

Chapter 9

Introduction to Singular Perturbation Theory

Section 9.1

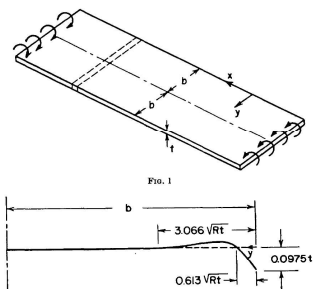
Roots of Polynomial Equations

9.1.0 Introduction

Boundary layer in Solid mechanics

Boundary layer in bent thin shells:

- H. Reissner 1912, Spannungen in Kugelschalen-fKuppeln, *Festschrift MueUer-Breslau*.
- Kelvin and Tait in 1867, Boussinesq in 1871. Love AEH: *A Treatise on the Mathematical Theory of Elasticity* (Cambridge, 1892)



- Fung, Y. C., & Wittrick, W. H. (1955). *A boundary layer phenomenon in the large deflexion of thin plates*. The Quarterly Journal of Mechanics and Applied Mathematics, 8(2), 191-210.

$$\zeta = -\frac{\mu t e^{-ay} (\cos ay - \sin ay)}{\sqrt{12(1-\mu^2)}}$$

boundary layer length $a = 1/\sqrt{Rt}$

9.1.0 Introduction

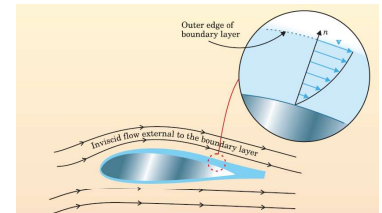
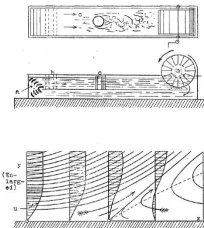
Boundary layer in Fluid mechanics

Prandtl Ludwig 1928. *Motion of fluids with very little viscosity*. Vier Abhandlungen zur Hydrodynamik und Aerodynamik.



Prandtl's paper gave the first description of the boundary-layer concept in fluid dynamics.

- The viscous effects only exist in a thin region (boundary layer) near the surface.
- Outside the boundary layer, the flow was essentially the inviscid flow.



9.1.1 Roots of Polynomial Equations

Example: regular quadratic

$$x^2 + \epsilon x - 1 = 0$$

ϵ : a small constant, say $\epsilon = 0.0000001$.

Exact solutions $x = \frac{-\epsilon \pm \sqrt{\epsilon^2 + 4}}{2}$

As $\epsilon=0$, we have the unperturbed solution

$$x^2 + \epsilon x - 1 = 0 \Rightarrow x = \pm 1$$

9.1.1 Roots of Polynomial Equations

Taylor expansion of the exact solution:

$$x = \begin{cases} 1 - \frac{\epsilon}{2} + \frac{\epsilon^2}{8} - \frac{\epsilon^4}{128} + O(\epsilon^6) \\ -1 - \frac{\epsilon}{2} - \frac{\epsilon^2}{8} + \frac{\epsilon^4}{128} + O(\epsilon^6) \end{cases} \quad \begin{array}{l} \text{converge if} \\ |\epsilon| < 2 \end{array}$$

Series method

power series expansion around $x = \pm 1$

$$x = \pm 1 + a_1 \epsilon + a_2 \epsilon^2 + a_3 \epsilon^3 + \dots$$

The same Taylor expansions can be reproduced.

Example: singular quadratic

$$\epsilon x^2 + x - 1 = 0$$

ϵ is a small constant, say $\epsilon = 0.0000001$.

Exact solutions $x = \frac{-1 \pm \sqrt{1+4\epsilon}}{2\epsilon}$

As $\epsilon \rightarrow 0$, we have the unperturbed solution

~~$$\epsilon x^2 + x - 1 = 0 \Rightarrow x = 1$$~~

Only one root !

$$\epsilon^{-1}(a_{-1}^2 + a_{-1}) + \epsilon^0(2a_{-1}a_0 + a_0 - 1) + \epsilon(2a_{-1}a_1 + a_0^2 + a_1) + \dots = 0$$

Comparing coefficients of ϵ of same order

$$\epsilon^{-1}: a_{-1}^2 + a_{-1} = 0 \quad a_{-1} = -1 \quad \text{or} \quad a_{-1} = 0$$

$$\epsilon^0: 2a_{-1}a_0 + a_0 - 1 = 0 \quad a_0 = -1 \quad a_0 = 1$$

$$\epsilon^1: 2a_{-1}a_1 + a_0^2 + a_1 = 0 \quad a_1 = 1 \quad a_1 = -1$$

$$x = -\frac{1}{\epsilon} - 1 + \epsilon + \dots$$

Singular root

$$x = 1 - \epsilon + 2\epsilon^2 \dots$$

Regular root

(2) ϵx^2 and -1 is comparable, assuming x is smaller than other two terms.

$$\epsilon x^2 + x - 1 = 0 \Rightarrow |x| \approx \frac{1}{\sqrt{\epsilon}} \gg 1$$

$$\epsilon x^2 \sim O(1) \ll x \sim O(\epsilon^{-1/2})$$

Bad guess

■ Taylor expansion of the exact solution:

$$x = \begin{cases} 1 - \epsilon + 2\epsilon^2 - 5\epsilon^3 + O(\epsilon^4) & \text{converge if} \\ -\frac{1}{\epsilon} - 1 + \epsilon - 2\epsilon^2 + 5\epsilon^3 + O(\epsilon^4) & |\epsilon| < 1/4 \end{cases}$$

■ Expansion method

Assuming power series

Why? be patient

$$x = \frac{a_{-1}}{\epsilon} + a_0 + a_1\epsilon + a_2\epsilon^2 + \dots$$

Substituting into $\epsilon x^2 + x - 1 = 0$

■ Balancing terms

balance the three terms $\epsilon x^2 + x - 1 = 0$

(1) x and -1 is comparable, assuming ϵx^2 is smaller than other two terms.

$$\epsilon x^2 + x - 1 = 0 \Rightarrow x \approx 1$$

$$\epsilon x^2 \sim O(\epsilon) \ll x \sim O(1)$$

Good guess

(3) ϵx^2 and x is comparable, assuming both terms $\gg 1$.

$$\epsilon x^2 + x - 1 = 0 \Rightarrow |x| \approx \frac{1}{\epsilon} \gg 1$$

$$\epsilon x^2 \sim O(\epsilon^{-1}) \quad x \sim O(\epsilon^{-1}) \gg O(1)$$

Self consistent

When ϵx^2 and x balance, x is very large $x \sim O(\epsilon^{-1})$

rescale x , $x = \frac{X}{\epsilon}$ with $X \sim O(1)$

We get a regular looking problem

$$X^2 + X - \varepsilon = 0$$

Using regular expansion

$$X = a_0 + a_1\varepsilon + a_2\varepsilon^2 + \dots$$

Comparing coefficients of ε of same order, we get

$$x = -\frac{1}{\varepsilon} - 1 + \varepsilon + \dots$$

Consider a quadratic equation

$$m^2 + \varepsilon m - 4 = 0, \quad 0 < \varepsilon \ll 1$$

Naïve simplification leads to

$$m^2 - 4 = 0 \rightarrow m_0 \approx \pm 2$$

assume a series solution around $m_0 = \pm 2$

$$m = \pm 2 + m_1\varepsilon^1 + m_2\varepsilon^2 + \dots$$

We can find

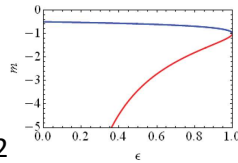
$$m_1 = 2 - \frac{1}{4}\varepsilon + \dots, \quad m_2 = -2 - \frac{3}{4}\varepsilon + \dots$$

Consider another quadratic equation

$$\varepsilon m^2 + 2m + 1 = 0, \quad 0 < \varepsilon \ll 1$$

Naïve simplification leads to

$$2m + 1 = 0 \rightarrow m_0 \approx -\frac{1}{2}$$



assume a series solution around $m_0 = -1/2$

$$m = -\frac{1}{2} + m_1\varepsilon^1 + m_2\varepsilon^2 + \dots$$

We can find

$$m = -\frac{1}{2} - \frac{1}{8}\varepsilon^1 - \frac{1}{16}\varepsilon^2 + \dots \quad \text{One root lost}$$

The naively approximated problem has a different qualitative nature from the original problem.

The Naïve simplification implies the root we are look for is

$$m \sim O(1), \quad \text{as } \varepsilon \rightarrow 0 \quad \varepsilon m^2 \sim O(\varepsilon)$$

so we can find the root $m_0 \approx -1/2$.

But, we may lost the second root if

$$m \rightarrow \infty, \quad \text{as } \varepsilon \rightarrow 0 \quad \varepsilon m^2 \gg 1$$

Let's compare the magnitude of three terms systemically

$$\varepsilon m^2 + 2m + 1 = 0$$

Case 0. All three terms have the same magnitude

- all three terms are important
- we have to solve the original full equation.

Case 1. terms 1 ~ term 3

$$\varepsilon m^2 + 1 = 0 \rightarrow m = i\varepsilon^{-1/2}$$

Check: $\varepsilon m^2 + 2m + 1 = 0$

$$\left| i2\varepsilon^{-1/2} \right| \gg 1 \quad \begin{cases} \bullet \text{ bigger than the rest two.} \\ \bullet \text{ Inconsistent.} \end{cases}$$

Case 2. terms 1 ~ term 2

$$\varepsilon m^2 + 2m = 0 \rightarrow m = -2\varepsilon^{-1}, \quad m = 0$$

$$1 \ll \left| \varepsilon m^2 \right| \approx 4\varepsilon^{-1} \quad 1 \ll \left| 2m \right| \approx 4\varepsilon^{-1} \quad \text{consistent}$$

root $m=0 \rightarrow$ inconsistent

To improve the accuracy of root m , we use iterative method.

$$F(x) = G(x) \quad m \rightarrow \infty$$

$$\varepsilon m^2 + 2m = 0 \quad \rightarrow \quad \varepsilon m + 2 = -m^{-1}$$

$$m_n = -2\varepsilon^{-1} - \varepsilon^{-1}m_{n-1}^{-1} \quad \leftarrow \quad m = -2\varepsilon^{-1} - \varepsilon^{-1}m^{-1}$$

$$m_1 = -\frac{2}{\varepsilon} - \frac{1}{\varepsilon m_0} = -\frac{2}{\varepsilon} - \frac{1}{\varepsilon(-2\varepsilon^{-1})} = -\frac{2}{\varepsilon} + \frac{1}{2}$$

$$m_2 = -\frac{2}{\varepsilon} - \frac{1}{\varepsilon(-2\varepsilon^{-1} + 2^{-1})} = -\frac{2}{\varepsilon} + \frac{1}{2 - \varepsilon/2}$$

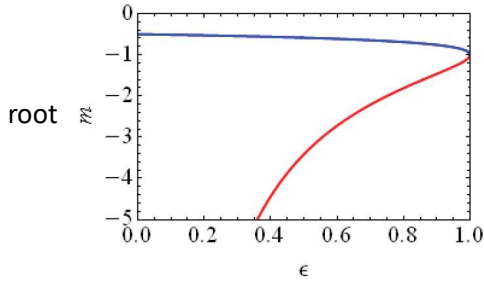
$$= -\frac{2}{\varepsilon} + \frac{1}{2} + \frac{1}{8}\varepsilon + \dots$$

9.1.1 Roots of Polynomial Equations

This indicates we may use series expansion to look for the two roots

$$m_1 = -\frac{1}{2} + m_1 \epsilon + \dots$$

$$m_2 = -\frac{2}{\epsilon} + m_0 + m_1 \epsilon + \dots$$

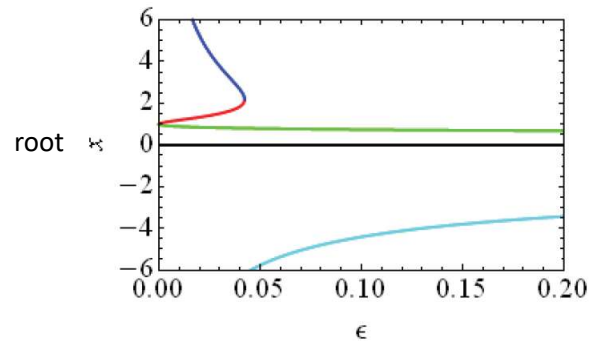


9.1.2 A more complicated problem

A more complicated problem

Consider a quartic equation

$$\epsilon x^4 + \epsilon x^3 - x^2 + 2x - 1 = 0, \quad 0 < \epsilon \ll 1$$



9.1.2 A more complicated problem

Example 1. Find approximations to those roots $\sim O(1)$

Assume naively

$$x = 1 + x_1 \epsilon + O(\epsilon^2)$$



$$\epsilon x^4 + \epsilon x^3 - x^2 + 2x - 1 = 0$$

$$\epsilon (1 + x_1 \epsilon + O(\epsilon^2))^4 + \epsilon (1 + x_1 \epsilon + O(\epsilon^2))^3 - (1 + x_1 \epsilon + O(\epsilon^2))^2 + 2(1 + x_1 \epsilon + O(\epsilon^2)) - 1 = 0$$



$$\epsilon (2 + O(\epsilon)) - x_1 \epsilon^2 + O(\epsilon^3) = 0$$



$$\epsilon^1: 2 = 0 \Rightarrow x = 1 + x_1 \epsilon + O(\epsilon^2)$$

Assumption is naive

9.1.2 A more complicated problem

Let's try with iterative method:

$$F(x) = G(x)$$

$$\epsilon x^4 + \epsilon x^3 - x^2 + 2x - 1 = 0 \rightarrow (x-1)^2 = \epsilon x^4 + \epsilon x^3$$

We can reach for the two roots iteratively

$$x_0 = 1, \quad x_n = 1 \pm \sqrt{\epsilon x_{n-1}^4 + \epsilon x_{n-1}^3}$$

For example

$$x_0 = 1, \quad x_1 = 1 + \sqrt{\epsilon x_0^4 + \epsilon x_0^3} = 1 + \sqrt{2\epsilon} \epsilon^{1/2}$$

$$x_2 = 1 + \sqrt{2\epsilon} \epsilon^{1/2} + \frac{7}{2} \epsilon + 2\sqrt{2\epsilon} \epsilon^{3/2} + \epsilon^2$$

Therefore we may use regular series expansion

$$x = 1 + a_1 \epsilon^{1/2} + a_2 \epsilon + a_3 \epsilon^{3/2} + \dots$$

Instead of $x = 1 + x_1 \epsilon + O(\epsilon^2)$

9.1.2 A more complicated problem

balancing procedure: look for the rest two roots

- The case that the first two terms are negligible has been considered.
- We consider the following balance:

$$\epsilon x^4 + \epsilon x^3 - x^2 + 2x - 1 = 0$$

(a) $\epsilon x^4 + \epsilon x^3 = 0$ (b) $\epsilon x^4 - x^2 = 0$

(c) $\epsilon x^4 + 2x = 0$ (d) $\epsilon x^4 - 1 = 0$

(e) $\epsilon x^3 - x^2 = 0$ (f) $\epsilon x^3 + 2x = 0$

(g) $\epsilon x^3 - 1 = 0$

9.1.2 A more complicated problem

For example, check case (a)

(a) $\epsilon x^4 + \epsilon x^3 = 0 \rightarrow x = 0$ $x = -1$

$$\epsilon x^4 + \epsilon x^3 - x^2 + 2x - 1 = 0$$

$$\epsilon x^4 = O(\epsilon) \quad \epsilon x^3 = O(\epsilon)$$

$$x^2 = O(1) \quad 2x = O(1) \quad 1 = O(1)$$

inconsistency

Check case (b)

(b) $\epsilon x^4 - x^2 = 0$

$p^{(0)} = \frac{1}{\sqrt{\epsilon}}, \quad n^{(0)} = -\frac{1}{\sqrt{\epsilon}}$

$\epsilon x^4 + \epsilon x^3 - x^2 + 2x - 1 = 0$

$\epsilon x^4 = O(\epsilon^{-1}) \quad x^2 = O(\epsilon^{-1})$

$\epsilon x^3 = O(\epsilon^{-1/2}) \quad 2x = O(\epsilon^{-1/2}) \quad 1 = O(1)$

consistency

Check all cases for consistency

	x	ϵx^4	ϵx^3	$-x^2$	$2x$	-1
$\epsilon x^4 \sim \epsilon x^3$	$O(1)$	$O(\epsilon)$	$O(\epsilon)$	$O(1)$	$O(1)$	$O(1)$
$\epsilon x^4 \sim x^2$	$O(\epsilon^{-1/2})$	$O(\epsilon^{-1})$	$O(\epsilon^{-1/2})$	$O(\epsilon^{-1})$	$O(\epsilon^{-1/2})$	$O(1)$
$\epsilon x^4 \sim 2x$	$O(\epsilon^{-1/3})$	$O(\epsilon^{-1/3})$	$O(1)$	$O(\epsilon^{-2/3})$	$O(\epsilon^{-1/3})$	$O(1)$
$\epsilon x^4 \sim 1$	$O(\epsilon^{-1/4})$	$O(1)$	$O(\epsilon^{-1/4})$	$O(\epsilon^{-1/2})$	$O(\epsilon^{-1/4})$	$O(1)$
$\epsilon x^3 \sim x^2$	$O(\epsilon^{-1})$	$O(\epsilon^{-3})$	$O(\epsilon^{-2})$	$O(\epsilon^{-2})$	$O(\epsilon^{-1})$	$O(1)$
$\epsilon x^3 \sim 2x$	$O(\epsilon^{-1/2})$	$O(\epsilon^{-1})$	$O(\epsilon^{-1/2})$	$O(\epsilon^{-1})$	$O(\epsilon^{-1/2})$	$O(1)$
$\epsilon x^3 \sim 1$	$O(\epsilon^{-1/3})$	$O(\epsilon^{-1/3})$	$O(1)$	$O(\epsilon^{-2/3})$	$O(\epsilon^{-1/3})$	$O(1)$

- (a), (c)-(g) are all inconsistent.
- Only (b) is consistent:

iterative method for more accurate roots:

$\epsilon x^4 + \epsilon x^3 - x^2 + 2x - 1 = 0 \rightarrow \epsilon x^2 - 1 = x^{-2} - 2x^{-1} - \epsilon x$

$\epsilon [x^{(0)}]^2 - 1 = 0 \quad F(x) = G(x)$

$\epsilon [x^{(n)}]^2 - 1 = [x^{(n-1)}]^{-2} - 2[x^{(n-1)}]^{-1} - \epsilon [x^{(n-1)}]$

$p^{(0)} = \frac{1}{\sqrt{\epsilon}}, \quad p^{(1)} = \frac{1}{\sqrt{\epsilon}} - \frac{3}{2} + \dots$

$n^{(0)} = -\frac{1}{\sqrt{\epsilon}}, \quad n^{(1)} = -\frac{1}{\sqrt{\epsilon}} - \frac{3}{2} + \dots$

p and n can be expanded in power series of $\sqrt{\epsilon}$

Section 9.2

Boundary Value Problems for Ordinary Differential Equations

For quadratic equation

$\epsilon m^2 + 2m + 1 = 0, \quad 0 < \epsilon \ll 1$

root-1 $m_1 = -\frac{1}{2} - \frac{1}{8}\epsilon + \dots \sim O(1)$

$\epsilon m^2 + 2m + 1 = 0$ correct scaling

↑
Indicate the magnitude

root-2 $m_2 = -\frac{2}{\epsilon} + \frac{1}{2} + \frac{1}{8}\epsilon + \dots \sim O(\epsilon^{-1})$

To get correct scaling, introduce $m = \epsilon^{-1}\hat{m}$

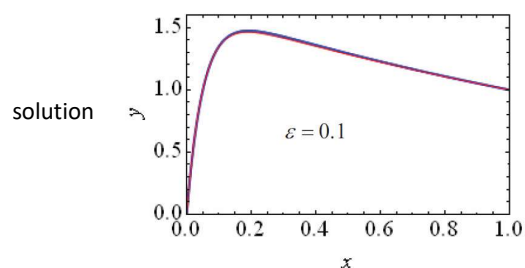
$\hat{m}^2 + 2\hat{m} + \epsilon = 0 \quad \hat{m} \sim O(1)$

no single scale in singular perturbation problem.

Consider a second order ordinary differential equation

$\epsilon \frac{d^2 y}{dx^2} + 2 \frac{dy}{dx} + y = 0, \quad 0 < x < 1, \quad 0 < \epsilon \ll 1$

B.C. $y(0) = 0, \quad y(1) = 1$



Consider a second order ordinary differential equation

$$\varepsilon \frac{d^2 y}{dx^2} + 2 \frac{dy}{dx} + y = 0, \quad 0 < x < 1, \quad 0 < \varepsilon \ll 1$$

B.C. $y(0) = 0, \quad y(1) = 1$

Naïve simplification

$$2 \frac{dy}{dx} + y = 0 \rightarrow y = K \exp\left(-\frac{1}{2}x\right)$$

$$y(0) = 0 \rightarrow y \equiv 0$$

$$y(1) = 1 \rightarrow y = \exp\left(\frac{1}{2} - \frac{x}{2}\right)$$

Which boundary conditions should be used?

We see in this BVP problem:

- singular features similar to those of the algebraic equations discussed above.
- neither simplified solutions can satisfy both boundary conditions
- Perhaps, the simplified solutions is a good in certain subdomain ?

Exact solution:

$$\varepsilon m^2 + 2m + 1 = 0 \quad \text{Characteristic equation}$$

$$y = Ae^{m_1 x} + Be^{m_2 x}$$

$$m_{1,2} = \left(-1 \pm \sqrt{1 - \varepsilon}\right) / \varepsilon$$

$$\begin{cases} y(0) = A + B = 0 \\ y(1) = Ae^{m_1} + Be^{m_2} = 1 \end{cases} \rightarrow y = \frac{\exp(m_1 x) - \exp(m_2 x)}{\exp(m_1) - \exp(m_2)}$$

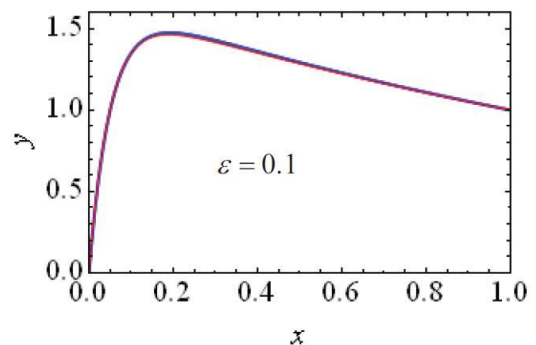
Approximate solution:

$$m_1 \approx -\frac{1}{2}, \quad m_2 \approx -\frac{2}{\varepsilon} \quad y(x, \varepsilon) \approx e^{1/2} \left(e^{-x/2} - e^{-2x/\varepsilon} \right)$$

$$y = \frac{\exp(m_1 x) - \exp(m_2 x)}{\exp(m_1) - \exp(m_2)}$$

$$y \approx e^{1/2} \left(e^{-x/2} - e^{-2x/\varepsilon} \right)$$

Very good approximate solution



Exam the behaviors of the approximate solution:

$$y(x, \varepsilon) \approx e^{1/2} \left(e^{-x/2} - e^{-2x/\varepsilon} \right)$$

- B.C. at $x = 0$ is satisfied

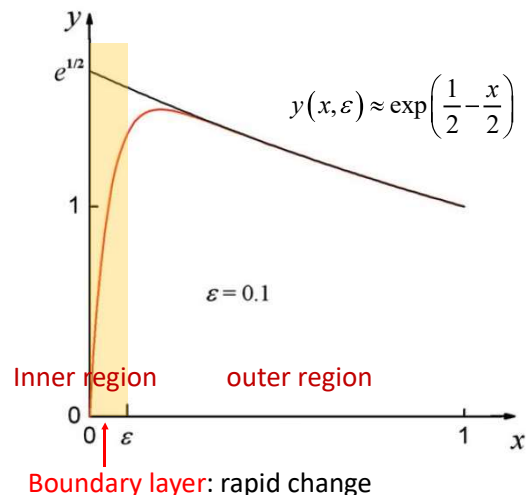
$$y(0) \approx e^{1/2} \left(e^{-x/2} - e^{-2x/\varepsilon} \right) \Big|_{x=0} = 0$$

- when $x \gg \varepsilon$, the second term decays rapidly.

$$y(x, \varepsilon) \approx e^{1/2} \left(e^{-x/2} - \cancel{e^{-2x/\varepsilon}} \right) = \exp\left(\frac{1}{2} - \frac{x}{2}\right)$$

Coincident with one of the naïve solutions

$$y(1) = 1 \rightarrow y = \exp\left(\frac{1}{2} - \frac{x}{2}\right)$$



9.2.1 Examination of the exact solution to a model problem

Consider the case as $\varepsilon < 0$

$$\varepsilon \frac{d^2 y}{dx^2} + 2 \frac{dy}{dx} + y = 0, \quad 0 < x < 1, \quad 0 < |\varepsilon| \ll 1, \quad \varepsilon < 0$$

B.C. $y(0) = 0, \quad y(1) = 1$



$$e \frac{d^2 y}{dx^2} - 2 \frac{dy}{dx} - y = 0, \quad 0 < e \ll 1, \quad e = -\varepsilon > 0$$

B.C. $y(0) = 0, \quad y(1) = 1$

9.2.1 Examination of the exact solution to a model problem

$$em^2 - 2m - 1 = 0$$

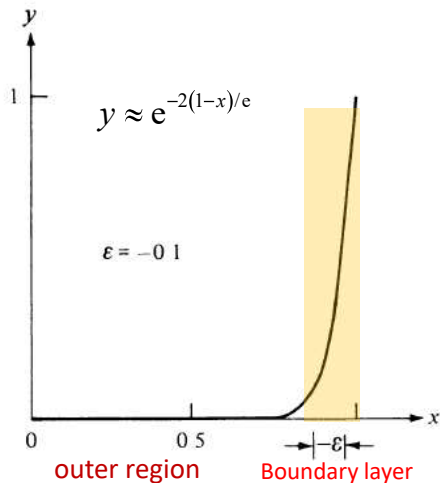
$$m_1 \approx -\frac{1}{2}, \quad m_2 \approx \frac{2}{e} = -\frac{2}{\varepsilon} > 0$$

$$y = \frac{e^{-x/2} - e^{2x/e}}{e^{-1/2} - e^{2/e}} \approx -e^{-2/e} (e^{-x/2} - e^{2x/e}) \approx e^{-2(1-x)/e}$$

Coincident with one of the naïve solutions as $\varepsilon > 0$

$$y \approx e^{-2(1-x)/e}$$

9.2.1 Examination of the exact solution to a model problem



9.2.1 Examination of the exact solution to a model problem

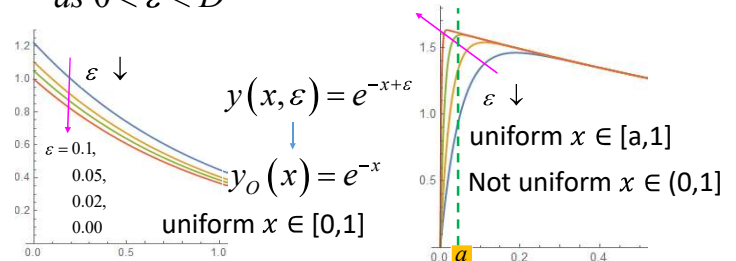
uniform limit 一致收敛

We say $\lim_{\varepsilon \rightarrow 0} y(x, \varepsilon) = y_0(x), x \in [a, 1], a > 0$

is an uniform limit, if

$$\forall E > 0, \exists D(E) > 0, |y(x, \varepsilon) - y_0(x)| < E,$$

as $0 < \varepsilon < D$



9.2.1 Examination of the exact solution to a model problem

Consider the following limit

$$y_0(x) = \exp\left[\frac{1}{2}(1-x)\right]$$

$$\lim_{x \rightarrow 0} \left[\lim_{\varepsilon \rightarrow 0} y(x, \varepsilon) \right] = \lim_{x \rightarrow 0} \left[\lim_{\varepsilon \rightarrow 0} e^{1/2} (e^{-x/2} - e^{-2x/\varepsilon}) \right]$$

$$= \lim_{x \rightarrow 0} e^{1/2} e^{-x/2} = \lim_{x \rightarrow 0} y_0(x) = y_0(0) = e^{1/2}$$

but

$$\lim_{\varepsilon \rightarrow 0} \left[\lim_{x \rightarrow 0} y(x, \varepsilon) \right] = \lim_{\varepsilon \rightarrow 0} \left[\lim_{x \rightarrow 0} e^{1/2} (e^{-x/2} - e^{-2x/\varepsilon}) \right]$$

$$= \lim_{\varepsilon \rightarrow 0} 0 = 0$$

The limit operator can be interchanged only if

$$\lim_{\varepsilon \rightarrow 0} y(x, \varepsilon) = y_0(x), x \in (0, b), b > 0$$

is an uniform limit

9.2.1 Examination of the exact solution to a model problem

Summary

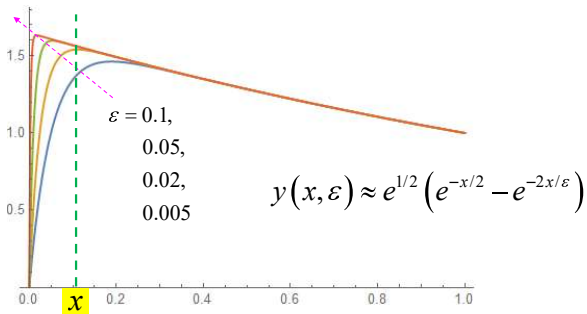
- the outer solution $y_0(x)$ does not satisfy the boundary condition at $x = 0$.
- $y_0(x)$ is the limit of $y(x, \varepsilon)$ in $(0, 1]$ as $\varepsilon \rightarrow 0$
- $y_0(x)$ is not the uniform limit of $y(x, \varepsilon)$ in $(0, 1]$ as $\varepsilon \rightarrow 0$

what is an appropriate approximation for small ε inside the boundary layer?

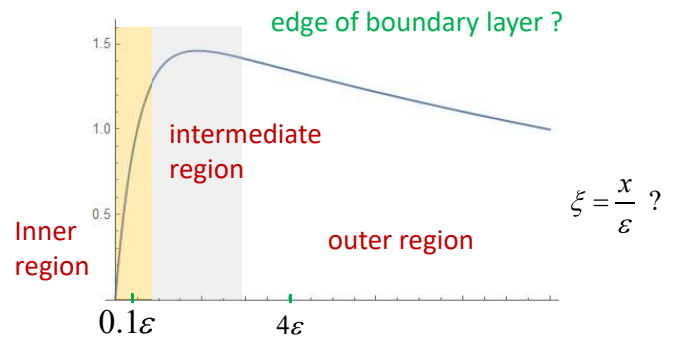
- Inner solution

9.2.1 Examination of the exact solution to a model problem

For a fixed x close to the origin, it will be out of the boundary layer as ϵ small enough



9.2.1 Examination of the exact solution to a model problem



The so-called "edge" of boundary layer is **obscure**, but

- $x=0.1\epsilon$ is well inside the boundary layer
- $x=4\epsilon$ is well outside the boundary layer

9.2.1 Examination of the exact solution to a model problem

Inner solution $y_I(\xi)$

Define $\xi = \frac{x}{\epsilon}$, for x in boundary layer

then $y(x, \epsilon) = y(\xi\epsilon, \epsilon) = Y(\xi, \epsilon)$

Inner solution $y_I(\xi) = \lim_{\epsilon \rightarrow 0} Y(\xi, \epsilon)$
 ξ fixed

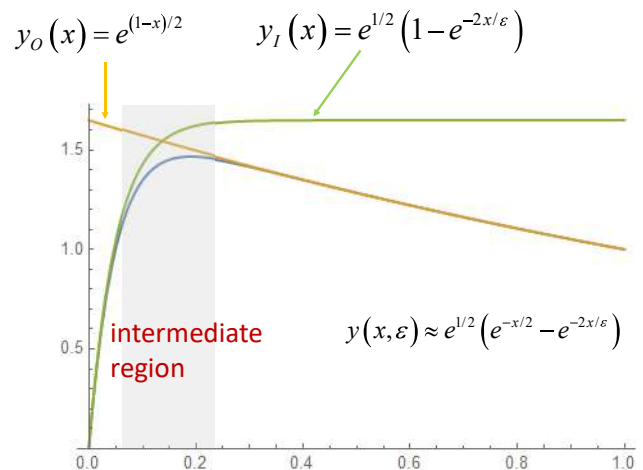
In our problem:

$$y(x, \epsilon) \approx e^{1/2} (e^{-x/2} - e^{-2x/\epsilon}) = e^{1/2} (e^{-\xi/2} - e^{-2\xi}) = Y(\xi, \epsilon)$$

$$y_I(\xi) = \lim_{\epsilon \rightarrow 0} Y(\xi, \epsilon) = e^{1/2} (1 - e^{-2\xi})$$

ξ fixed

9.2.1 Examination of the exact solution to a model problem



9.2.2 Finding an approximate solution by singular perturbation methods

Let's solve ODE

$$\epsilon \frac{d^2 y}{dx^2} + 2 \frac{dy}{dx} + y = 0, \quad 0 < x < 1, \quad 0 < \epsilon \ll 1$$

B.C. $y(0) = 0, \quad y(1) = 1$

• outer solution

$$2 \frac{dy_o}{dx} + y_o = 0, \quad y_o(1) = 1 \quad \rightarrow \quad y_o(x) = e^{(1-x)/2}$$

• inner solution

Let $\xi = \frac{x}{\delta}$ $y(x, \epsilon) = y(\xi\delta, \epsilon) = Y(\xi, \epsilon)$

thus $\frac{\epsilon}{\delta^2} \frac{d^2 Y}{d\xi^2} + \frac{2}{\delta} \frac{dY}{d\xi} + Y = 0$

9.2.2 Finding an approximate solution by singular perturbation methods

Pairwise balance:

$$\frac{\epsilon}{\delta^2} \frac{d^2 Y}{d\xi^2} + \frac{2}{\delta} \frac{dY}{d\xi} + Y = 0$$

$\delta = \sqrt{\epsilon}$ $O(1)$ $O(\epsilon^{-1/2})$ $O(1)$

inconsistent

$$\frac{\epsilon}{\delta^2} \frac{d^2 Y}{d\xi^2} + \frac{2}{\delta} \frac{dY}{d\xi} + Y = 0$$

$\delta = \epsilon$ $O(\epsilon^{-1})$ $O(\epsilon^{-1})$ $O(1)$

consistent

$$\frac{d^2 Y}{d\xi^2} + 2 \frac{dY}{d\xi} + \epsilon Y = 0$$

$$y_I(\xi) = \lim_{\epsilon \rightarrow 0} Y(\xi, \epsilon)$$

ξ fixed

$$\frac{d^2 y_I}{d\xi^2} + 2 \frac{dy_I}{d\xi} = 0 \quad y_I(0) = 0$$

inner solution $y_I(x) = C(1 - e^{-2\xi})$

9.2.3 Matching

The scales of three sub-domain :

- Inner region $x = O(\delta)$
 - Intermediate region $x = O[\Theta(\varepsilon)]$
 - Outer region $x = O(1)$
- $$\left. \begin{array}{l} x = O(\delta) \\ x = O[\Theta(\varepsilon)] \\ x = O(1) \end{array} \right\} \begin{array}{l} \lim_{\varepsilon \rightarrow 0} \frac{\Theta}{\delta} = \infty \\ \lim_{\varepsilon \rightarrow 0} \frac{\Theta}{1} = 0 \end{array}$$

In intermediate region, x can be rescaled by

$$\eta = \frac{x}{\Theta(\varepsilon)} \quad \longrightarrow \quad \begin{array}{l} x = \eta\Theta \\ \xi = \frac{x}{\delta} = \frac{\eta\Theta}{\delta} \end{array}$$

The inner and outer solutions matches if

$$\lim_{\varepsilon \rightarrow 0} \left[y_o(x) \Big|_{x=\eta\Theta} \right] = \lim_{\varepsilon \rightarrow 0} \left[y_i(\xi) \Big|_{\xi=\eta\Theta/\delta} \right], \quad \eta \text{ fixed}$$

9.2.3 Matching

In the present problem

$$\left. \begin{array}{l} \lim_{\substack{\varepsilon \rightarrow 0 \\ \eta \text{ fixed}}} \left[y_o(x) \Big|_{x=\eta\Theta} \right] = \lim_{\substack{\varepsilon \rightarrow 0 \\ \eta \text{ fixed}}} e^{(1-\eta\Theta)/2} = e^{1/2} \\ \lim_{\substack{\varepsilon \rightarrow 0 \\ \eta \text{ fixed}}} \left[y_i(\xi) \Big|_{\xi=\eta\Theta/\delta} \right] = \lim_{\substack{\varepsilon \rightarrow 0 \\ \eta \text{ fixed}}} C(1 - e^{-2\eta\Theta/\delta}) = C \end{array} \right\} C = e^{1/2}$$

We find the uniform approximation for the whole domain

$$\begin{aligned} y_U(x) &= y_i\left(\frac{x}{\delta}\right) + y_o(x) - \lim_{\varepsilon \rightarrow 0} y_o(\eta\Theta) \\ &= e^{1/2} (1 - e^{-2x/\varepsilon}) + e^{(1-x)/2} - e^{1/2} \\ &= e^{1/2} (e^{-x/2} - e^{-2x/\varepsilon}) \end{aligned}$$

9.2.3 Matching

Another way for the uniform approximation

$$\begin{aligned} y_U(x) &= y_i\left(\frac{x}{\delta}\right) y_o(x) / \lim_{\varepsilon \rightarrow 0} y_o(\eta\Theta) \\ &= e^{1/2} (1 - e^{-2x/\varepsilon}) e^{(1-x)/2} / e^{1/2} \\ &= e^{1/2} \left(e^{-\frac{x}{2}} - e^{-\frac{2x}{\varepsilon} - \frac{x}{2}} \right) \\ &\approx e^{1/2} \left(e^{-\frac{x}{2}} - e^{-\frac{2x}{\varepsilon}} \right) \end{aligned}$$

9.2.4 Further examples

Example 1. as $\varepsilon < 0$, we know the BL is at $x=1$. But if we still pretend BL locates at $x=0$, matching will be impossible

$$e \frac{d^2 y}{dx^2} - 2 \frac{dy}{dx} - y = 0, \quad 0 < e \ll 1, \quad e = -\varepsilon > 0$$

B.C. $y(0) = 0, \quad y(1) = 1$

9.2.4 Further examples

• outer solution

$$2 \frac{dy_o}{dx} + y_o = 0, \quad y_o(1) = 1 \quad \longrightarrow \quad y_o(x) = e^{(1-x)/2}$$

• inner solution

Let $\xi = \frac{x}{\delta}$ $y(x, e) = y(\xi\delta, e) = Y(\xi, e)$

thus $\frac{e}{\delta^2} \frac{d^2 Y}{d\xi^2} - \frac{2}{\delta} \frac{dY}{d\xi} - Y = 0$

$\delta = e$ $O(e^{-1}) \quad O(e^{-1}) \quad O(1)$

9.2.4 Further examples

$$\frac{d^2 Y}{d\xi^2} - 2 \frac{dY}{d\xi} - eY = 0$$

$$\frac{d^2 y_i}{d\xi^2} - 2 \frac{dy_i}{d\xi} = 0$$

$$y_i(x) = C(1 - e^{2\xi})$$

$$y_i(\xi) = \lim_{\substack{e \rightarrow 0 \\ \xi \text{ fixed}}} Y(\xi, e)$$

$$y_i(0) = 0$$

In intermediate region, x can be rescaled by

$$\eta = \frac{x}{\Theta(e)} \quad \longrightarrow \quad x = \eta\Theta \quad \xi = \frac{e}{\delta} = \frac{\eta\Theta}{\delta}$$

9.2.4 Further examples

$$\lim_{\epsilon \rightarrow 0} \frac{\Theta}{\delta} = \infty \quad \lim_{\epsilon \rightarrow 0} \frac{\Theta}{1} = 0$$

$$\left. \begin{aligned} \lim_{\substack{\epsilon \rightarrow 0 \\ \eta \text{ fixed}}} \left[y_o(x) \Big|_{x=\eta\Theta} \right] &= \lim_{\substack{\epsilon \rightarrow 0 \\ \eta \text{ fixed}}} e^{(1-\eta\Theta)/2} = e^{1/2} \\ \lim_{\substack{\epsilon \rightarrow 0 \\ \eta \text{ fixed}}} \left[y_I(\xi) \Big|_{\xi=\eta\Theta/\delta} \right] &= \lim_{\substack{\epsilon \rightarrow 0 \\ \eta \text{ fixed}}} C(1 - e^{-2\eta\Theta/\delta}) = \infty \end{aligned} \right\} \infty \neq e^{1/2}$$

matching is impossible

9.2.4 Further examples

Example 2. as $\epsilon < 0$, we know the BL is at $x=1$.

$$e \frac{d^2 y}{dx^2} - 2 \frac{dy}{dx} - y = 0, \quad 0 < e \ll 1, \quad e = -\epsilon > 0$$

B.C. $y(0) = 0, \quad y(1) = 1$

9.2.4 Further examples

- outer solution

$$2 \frac{dy_o}{dx} + y_o = 0, \quad y_o(0) = 0 \quad \rightarrow \quad y_o(x) \equiv 0$$

- inner solution

Let $\xi = \frac{1-x}{\delta}$ $y(x, e) = y(\xi\delta, e) = Y(\xi, e)$

thus $\frac{e}{\delta^2} \frac{d^2 Y}{d\xi^2} + \frac{2}{\delta} \frac{dY}{d\xi} - Y = 0$

$\delta = e$ $O(e^{-1}) \quad O(e^{-1}) \quad O(1)$

9.2.4 Further examples

$$\frac{d^2 Y}{d\xi^2} + 2 \frac{dY}{d\xi} - eY = 0$$



$$\frac{d^2 y_I}{d\xi^2} + 2 \frac{dy_I}{d\xi} = 0 \quad y_I(1) = 1$$

$$y_I(\xi) = \lim_{\substack{\epsilon \rightarrow 0 \\ \xi \text{ fixed}}} Y(\xi, e)$$

$$y_I(x) = C(e^{-2\xi} - 1) + 1$$

In intermediate region, x can be rescaled by

$$\eta = \frac{1-x}{\Theta(e)} \quad \rightarrow \quad \begin{aligned} x &= 1 - \eta\Theta \\ \xi &= \frac{1-x}{\delta} = \frac{\eta\Theta}{\delta} \end{aligned}$$

9.2.4 Further examples

$$\lim_{\epsilon \rightarrow 0} \frac{\Theta}{\delta} = \infty \quad \lim_{\epsilon \rightarrow 0} \frac{\Theta}{1} = 0$$

$$\left. \begin{aligned} \lim_{\substack{\epsilon \rightarrow 0 \\ \eta \text{ fixed}}} \left[y_o(x) \Big|_{x=1-\eta\Theta} \right] &= 0 \\ \lim_{\substack{\epsilon \rightarrow 0 \\ \eta \text{ fixed}}} \left[y_I(\xi) \Big|_{\xi=\eta\Theta/\delta} \right] &= \lim_{\substack{\epsilon \rightarrow 0 \\ \eta \text{ fixed}}} C \left(e^{-2\eta\Theta/\delta} - 1 \right) + 1 = C - 1 \end{aligned} \right\} C = 1$$

the uniform approximation

$$\begin{aligned} y_U(x) &= y_I \left(\frac{1-x}{\delta} \right) + y_o(x) - \lim_{\epsilon \rightarrow 0} y_o(\eta\Theta) \\ &= e^{-2(1-x)/e} = e^{2(1-x)/\epsilon} \end{aligned}$$

9.2.4 Further examples

General procedure:

- for singular perturbation BVP for 2nd ODE
- works often but may fail

Consider the boundary value problem

$$\epsilon y'' + f(x, y, y') = 0, \quad y(a) = A, \quad y(b) = B, \quad \epsilon > 0$$

Suppose boundary layer locates at $x=a$

(a) Determine the outer approximation y_o by solving

$$f(x, y, y') = 0, \quad y(b) = B$$

(b) Introduce the boundary layer variable

$$\xi = \pm \frac{x-a}{\delta(\varepsilon)}$$

where the sign is chosen to make ξ positive

With the notation $Y(\xi, \varepsilon) = y(a \pm \xi\delta, \varepsilon)$

$$\frac{\varepsilon}{\delta^2} \frac{d^2 Y}{d\xi^2} + f\left(a \pm \xi\delta, Y, \pm \frac{1}{\delta} \frac{dY}{d\xi}\right) = 0$$

Suppose the second term has the form $\delta^s F\left(\xi, Y, \frac{dY}{d\xi}\right)$

$$\text{Balance: } \frac{\varepsilon}{\delta^2} = \delta^s \rightarrow \delta = \varepsilon^{1/(2+s)} \quad s > -2$$

(c) Determine $y_I(\xi)$ by solving

$$\frac{d^2 y_I}{d\xi^2} + F\left(\xi, y_I, \frac{dy_I}{d\xi}\right) = 0, \quad y_I(0) = A$$

(c) To find the arbitrary constant in the above solution introduce the intermediate variable

$$\eta = \pm \frac{x-a}{\Theta(\varepsilon)}$$

$$\lim_{\varepsilon \rightarrow 0} \frac{\Theta}{\delta} = \infty \quad \lim_{\varepsilon \rightarrow 0} \frac{\Theta}{1} = 0$$

(c) Determine $y_I(\xi)$ by solving

$$\frac{d^2 y_I}{d\xi^2} + F\left(\xi, y_I, \frac{dy_I}{d\xi}\right) = 0, \quad y_I(0) = A$$

(c) To find the arbitrary constant in the above solution introduce the intermediate variable

$$\eta = \pm \frac{x-a}{\Theta(\varepsilon)} \quad \lim_{\varepsilon \rightarrow 0} \frac{\Theta}{\delta} = \infty \quad \lim_{\varepsilon \rightarrow 0} \frac{\Theta}{1} = 0$$

Impose the matching requirement

$$\lim_{\varepsilon \rightarrow 0} \left[y_O(x) \Big|_{x=a \pm \eta\Theta} \right] = \lim_{\varepsilon \rightarrow 0} \left[y_I(\xi) \Big|_{\xi=\eta\Theta/\delta} \right], \quad \eta \text{ fixed}$$

Example 3. Find outer and inner approximations for small positive ε to the solution of

$$\varepsilon \frac{d^2 y}{dx^2} + \alpha(x) \frac{dy}{dx} + \beta(x)y = 0, \quad \text{BL: } x = 0$$

$$y(0) = 0, \quad y(1) = 1, \quad \alpha(0) \equiv a_0 > 0, \quad \beta(0) \text{ finite}$$

• outer solution

$$\alpha(x) \frac{dy}{dx} + \beta(x)y = 0, \quad y(1) = 1$$

$$y_O(x) = \exp\left[\int_x^1 \frac{\beta}{\alpha} ds\right]$$

• inner solution

$$\text{Let } \xi = \frac{x}{\delta} \quad \delta(\varepsilon) \rightarrow 0 \text{ as } \varepsilon \rightarrow 0$$

$$\frac{\varepsilon}{\delta^2} \frac{d^2 Y}{d\xi^2} + \frac{\alpha(\delta\xi)}{\delta} \frac{dY}{d\xi} + \beta(\delta\xi)Y = 0$$

$$\delta = \varepsilon \quad O(\varepsilon^{-1}) \quad O(\varepsilon^{-1}) \quad O(1)$$

$$\frac{d^2 y_I}{d\xi^2} + \alpha_0 \frac{dy_I}{d\xi} = 0, \quad y_I(0) = 0$$

$$y_I(x) = C(1 - e^{-\alpha_0 \xi}) + 1$$

Matching

$$C = \exp\left[\int_0^1 \frac{\beta}{\alpha} ds\right]$$