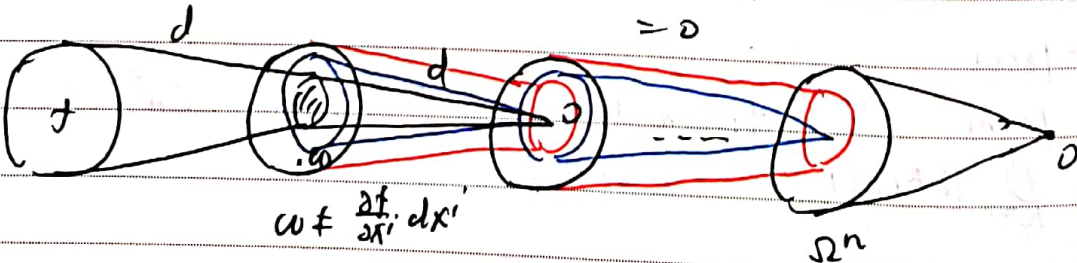
$$f \in C^0(M) \xrightarrow{d} df = \left( \frac{\partial f}{\partial x^i} \right) dx^i \xrightarrow{d} \dots \quad \text{大正合序列}$$

$$w = g^i da^i$$

$$\Omega^0 \rightarrow \Omega^1 \rightarrow \Omega^2 \rightarrow \dots$$

$$f \quad df = \frac{\partial f}{\partial x^i} dx^i \quad d(df) = 0$$

$$d\left(\frac{\partial f}{\partial x^i} dx^i\right) = \frac{\partial^2 f}{\partial x^i \partial x^j} dx^j \wedge dx^i + \frac{\partial^2 f}{\partial x^i \partial x^j} dx^i \wedge dx^j = 0$$



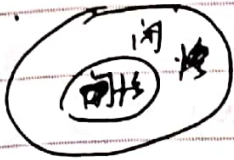
$$w = \frac{\partial f}{\partial x^i} dx^i$$

$$dw = 0$$

exact form : 恰当形式  $\alpha = d\eta \Rightarrow d^2\eta = 0$

closed form : 闭形式  $d\alpha = 0$





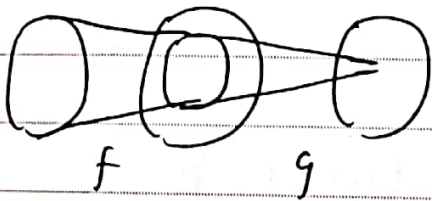
$$\alpha = P dx + Q dy$$

$$d\alpha = 0 \text{ 解 } \frac{\partial P}{\partial y} = \frac{\partial Q}{\partial x}$$

1) 正合序列

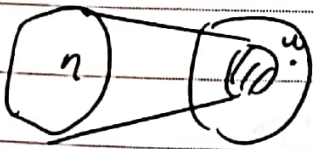
2) 为什么要做这个分类?

$$A \xrightarrow{f} B \xrightarrow{g} C$$



$$\text{im} f = \text{ker} g$$

$$\int_V d\omega = \langle V, d\omega \rangle = \langle \partial V, \omega \rangle = \langle \partial V, \omega \rangle + \langle \partial V, d\eta \rangle$$



群与同伦群

manifold



$$\bar{x} \quad f(\bar{x})$$

$$f: \bar{x} \rightarrow \mathbb{R}$$

两个空间

$$\left\{ \begin{array}{l} \text{切空间 } T_p(M) \Rightarrow \frac{\partial}{\partial x^i} \rightarrow \text{Lie 代数} \\ \text{余切... } T_p^*(M) \Rightarrow dx^i \rightarrow \text{Calculus} \end{array} \right.$$

余切  $\Rightarrow$

$$\left\{ \begin{array}{l} 1 \text{元} \text{ 线} \quad \oint \vec{F} \cdot d\vec{l} \quad 1 \text{ form} \in \Omega^1 \\ 2 \text{元} \text{ 面} \quad \oint \vec{B} \cdot d\vec{S} \quad \oint \vec{E} \cdot d\vec{S} \quad 2 \text{ form} \in \Omega^2 \\ 3 \text{元} \text{ 体} \quad \oint f dx dy dz \quad \oint f dV \quad 3 \text{ form} \in \Omega^3 \end{array} \right.$$

